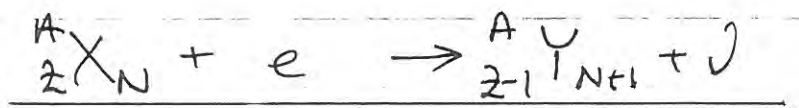


Problem 1.

a) EC disintegration:



$$Q = (\sum m_{\text{initial}} - \sum m_{\text{final}}) c^2$$

$$Q = c^2 \left\{ \underbrace{m({}^A_Z X)}_{\text{atomic mass}} - \underbrace{Z m_e}_{\approx \text{nuclear mass}} + \left(m_e - \frac{E_B}{c^2} \right) - \left[\underbrace{m({}^A_{Z-1} Y)}_{\text{binding energy of electron}} + \underbrace{(Z-1)m_e}_{m_\nu \approx 0} \right] \right\}$$

$$Q = \left[m({}^A_Z X) - m({}^A_{Z-1} Y) \right] + \left(-Z + 1 + Z - 1 \right) m_e - \frac{E_B}{c^2} \Big] c^2$$

$$\underline{\underline{Q = \left[m({}^A_Z X) - m({}^A_{Z-1} Y) \right] c^2 - E_B}}$$

If transition goes to an excited state Y^* of the daughter nucleus:

$$\Rightarrow \underline{\underline{Q^* = \left[m({}^A_Z X) - m({}^A_{Z-1} Y^*) \right] c^2 - E_B}}$$

$$\underline{\underline{\text{where } m({}^A_{Z-1} Y^*) = m({}^A_{Z-1} Y) + \frac{E_{ex}}{c^2}, \quad E_{ex} = E_2, E_3}}$$

(2)

Transition to $E_3 = 2340 \text{ keV}$:

$$Q_3 = (m_p - m_d)c^2 - E_3 - E_B$$

parent
daughter

$$(m_p - m_d) = \left[\left(m\left(\begin{smallmatrix} A \\ Z \\ X \end{smallmatrix}\right) - A \right) - \left(m\left(\begin{smallmatrix} A \\ Z-1 \\ Y \end{smallmatrix}\right) - A \right) \right] = \Delta_P - \Delta_D$$

$$\Rightarrow (m_p - m_d) = (\Delta_{\text{Bi-207}} - \Delta_{\text{Po-207}}) = [-21545 - (-24119)] \mu\text{u}$$

$$\underline{(m_p - m_d) = \Delta m = 2574 \mu\text{u}}$$

Capture of K-electrons: $E_K(\text{Bi-207}) = E_K(Z=83) = \underline{90.5 \text{ keV}}$.

$$\underline{Q_{3K}} = 2574 \cdot 10^{-6} \text{ u} \cdot 931,502 \frac{\text{keV}}{\text{u}} - 2340 \text{ keV} - 90,5 \text{ keV} = \underline{-33 \text{ keV}}$$

EC transition to E_3 by K-electron capture is not possible

Capture of L-electrons: $E_L(\text{Bi-207}) = E_L(Z=83) = \underline{15.2 \text{ keV}}$.

$$\underline{Q_{3L}} = \underbrace{2574 \cdot 10^{-6} \text{ u} \cdot 931,502 \frac{\text{keV}}{\text{u}}}_{2397.7 \text{ keV}} - 2340 \text{ keV} - 15.2 \text{ keV} = \underline{43 \text{ keV}}$$

EC transition to E_3 by L-electron capture is possible since $Q_{3L} > 0$.

(3)

Transition to level $E_2 = 1630$ keV :

$$Q_2 = \Delta mc^2 - E_2 - E_B \quad E_B = \begin{cases} E_K = 90.5 \text{ keV} \\ E_L = 15.2 \text{ keV} \end{cases}$$

$$Q_{2K} = (2397.7 - 1630 - 90.5) \text{ keV} = \underline{677 \text{ keV}} > 0$$

$$Q_{2L} = (2397.7 - 1630 - 15.2) \text{ keV} = \underline{750 \text{ keV}} > 0$$

EC transition to E_2 is possible by both
K and L electron capture.

b) β^- transitions.

$$\beta_1 : \underline{\frac{5}{2}^- \rightarrow \frac{1}{2}^-}$$

Parity : $\Delta\pi = No \Rightarrow$ Possible multipoles : $M1, E2, M3, \dots$

$$|I_f - I_i| \leq \sigma_h \leq |I_f + I_i|$$

$$|\frac{1}{2} - \frac{5}{2}| \leq \sigma_h \leq |\frac{1}{2} + \frac{5}{2}|$$

$$2 \leq \sigma_h \leq 3$$

Lowest possible multipole will dominate, i.e. $E2$

Probability of $M(L+1) \ll$ Probability of $E(L)$,

therefore no contribution from $M3$.

(4)

$$J_2: \quad \underline{\frac{13}{2}^+ \rightarrow \frac{5}{2}^-}$$

Parity: $\Delta\pi = \text{Yes} \Rightarrow$ Possible multipoles: $E1, M2, E3, M4, \dots$

$$\left| \frac{5}{2} - \frac{13}{2} \right| \leq \sigma_L \leq \left| \frac{5}{2} + \frac{13}{2} \right|$$

$$4 \leq \sigma_L \leq 9$$

Lowest multipole dominates i.e. $M4$, possibly

with same contribution from $E5$

since $E(L+1)$ is not too improbable compared to $M(L)$

$$J_3: \quad \underline{\frac{7}{2}^- \rightarrow \frac{5}{2}^-}$$

Parity: $\Delta\pi = \text{No} \Rightarrow$ Possible multipoles: $M1, E2, M3, \dots$

$$\left| \frac{5}{2} - \frac{7}{2} \right| \leq \sigma_L \leq \left| \frac{5}{2} + \frac{7}{2} \right|$$

$$1 \leq \sigma_L \leq 6$$

Lowest multipole dominates i.e. $M1$, possibly

with same contribution from $E2$.

J_2 is of significant higher order than the two other transitions. Thus, E_2 has longer half-life than E_1 and $E_3 \Rightarrow E_2$ is the metastable state

(5)

c) Conversion electrons in parallel to β .

$$E_{\beta i} = \Delta E = E_i - E_j$$

here, the recoil energy of the daughter nucleus is neglected

Conversion electron by deexcitation from E_1 :

$$\underline{E_{1K}} = (E_1 - E_0) - E_K(z=82) = (570 - 0 - 88.0) \text{ keV} = \underline{482 \text{ keV}}$$

binding energy in Pb-207

$$\underline{E_{1L}} = (E_1 - E_0) - E_L(z=82) = (570 - 0 - 14.7) \text{ keV} = \underline{555 \text{ keV}}$$

CE from E_2 :

$$\underline{E_{2K}} = E_2 - E_1 - E_K(z=82) = (1630 - 570 - 88.0) \text{ keV} = \underline{972 \text{ keV}}$$

$$\underline{E_{2L}} = E_2 - E_1 - E_L(z=82) = (1630 - 570 - 14.7) \text{ keV} = \underline{1045 \text{ keV}}$$

CE from E_3 :

$$\underline{E_{3K}} = E_3 - E_1 - E_K(z=82) = (2340.0 - 570 - 88.0) \text{ keV} = \underline{1682 \text{ keV}}$$

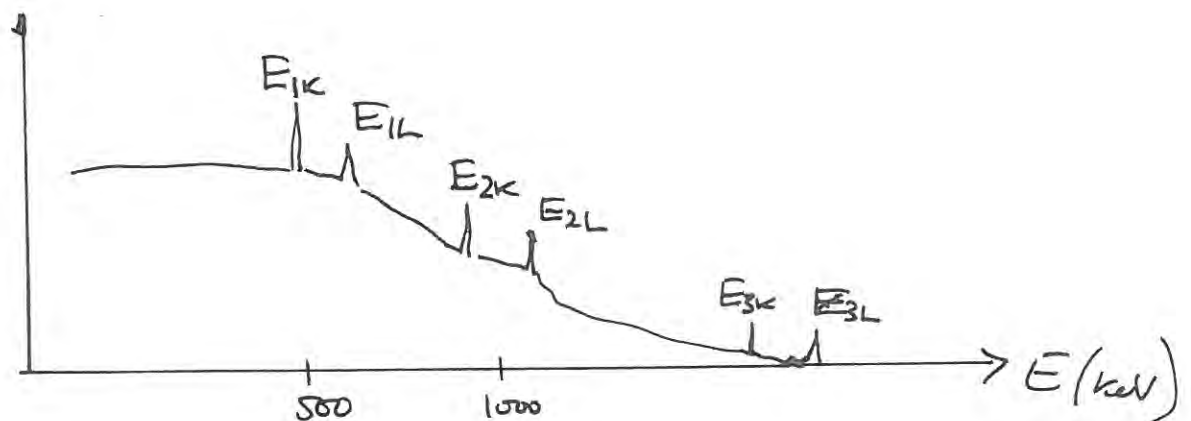
$$\underline{E_{3L}} = E_3 - E_1 - E_L(z=82) = (2340.0 - 570 - 14.7) \text{ keV} = \underline{1755.6 \text{ keV}}$$

(6)

Principle of determining stopping power of electrons (500keV, 1000keV) in e.g. aluminium.

Convulsion electrons from E_1 and E_2 have suitable energies.

Spectrum of emitted electrons



Determine shift (to the left) in peak positions between spectrum 1 acquired without material between source and detector, and spectrum 2 acquired after material of thickness Δx is inserted between source and detector. Otherwise same detection geometry.

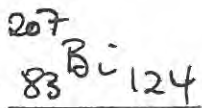
Energy decrement after passage through material thickness Δx :

$-\Delta E = E(\text{spectrum 2}) - E(\text{spectrum 1})$ recorded for peak E_{1K} and peak E_{2K} .

Estimate of stopping power $S_{col} = \left(\frac{-\Delta E}{\Delta x} \right)$ for peak $E_{1K} \approx 500 \text{ keV}$ and peak $E_{2K} \approx 1000 \text{ keV}$.

Suitable detector for β^- spectroscopy will be a semiconductor diode detector of sufficient thickness to obtain full energy deposition from the most energetic electrons in the emission.

d) Nucleon configuration in ground state:

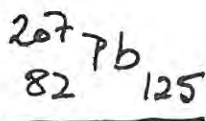


83P: Filling all shells $1s^{1/2}$... through $1h^{11/2}$, (total of 82)
single proton (no. 83) in $1h^{9/2}$ i.e. $(1h^{9/2})^1$

124N: Filling all shells $1s^{1/2}$... through $2f^{5/2}$ (total of 124)

Pairwise spin cancellation for all nucleons/states, except for single proton in state $(1h^{9/2})$, which therefore will determine net spin and parity of the ground state.

Spin: $9/2$, parity $(-1)^l = (-1)^5 = -1$ i.e. $9/2^-$



82P: Filled shells $1s^{1/2}$ through $1h^{11/2}$ (total of 82)

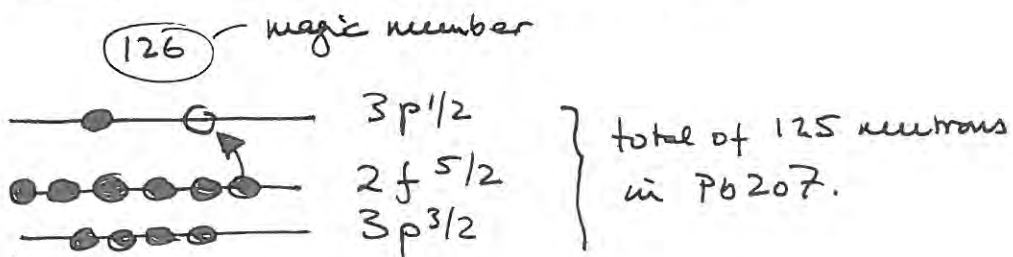
125N: Filled shells $1s^{1/2}$ through $2f^{5/2}$ (total of 124),
 and last neutron (no. 125) single in state $(3p^{1/2})^1$.

(8)

Pairwise spin cancellation except for single neutron (valence neutron) in state $(3p^{1/2})$.

Spin: $1/2$, parity $(-1)^l = (-1)^1 = -1$ i.e. $1/2^-$

Excited state E_1 of Pb_{207} is $5/2^-$



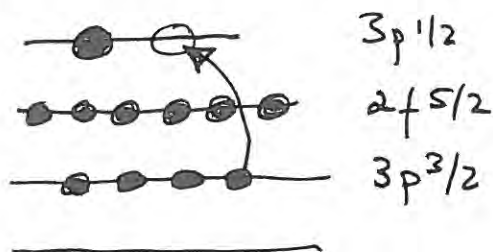
Excite one neutron from $2f^{5/2}$ to $3p^{1/2}$

⇒ Spin cancellation in $3p^{1/2}$,

single vacancy in $2f^{5/2}$ gives spin $5/2$,

parity $(-1)^l$ $l=3$ for $f \Rightarrow \pi = -1$ i.e. $5/2^-$ q.e.d.

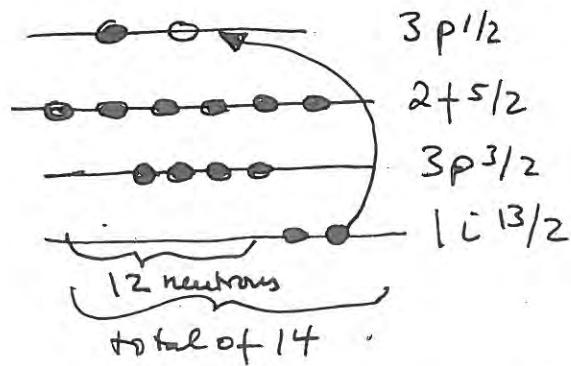
Excited state $3/2^-$ of Pb_{207} :



Excite one neutron from $3p^{3/2}$ to $3p^{1/2}$.

Single vacancy in $3p^{3/2}$ determines spin and parity i.e. $3/2^-$ q.e.d.

Excited state E_2 of Pb_{207} : $\frac{13}{2}^+$



Excite one neutron from $1i_{13/2}$ to $3p_{1/2}$

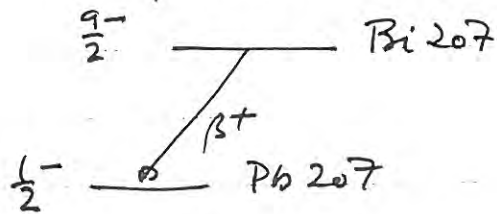
\Rightarrow Spin cancellation in $3p_{1/2}$, single vacancy in $1i_{13/2}$
 contributes spin $\frac{13}{2}$, parity $(-1)^l$ for $l=6 \Rightarrow +1$
 i.e. Excited state E_2 is $\frac{13}{2}^+$ q.e.d.

Excited state E_3 is $\frac{7}{2}^-$

Excite one neutron from $2f_{7/2}$ to $3p_{1/2}$

\Rightarrow Spin cancellation in $3p_{1/2}$,
 single vacancy in $2f_{7/2}$, \Rightarrow spin $\frac{7}{2}$,
 parity $(-1)^l$ for $l=3 \Rightarrow \pi = -1$ i.e. E_3 state is $\frac{7}{2}^-$ q.e.d.

e) Possible β^+ emission from ground state of $\text{Bi} 207$ to ground state of $\text{Pb} 207$



Classification of β^+ disintegration:

No parity change \Rightarrow Possible L -values $L = 0, 2, 4, \dots$

$(\vec{L} + \vec{S})$ is total angular momentum represented by the (β^+, ν) particle system.

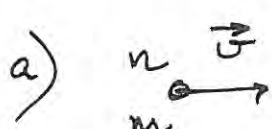
$$|I_f - I_i| \leq |\vec{L} + \vec{S}| \leq |I_f + I_i|$$

$$\left| \frac{1}{2} - \frac{9}{2} \right| \leq |\vec{L} + \vec{S}| \leq \frac{9}{2} + \frac{1}{2}$$

$$\underline{4 \leq |\vec{L} + \vec{S}| \leq 5}, \quad \underline{L \text{ even.}}$$

Lowest possible value of L is $L=4$, which in combination with both $S=0$ (Fermi) and $S=1$ (Gamow/Teller) satisfy the above requirement.

i.e. 4th forbidden transition of mixed Fermi and G/T type.

Problem 2.

M at rest in lab system.
 $\Rightarrow T_{\text{lab}} = T_n = \frac{1}{2} m v^2$

Velocity of CM (coordinate) system:

$$\underline{\vec{v}_{\text{CM}}} = \frac{\sum \vec{p}_i}{\sum m_i} = \frac{m\vec{v}}{(m+M)} = \underline{\frac{m}{m+M} \cdot \vec{v}}$$

Velocities in CM coordinate system:

$$\underline{\vec{v}'_m} = \vec{v} - \vec{v}_{\text{CM}} = \frac{m+M-m}{m+M} \cdot \vec{v} = \underline{\frac{M}{m+M} \cdot \vec{v}}$$

$$\underline{\vec{v}'_M} = 0 - \vec{v}_{\text{CM}} = \underline{-\frac{m}{m+M} \cdot \vec{v}}$$

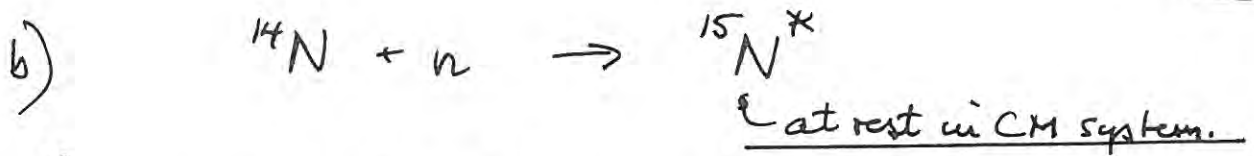
Total kinetic energy in CM system (i.e. "collision energy"):

$$T_c = \sum T'_i = \frac{1}{2} m (v'_m)^2 + \frac{1}{2} M (v'_M)^2 = T'$$

$$T_c = \frac{1}{2} m \left(\frac{M}{m+M} \right)^2 v^2 + \frac{1}{2} M \left(\frac{m}{m+M} \right)^2 v^2 = \frac{1}{2} m v^2 \cdot \frac{M}{m+M}$$

$$\underline{\underline{T_c = T_n \cdot \frac{M}{m+M} \quad \text{q.e.d.}}}$$

(It is easily seen that kinetic energy of the center of mass, $T_{\text{CM}} = \frac{1}{2} (m+M) v_{\text{CM}}^2 = T_n \frac{m}{m+M}$, thus total energy $T = T' + T_{\text{CM}} = \frac{1}{2} m v^2 = T_{\text{lab}}$.)



Conservation of total energy in CM coordinate system:

$$m(^{14}\text{N}) + m_n + T_c/c^2 = M(^{15}\text{N}^*) = M(^{15}\text{N}) + E^*/c^2$$

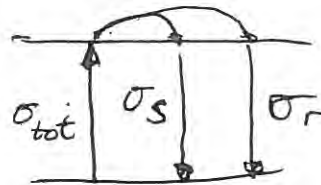
$$\Rightarrow E^* = [m(^{14}\text{N}) + m_n - m(^{15}\text{N})]c^2 + T_c$$

Excitation energy of $^{15}\text{N}^*$: $(T_c = 1.779 \cdot \frac{14}{14+1} \text{ MeV})$

$$E^* = \left[(14.003074 + 1.008665 - 15.000109) \cdot 931.5 + \frac{14}{14+1} \cdot 1.779 \right] \text{ MeV}$$

$$\underline{E^* = (10.8333 + 1.6604) \text{ MeV} = 12.49 \text{ MeV.}}$$

c) Branching for reactions, i.e. interactions different from elastic scattering (s)

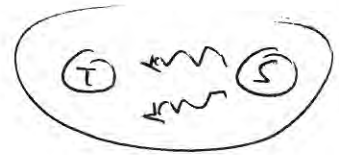


$$y_r = \frac{\sigma_r}{\sigma_{tot}} \quad \sigma_{tot} = \sigma_s + \sigma_r$$

$$\sigma_r = \sigma_{tot} - \sigma_s = 406 - 246 = 166.$$

$$\underline{y_r = \frac{\sigma_r}{\sigma_{tot}} = \frac{166}{406} = 0.4 \text{ i.e. } 40\%}$$

40% of interactions will be something else than elastic scattering.

Problem 3.

a) Specific effective energy:

$$SEE = \frac{1}{m_T} \sum_i k_i \bar{E}_i \varphi_i (T \leftarrow S)$$

m_T : mass of target organ T. (here T = whole body)

k_i : yield of radiation type i from source S.

\bar{E}_i : average quantum energy of radiation i .
(= E_γ for gamma, $\approx \frac{1}{3} E_{max}$ for beta)

φ_i : absorbed fraction of radiation i in target T.

For α and β : $\varphi = 0$ if $S \neq T$
 $\varphi = 1$ if $S \equiv T$

For γ : φ must be known (here 0.75 for 662 keV)

	k_i	\bar{E}_i	φ_i	$k_i \bar{E}_i \varphi_i$
β_1	94.4%	514/3 keV	1	161.7 keV
γ	94.4%	662 keV	0.75	468.7 keV
β_2	5.6%	1176/3 keV	1	21.9 keV
				<u>$\Sigma = 652.4 \text{ keV/dis.}$</u>

$$SEE = \frac{1}{70 \text{ kg}} \cdot 652.4 \text{ keV/dis} \cdot 1.602 \cdot 10^{-16} \frac{\text{J}}{\text{keV}}$$

$$SEE = 14.9 \cdot 10^{-16} \frac{\text{J}}{\text{kg} \cdot \text{dis.}} = 1.5 \cdot 10^{-15} \text{ Gy/dis. q.e.d.}$$

b) Committed effective dose after intake of radioactivity A_0 at time $t=0$:

$$E(50) = \sum_T \omega_T H_T(50) = \sum_T \omega_T \sum_R \omega_R \cdot D_{TR}(50)$$

tissue weighting factor
radiation weighting factor.

$$D_{TR}(50) = \int_0^{50 \text{ yrs}} \dot{D}_{TR}(t) dt = \int_0^{50} SEE \cdot A_0 e^{-\lambda_{\text{eff}} \cdot t} dt.$$

Here, $\omega_T = 1,0$ (whole body as source and target organ),
 $\omega_R = 1$ (β and γ radiation)
 $\lambda_{\text{eff}} = \lambda_{\text{rad}} + \lambda_{\text{biol}}$, $\lambda = \ln 2 / T_{1/2}$.

$T_{1/2 \text{ biol}} \ll T_{1/2 \text{ rad}}$ and $T_{1/2 \text{ biol}} \ll 50 \text{ yrs}$

$$\Rightarrow \lambda_{\text{eff}} \approx \lambda_{\text{biol}} \text{ and } \int_0^{50 \text{ yrs}} e^{-\lambda_{\text{eff}} t} dt \approx \int_0^{\infty} e^{-\lambda_{\text{eff}} t} dt = \frac{1}{\lambda_{\text{eff}}}.$$

With the above simplifications:

$$e(50) = \frac{E(50)}{A_0} = \frac{1}{A_0} \cdot \omega_T \omega_R \cdot SEE \cdot A_0 \cdot \frac{1}{\lambda_{\text{biol}}},$$

$$\Rightarrow T_{1/2 \text{ biol}} = \frac{\ln 2}{\lambda_{\text{biol}}} = \frac{\ln 2 \cdot e(50)}{\omega_T \omega_R \cdot SEE}$$

$$T_{1/2 \text{ biol}} = \frac{0,693 \cdot 13 \cdot 10^{-6} \text{ Sv} \cdot \frac{1}{10^3 \text{ Bq}}}{1,0 \cdot 1 \text{ Sv/gy} \cdot 1,5 \cdot 10^{-15} \frac{\text{Sv}}{\text{Bq}} \cdot \frac{1}{89,5}} = 6 \cdot 10^6 \text{ s}$$

$$\underline{T_{1/2 \text{ biol}}} = \frac{6 \cdot 10^6 \text{ s}}{24 \cdot 3600 \text{ s/day}} = \underline{69 \text{ days}} \quad (\ll 30 \text{ yrs})$$