

$$1) \frac{\alpha_p \beta_T}{\alpha_V} = \frac{1}{P} \Rightarrow \frac{RT}{P^2 V} = \frac{1}{P} \Rightarrow \underline{\underline{P = \frac{RT}{V}}}$$

2) C 3) B

$$4) \alpha_L = \frac{1}{L} \left(\frac{\partial L}{\partial T} \right)_P \Rightarrow \alpha_L = \frac{\Delta L/L}{\Delta T} \quad \text{original length}$$

$$\Rightarrow \Delta L = L \cdot \alpha_L \Delta T$$

$$\Rightarrow t_2 = 2\pi \sqrt{\frac{L + L\alpha_L \Delta T}{g}} = 2\pi \sqrt{\frac{L(1 + \alpha_L \Delta T)}{g}} = t_1 \sqrt{1 + \alpha_L \Delta T}$$

$$= 1s \sqrt{1 + 20 \cdot 10^{-6} \cdot 10} \approx \underline{\underline{1.0001 s}}$$

5) B, C, E 6) B, C 7) A

$$8) L_f = P \cdot t = 500 \frac{J}{s} \cdot 15s = 500 \cdot 15 J = 7500 J = \underline{\underline{7.5 kJ}}$$

9) Adiabats $PV^\gamma = \text{konst}$



$$W = \int_{V_1}^{V_2} P dV = \int_{V_1}^{V_2} \frac{\text{konst}}{V^\gamma} dV = \text{konst} \left[\frac{V^{1-\gamma}}{1-\gamma} \right]_{V_1}^{V_2} = \frac{\text{konst}}{1-\gamma} \left[\frac{V_2}{V_2^\gamma} - \frac{V_1}{V_1^\gamma} \right]$$

$$= \underline{\underline{\frac{1}{1-\gamma} [P_2 V_2 - P_1 V_1]}}$$

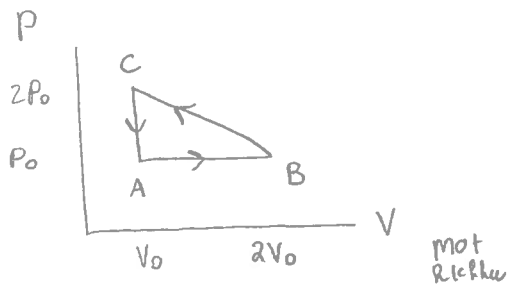
10) H_2 : Diatomisch $f = f_{\text{trans}} + f_{\text{rot}} + f_{\text{vib}} = 3 + 2 + 2$

Ved 40K bidrar kun translationen $\Rightarrow \underline{\underline{f = 3}}$

$$\Rightarrow C_V = \underline{\underline{\frac{3}{2} R}}$$

11) $\underline{\underline{e_R = \frac{T_1}{T_2 - T_1}}}$ 12) Carnetmaskin $\Rightarrow \underline{\underline{\eta = 1 - \frac{T_1}{T_2} = \frac{T_2 - T_1}{T_2}}}$

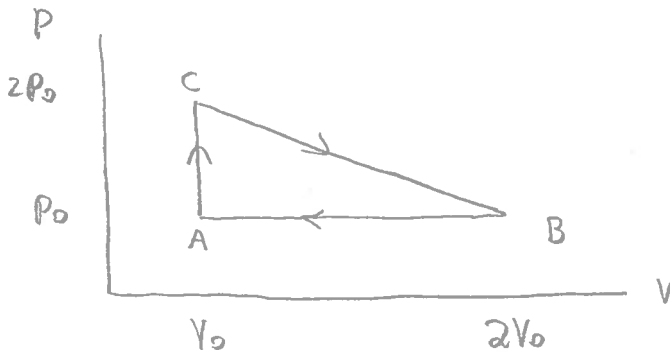
13)



$$Q_t = \Delta U_t + W_t = W_t = \frac{1}{2} V_0 P_0 \frac{1}{2} = \underline{\underline{-\frac{1}{2} P_0 V_0}}$$

$L > 0$ per syklus

14)



Carnots teorem: $\eta < \eta_c = 1 - \frac{T_{\text{Lavest}}}{T_{\text{Højest}}}$

Vi må bestemme laveste og højeste temperatur.

$$P = \frac{nRT}{V}$$

Laveste temperatur: $T_A \Rightarrow nRT_A = P_0 V_0$

Højeste temp: et eller andet sted på linje fra C til B.

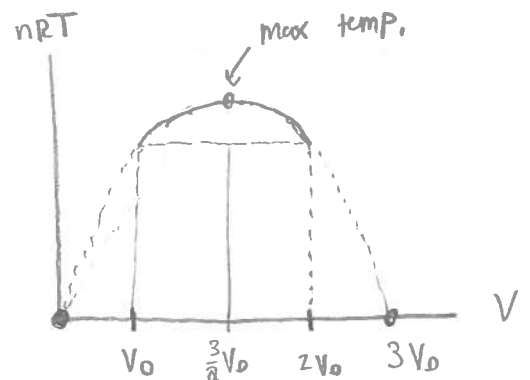
Parametrisering: $P_{CB} = -\frac{P_0}{V_0} V_{BC} + 3P_0$ (rett linje giv 2 punkt)

Temperatur på linje CB:

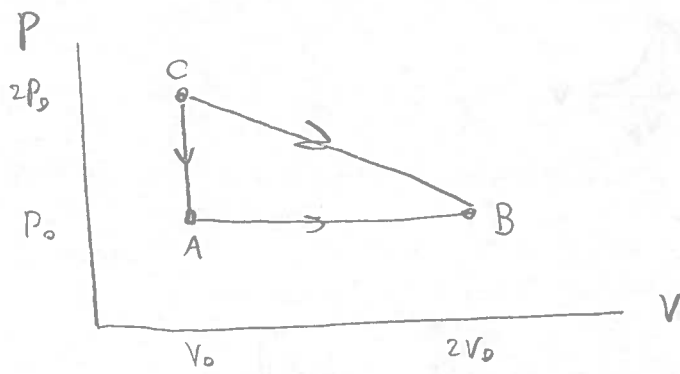
$$nRT = PV = -\frac{P_0}{V_0} V^2 + 3P_0 V = V \left(3P_0 - \frac{P_0}{V_0} V \right) \Rightarrow V = 3V_0$$

$$nRT_{\text{max}} = \frac{3}{2} V_0 \left(3P_0 - \frac{3}{2} P_0 \right) = \frac{3}{2} V_0 \left(\frac{3}{2} P_0 \right) = \frac{9}{4} P_0 V_0$$

$$\eta_c = 1 - \frac{P_0 V_0}{\frac{9}{4} P_0 V_0} = 1 - \frac{4}{9} = \frac{5}{9} \approx \underline{\underline{0.56}}$$



15)



$$pV = nRT$$

$$\frac{P}{nR} = \frac{T}{V}$$

$$ds = C_V \frac{dT}{T} + \underbrace{\left(\frac{\partial P}{\partial T}\right)_V}_{\frac{nR}{V}} dV$$

$$\Rightarrow \Delta S = C_V \ln \frac{T_{\text{slutt}}}{T_{\text{start}}} + nR \ln \frac{V_{\text{slutt}}}{V_{\text{start}}}$$

$$CA: \text{Isoker, } \Delta T < 0 \Rightarrow \underline{\underline{\Delta S(C \rightarrow A) < 0}}$$

$$AB: \text{Isobar: } \Delta S = C_V \ln \frac{T_B}{T_A} + nR \ln 2; \quad \frac{T_A}{V_0} = \frac{T_B}{2V_0} \Rightarrow \frac{T_B}{T_A} = 2$$

$$\underline{\underline{= (C_V + nR) \ln 2 = C_P \ln 2}}$$

CB: Samme temperatur i C & B (symmetri, evt ved utregning som i forrige oppgave)

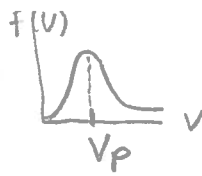
$$\underline{\underline{\Delta S = nR \ln 2}}$$

$$\underbrace{(C_V + nR) \ln 2}_{\Delta S(A \rightarrow B)} > \underbrace{nR \ln 2}_{\Delta S(C \rightarrow B)} > \Delta S(C \rightarrow A)$$

$$\Rightarrow \underline{\underline{\Delta S(A \rightarrow B) > \Delta S(C \rightarrow B) > \Delta S(C \rightarrow A)}}$$

16) C

17) Maxwell's farts fordeling



$$f(v) = 4\pi \left(\frac{m}{2\pi kT} \right)^{3/2} v^2 e^{-\frac{mv^2}{2kT}}$$

$$\frac{df}{dv} = \cancel{0} = 4\pi \left(\frac{m}{2\pi kT} \right)^{3/2} \left\{ 2ve^{-\frac{mv^2}{2kT}} + v^2 \cdot e^{-\frac{mv^2}{2kT}} \cdot -\frac{mv}{kT} \right\}$$

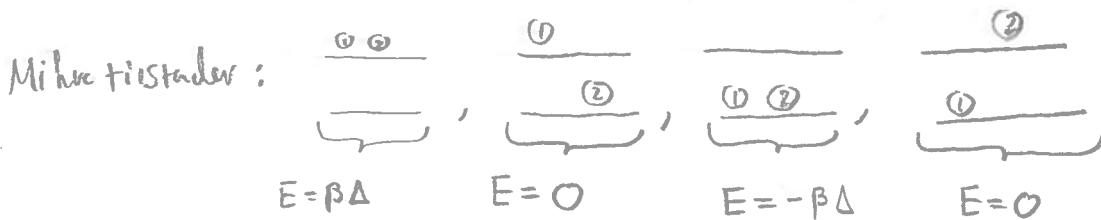
$$= 4\pi \left(\frac{m}{2\pi kT} \right)^{3/2} e^{-\frac{mv^2}{2kT}} \underbrace{\left\{ 2 - \frac{mv^2}{kT} \right\}}_{0 \text{ (v=v_p)}} v$$

$$\Rightarrow \underline{\underline{v_p = \sqrt{\frac{2kT}{m}}}}$$

18) ——— $E_2 = \frac{\beta\Delta}{2}$

to partikler ① ②

————— $E_1 = -\frac{\beta\Delta}{2}$



$$\Delta S = S(E = \beta\Delta) - S(E = 0) = \underbrace{k \ln 1}_0 - k \ln 2 = \underline{\underline{-k \ln 2}}$$

19)

$$E = J/2, \quad 0, \quad -J/2, \quad 0$$

$$Z = \sum_j e^{-\beta E_j} = 2 e^{-\beta \frac{J}{2}} + 2 e^{\beta \frac{J}{2}} + 4 e^0$$

$$= \underline{\underline{4 \left(1 + \cosh \left(\beta \frac{J}{2} \right) \right)}}$$

20)

$$F = F_i + kT B_2(T) \frac{N^2}{V}$$

$$P = - \left(\frac{\partial F}{\partial V} \right)_T \quad \text{From TDI}$$

$$\left(\frac{\partial F}{\partial V} \right)_T = \left(\frac{\partial F_i}{\partial V} \right)_T - kT B_2(T) \frac{N^2}{V^2}$$

$$P = \frac{NkT}{V} + \frac{NkT}{V} \cdot \frac{N}{V} B_2(T) = \underline{\underline{\frac{NkT}{V} \left(1 + \frac{N}{V} B_2(T) \right)}}$$

21)

$$S = - \left(\frac{\partial F}{\partial T} \right)_V$$

$$\left(\frac{\partial F}{\partial T} \right)_V = \left(\frac{\partial F_i}{\partial T} \right)_V + k \frac{N^2}{V} \left\{ B_2(T) + T \frac{dB_2}{dT} \right\}$$

$$\Rightarrow S = S_i - \frac{kN^2}{V} \left\{ B_2(T) + T \frac{dB_2}{dT} \right\}$$

$$22) B_2(T) = -\frac{a}{RT}$$

$$\frac{dB_2}{dT} = \frac{a}{RT^2} \quad ; \quad \frac{d^2B_2}{dT^2} = -\frac{2a}{RT^3}$$

$$C_v = C_{v,i} - R \frac{N^2}{V} T \left(\frac{2a}{RT^2} + -\frac{2a}{RT^2} \right) = C_{v,i} = \frac{3}{2} R$$

EPP: Kun bidrag fra translation ved alle temperaturer

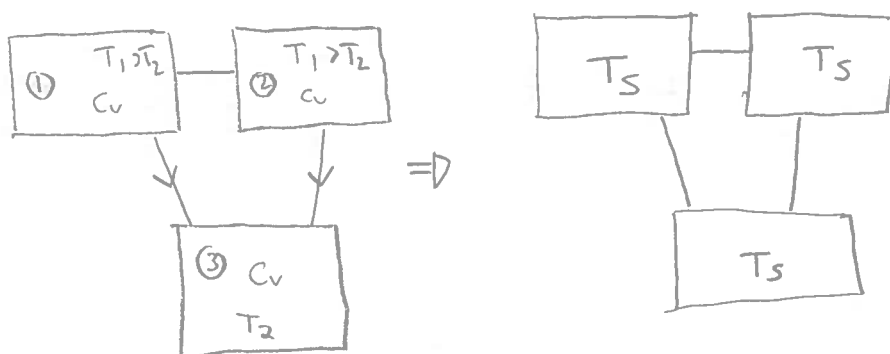
\Rightarrow gasser må være monoatomisk

23) A, B, C, D, E Alt er ekvivalent

$$24) dS = C_v \frac{dT}{T} + \frac{nR}{V} \underbrace{dV}_0 = 0 \quad \Delta S = C_v \ln \frac{T_{\text{sup}}}{T_{\text{inf}}}$$

$$\Rightarrow \Delta S = C_v \ln \frac{T_{012}}{T_0} = C_v \ln \frac{1}{2} = -C_v \ln 2 = \underline{\underline{-\frac{3}{2} Nk \ln 2}}$$

25)



$$Q = \Delta U + W \Rightarrow \Delta U = \Delta U_{(1)} + \Delta U_{(2)} + \Delta U_{(3)} = 0$$

$$\downarrow \rightarrow 0$$

$$\Rightarrow T_s = \frac{T_1 + T_1 + T_2}{3} = \frac{2T_1 + T_2}{3}$$

$$\Delta S = \Delta S_{(1)} + \Delta S_{(2)} + \Delta S_{(3)} = 2C_v \ln \frac{T_s}{T_1} + C_v \ln \frac{T_s}{T_2} = C_v \ln \frac{T_s^3}{T_1^2 T_2}$$

$$= C_v \ln \frac{(2T_1 + T_2)^3}{9T_1^2 T_2}$$

$$26) \quad dW = -f dL$$

$$P \rightarrow -f, \quad V \rightarrow L$$

$$S(L, T) = g(T) - c \left(\frac{L^2}{2L_0} + \frac{L_0^2}{L} \right)$$

$$dF = -S dT + P dV$$

$$\Rightarrow \left(\frac{\partial S}{\partial V} \right)_T = \left(\frac{\partial P}{\partial T} \right)_V$$

$$\Rightarrow \left(\frac{\partial S}{\partial L} \right)_T = - \left(\frac{\partial f}{\partial T} \right)_L$$

$$+ \left(\frac{\partial f}{\partial T} \right)_L = + c \left(\frac{L}{L_0} - \frac{L_0^2}{L^2} \right)$$

$$\Rightarrow f(T, L) = cT \left(\frac{L}{L_0} - \frac{L_0^2}{L^2} \right) + h(L)$$

$$f(T, L=L_0) = 0 = cT(1-1) + h(L) \Rightarrow h(L) = 0$$

$$\Rightarrow \underline{\underline{f(T, L) = cT \left(\frac{L}{L_0} - \frac{L_0^2}{L^2} \right)}}$$

$$27) U = aVT^4$$

$$T ds = dU + PdV = \left(\frac{\partial U}{\partial V} \right)_T dV + \left(\frac{\partial U}{\partial T} \right)_V dT + PdV$$

$$= (aT^4 + p) dV + 4aVT^3 dT$$

$$\Rightarrow ds = \left(\frac{p}{T} + aT^3 \right) dV + \underline{\underline{4aVT^2 dT}}$$

$$S = V \cdot s(T)$$

$$dS = \left(\frac{\partial S}{\partial V} \right)_T dV + \left(\frac{\partial S}{\partial T} \right)_V dT = \underline{s(T) dV} + \underline{\underline{\left(\frac{\partial S}{\partial T} \right)_V dT}}$$

$$\Rightarrow s(T) = aT^3 + \frac{p}{T} \quad \text{og} \quad \left(\frac{\partial S}{\partial T} \right)_V = 4aVT^2$$

$$\Rightarrow S'(T, V) = \frac{4}{3} aVT^3 + f(V)$$

$L > 0$ (initialbeding)

$$= \underline{\underline{\frac{4}{3} aVT^3}}$$

$$28) P = \frac{a}{3} T^4$$

$$H = U + PV = aVT^4 + \frac{a}{3} VT^4 = \frac{4}{3} aVT^4 = 4PV$$

$$C_P = \left(\frac{\partial H}{\partial T} \right)_P = 4P \left(\frac{\partial V}{\partial T} \right)_P = 4P \left(\frac{\partial T}{\partial V} \right)_P^{-1}$$

$$\left(\frac{\partial T}{\partial V} \right)_P \Rightarrow T(V, P) = \left(\frac{3P}{a} \right)^{1/4} \Rightarrow \left(\frac{\partial T}{\partial V} \right)_P = 0 \Rightarrow \left(\frac{\partial V}{\partial T} \right)_P = \underline{\underline{\infty}}$$

$$\Rightarrow \underline{\underline{C_P = \infty}}$$

$$29) dG = Vdp - SdT + \mu dN \quad ; \quad G = \mu N$$

$$= N d\mu + \mu dN$$

$$pV = NkT \Rightarrow \frac{N}{V} = \frac{p}{kT}$$

$$\Rightarrow N d\mu = V dp - S dT$$

$$\int \frac{dy}{dx} dx$$

$$\Rightarrow \frac{N}{V} = \left(\frac{dp}{d\mu} \right)_T = \frac{p}{kT}$$

$$\Rightarrow \left(\frac{\partial \mu}{\partial p} \right)_T = \frac{kT}{p}$$

$$\Rightarrow \mu = \mu_0 = kT \ln \frac{p}{p_0}$$

$$\Rightarrow \underline{\underline{\mu = \mu_0 + kT \ln \frac{p}{p_0}}}$$

30) Siden de støkiometriske koeff. alle er 1
er LV-betingelsen ~~at alle kemiske potensial~~
~~skal være lige. Dvs alternativ C~~

$$\mu(A) + \mu(B) = \mu(C) + \mu(D)$$

\Rightarrow eneste mulighed er C

31) Ideell blanding:

$$\mu_j = \mu_j^{\circ} + RT \ln X_j$$

molebrøk salt: x

\Rightarrow molebrøk vann: $1-x$

$$\begin{aligned} \text{LV: } \mu_{\text{H}_2\text{O}}^{\circ}(P, T) &= \underbrace{\mu_{\text{H}_2\text{O}}(P + \Delta P, T)} \\ &= \underbrace{\mu_{\text{H}_2\text{O}}^{\circ}(P + \Delta P, T)} + \underbrace{RT \ln(1-x)}_{-RTx} \quad (\text{ideell blanding}) \\ &= \mu_{\text{H}_2\text{O}}^{\circ}(P, T) + \Delta P \left(\frac{\partial \mu_{\text{H}_2\text{O}}^{\circ}}{\partial P} \right)_T \end{aligned}$$

$$\begin{aligned} \Rightarrow \frac{RTx}{\Delta P} &= \underbrace{\left(\frac{\partial \mu_{\text{H}_2\text{O}}^{\circ}}{\partial P} \right)_T} \\ &= \frac{V_{\text{H}_2\text{O}}}{N_{\text{H}_2\text{O}}} \approx \frac{V}{N} \end{aligned}$$

$$\Rightarrow \frac{RT}{\Delta P} \frac{N_{\text{salt}}}{N} = \frac{V}{N} \Rightarrow \underline{\underline{\Delta P = \frac{N_{\text{salt}} RT}{V}}}$$

32) A, B

C er feil fordi at kvasistatiske prosesser hvor vi ikke ser bort i fra dissipative effekter (f.eks friksjon) ikke er reversible.

$$33) T ds = C_v dT + T \left(\frac{\partial P}{\partial T} \right)_v dv$$

$$P \rightarrow -\mu_0 \mathcal{H}, \quad V \rightarrow \mathcal{M}$$

$$\mu = C \frac{\mathcal{H}}{T} \Rightarrow \mathcal{H} = \frac{\mu T}{C}$$

$$\Rightarrow T ds = \underbrace{C_{\mu}}_{\left(\frac{\partial U}{\partial T} \right)_{\mu} = 0} dT - T \mu_0 \underbrace{\left(\frac{\partial \mathcal{H}}{\partial T} \right)_{\mu}}_{\frac{\mu}{C}} d\mu = -T \frac{\mu_0}{C} \mu d\mu$$

$$\Rightarrow \int_{S_0}^S ds = -\frac{\mu_0}{C} \int_0^{\mu_1} \mu d\mu = -\frac{\mu_0 \mu_1^2}{2C}$$

$$\Rightarrow \underline{\underline{\Delta S = -\frac{\mu_0 \mu_1^2}{2C}}}$$

34) B

$$35) dT = \frac{\mu_0 T}{C \mathcal{H}} \left| \left(\frac{\partial \mu}{\partial T} \right)_{\mathcal{H}} \right| d\mathcal{H}$$

$$C_{\mathcal{H}} = \left(\frac{\partial H}{\partial T} \right)_{\mathcal{H}} = \frac{\partial}{\partial T} \left\{ -\mu_0 \mathcal{H} \cdot C \left(\frac{\mathcal{H}}{T} \right)^{\alpha} \right\}$$

$$= -\mu_0 C \mathcal{H}^{\alpha+1} \cdot -\alpha \frac{1}{T^{\alpha+1}} = \alpha \mu_0 C \left(\frac{\mathcal{H}}{T} \right)^{\alpha+1}$$

$$\left(\frac{\partial \mu}{\partial T} \right)_{\mathcal{H}} = \frac{\partial}{\partial T} \left\{ C \frac{\mathcal{H}^{\alpha}}{T^{\alpha}} \right\} = C \mathcal{H}^{\alpha} \frac{\partial}{\partial T} \left\{ T^{-\alpha} \right\} = C \mathcal{H}^{\alpha} \cdot -\alpha \frac{1}{T^{\alpha+1}}$$

$$dT = \mu_0 T \cdot \frac{T^{\alpha+1}}{\mathcal{H}^{\alpha+1} \alpha \mu_0 C} \cdot \frac{\alpha C \mathcal{H}^{\alpha}}{T^{\alpha+1}} d\mathcal{H} = \frac{T}{\mathcal{H}} d\mathcal{H}$$

$$\Rightarrow \frac{dT}{T} = \frac{d\mathcal{H}}{\mathcal{H}} \Rightarrow \ln T = \ln \mathcal{H} + C$$

$$\underline{\underline{T = A \mathcal{H}}}$$

36)

$$V = \frac{RT}{P} + \left(b - \frac{a}{RT}\right)$$

$$\mu_{JT} = \frac{1}{C_p} \left[T \left(\frac{\partial V}{\partial T} \right)_P - V \right]$$

$$= \frac{1}{C_p} \left[T \left(\frac{R}{P} + \frac{a}{RT^2} \right) - V \right] = \frac{1}{C_p} \left[\frac{RT}{P} + \frac{a}{RT} - V \right]$$

$$= \frac{1}{C_p} \left[\cancel{V} - b + \frac{a}{RT} + \frac{a}{RT} - \cancel{V} \right] = \frac{1}{C_p} \left[\frac{2a}{RT} - b \right]$$

Inversions temperature: $\mu_{JT}(T_x) = 0$

$$\Rightarrow \frac{2a}{RT_x} - b = 0 \Rightarrow \underline{\underline{T_x = \frac{2a}{Rb}}}$$

37)
$$P = \frac{NkT}{V-Nb} - a \frac{N^2}{V^2}$$

$$\frac{\partial P}{\partial V} = -\frac{NkT}{(V-Nb)^2} + 2a \frac{N^2}{V^3} \Rightarrow \frac{NkT_c}{(V-Nb)^2} = 2a \frac{N^2}{V^3} \quad (1)$$

$$\frac{\partial^2 P}{\partial V^2} = 2 \frac{NkT}{(V-Nb)^3} - 6a \frac{N^2}{V^4} \Rightarrow \frac{NkT_c}{(V-Nb)^3} = 3a \frac{N^2}{V^4} \quad (2)$$

$$\Rightarrow \frac{(1)}{(2)} \Rightarrow \frac{NkT_c}{(V_c-Nb)^2} \cdot \frac{(V_c-Nb)^3}{NkT_c} = (V_c-Nb) = \frac{2aN^2}{V_c^3} \cdot \frac{V_c^4}{3aN^2} = \frac{2}{3} V_c$$

$$\Rightarrow \frac{V_c}{3} = Nb \Rightarrow V_c = \underline{\underline{3Nb}}$$

$$\Rightarrow NkT_c = 2a \frac{N^2}{3^3 N^2 b^3} \cdot (2Nb)^2 = \frac{2a}{27Nb^3} \cdot 4N^2 b^2$$

$$= \frac{8aN}{27b} \Rightarrow \underline{\underline{T_c = \frac{8a}{27Rb}}}$$

$$38) P = \frac{RT}{V-b} - \frac{a}{V^2}$$

$$dU = \underbrace{\left(\frac{\partial U}{\partial T}\right)}_{C_V} dT + \underbrace{\left(\frac{\partial U}{\partial V}\right)}_{T\left(\frac{\partial P}{\partial T}\right)_V - P} dV$$

$$T\left(\frac{\partial P}{\partial T}\right)_V - P = \frac{RT}{V-b} - P = \frac{a}{V^2}$$

$$dU = C_V dT + a \frac{dV}{V^2}$$

$$\Rightarrow U - U_0 = C_V T - \frac{a}{V} \Rightarrow \underline{U = C_V T - \frac{a}{V} + U_0}$$

$$39) \frac{dP}{dT} = \frac{l}{T \Delta V} = \frac{L_f}{T \Delta V} \quad \Delta V = V_g - V_v \approx V_g = \left(\frac{RT}{P}\right)^2$$

$$\frac{dP}{dT} = \frac{L_f P^2}{A^2 T^3} \Rightarrow \int_{P_1}^P \frac{dP}{P^2} = \frac{L_f}{A^2} \int_{T_1}^T \frac{dT}{T^3}$$

$$= - \left[\frac{1}{P} - \frac{1}{P_1} \right] = \frac{L_f}{A^2} \left[-\frac{1}{2} T^{-2} \right]_{T_1}^T = \frac{-L_f}{2A^2} \left[\frac{1}{T^2} - \frac{1}{T_1^2} \right]$$

$$\Rightarrow \underline{\underline{\frac{1}{P} - \frac{1}{P_1} = \frac{L_f}{2A^2} \left(\frac{1}{T^2} - \frac{1}{T_1^2} \right)}}$$

$$40) P \rightarrow -\mu_0 H \quad V \rightarrow M$$

$$-\mu_0 \frac{dH}{dT} = \frac{\Delta S}{\Delta M} \Rightarrow \Delta S = -\underbrace{\mu_0}_{\text{positiv}} \Delta M \cdot \frac{dH}{dT}$$

$$A: \Delta S > 0 \quad B: \Delta S = 0 \quad C: \Delta S < 0 \quad D: \Delta S = 0 \quad E: \Delta S = +\infty$$

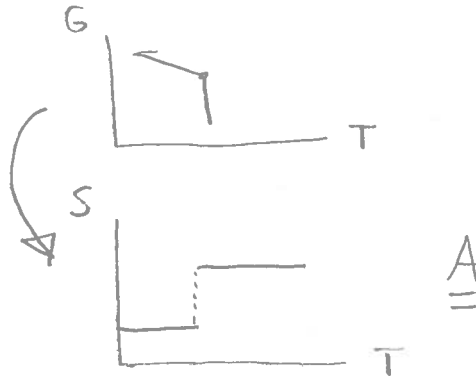
$$\underline{\underline{\Delta S_E > \Delta S_A > \Delta S_B = \Delta S_D > \Delta S_C}}$$

$$41) G = U + PV - TS$$

$$Tds = du + pdv$$

$$dG = du + p dv + v dp - T ds - s dT = v dp - s dT$$

$$\Rightarrow S = - \left(\frac{\partial G}{\partial T} \right)$$

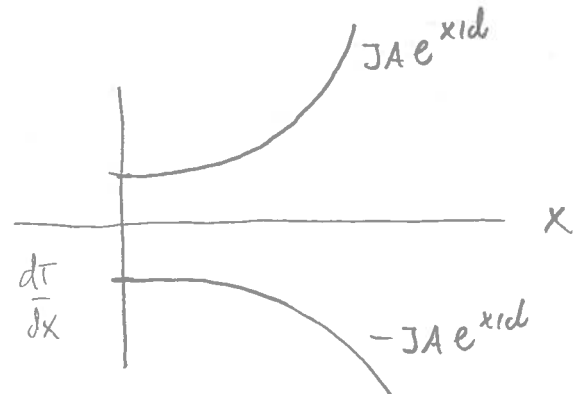


$$42) \Delta T = J \sum_{i=1}^N \frac{a_i}{h_i} = J \left(\frac{b_1}{h_1} + \frac{b_2}{h_2} \right)$$

$$J = \frac{20}{\frac{1}{0.14} + \frac{2}{0.047}} = \underline{\underline{0.40}} \frac{\text{W}}{\text{m}^2}$$

$$43) J = -h(x) \frac{dT}{dx} \quad \text{Siden ensi er bevart er } J = \text{konst. i } \underline{\underline{1D}}$$

$$\Rightarrow \frac{dT}{dx} = - \frac{J}{h(x)} = -JA e^{x/d}$$



$\frac{dT}{dx}$ blir mindre og mindre med økende x

\Rightarrow Beste mulighet er B

$$44) \nabla \cdot \mathbf{J} = 0 \quad \mathbf{J} = -\mathcal{H} \nabla T$$

$$0 = \frac{d\mathbf{J}}{dx} = \frac{d}{dx} \left(-\mathcal{H} \frac{dT}{dx} \right)$$

$$\Rightarrow 0 = \frac{d\mathcal{H}}{dx} \frac{dT}{dx} + \mathcal{H} \frac{d^2T}{dx^2}$$

$$\mathcal{H}(x) = A e^{-x/d}$$

$$\frac{d\mathcal{H}}{dx} = A e^{-x/d} \cdot -\frac{1}{d} = -\frac{A}{d} e^{-x/d}$$

$$\Rightarrow 0 = -\frac{A}{d} e^{-x/d} \frac{dT}{dx} + A e^{-x/d} \frac{d^2T}{dx^2}$$

$$\Rightarrow \boxed{0 = \frac{d^2T}{dx^2} - \frac{1}{d} \frac{dT}{dx}}$$

45) Symmetrisk fordeling $\Rightarrow \langle \vec{r} \rangle = \langle \vec{r} \rangle^2 = 0$

$$\langle \vec{r}^2 \rangle = \langle x^2 + y^2 \rangle = \iint_{-\infty}^{\infty} (x^2 + y^2) \tilde{p}(x, y, t) dx dy$$

$$= \int_0^{\infty} \int_0^{2\pi} r^2 \cdot \frac{1}{4\pi Dt} e^{-\frac{r^2}{4Dt}} r dr d\theta$$

$$= \frac{2\pi}{4\pi Dt} \int_0^{\infty} r^3 e^{-\frac{r^2}{4Dt}} dr$$

$\underbrace{\hspace{10em}}_{\frac{1}{2\left(\frac{1}{4Dt}\right)^2}}$

$$= \frac{1}{2Dt} \frac{16D^2t^2}{2} = 4Dt$$

$$\Rightarrow \text{Var}(\vec{r}) = \langle \vec{r}^2 \rangle - \langle \vec{r} \rangle^2 = \underline{\underline{4Dt}}$$