



**Problem 1. Dielectrics and Ferroelectrics****1a)**

Assume that we are studying a dielectric sample applying an external electric field  $E$  to it, such that the responding polarization  $P = \epsilon_0 \chi E$ , where  $\chi = \epsilon - 1$  is the dielectric susceptibility,  $\epsilon$  is the dielectric constant and  $\epsilon_0$  is the vacuum permittivity.

- Discuss the distinction between orientational (or)  $P_{or}$  and distortional (d)  $P_d$  polarizations respectively, in dielectrics.

In a dielectric, one needs to distinguish between the externally applied field,  $E$ , and the local field,  $E_{loc}$ , in the sample. It is the local field at the site of a molecule that acts to polarize or reorient that molecule. The relation between the polarization  $P$  and the local field for monomolecular dielectrics is

$$P = P_{or} + P_d = \rho \beta E_{loc} \quad (1.1)$$

where  $\rho$  is the particle density, and  $\beta = \beta_{or} + \beta_d$  is the polarisability.

- Discuss without entering into details, how it can be shown that for isotropic liquids and symmetric crystals

$$E_{loc} = E + P/(3\epsilon_0) \quad (1.2)$$

where  $E$  is the externally applied electric field, and  $P/(3\epsilon_0) = E_L$  is the so-called Lorentz-field, which represents the average (mean-field) contribution to  $E_{loc}$  acting on one molecule from all the other molecules in the sample.

- Show that for this case:

$$(\epsilon - 1)/(\epsilon + 2) = \rho \beta / (3\epsilon_0) \quad (1.3)$$

which is the Clausius-Mosotti equation for dielectrics.

**1b)**

$P_{or}$  is related to the orientation of molecular dipoles in the dielectric. The potential energy,  $u$ , of a molecule with dipole moment  $p_0$  in an electric field  $E_{loc}$  is:

$$u = -p_0 E_{loc} \cos \theta \quad (1.4)$$

where  $\theta$  is the angle between the dipole - direction and the  $E_{loc}$  - direction.

Defining the orientational distribution function for dipolar molecules in the dielectric as  $w(\varphi, \theta)$ , and thus obtaining the fraction of molecular dipoles in the angular interval  $d\varphi d\theta$  as  $w(\varphi, \theta) \sin(\theta) d\varphi d\theta$  (for  $0 < \theta < \pi$  and  $0 < \varphi < 2\pi$ ), we get by using standard Boltzmann statistics i.e.  $w(\theta, \varphi) \propto \exp(-u/k_B T)$ , where  $k_B$  is the Boltzmann constant and  $T$  is the temperature, that:

$$P_{or} = \rho p_0 \langle \cos\theta \rangle = \rho p_0 \frac{\int d\varphi \int d\theta w(\varphi, \theta) \sin(\theta) \cos\theta}{\int d\varphi \int d\theta w(\varphi, \theta) \sin(\theta)} \\ = \rho p_0 \frac{\int d\theta [2\pi \sin(\theta) \cos(\theta) \exp(x \cos(\theta))] / Z}{\int d\theta [2\pi \sin(\theta) \exp(x \cos(\theta))] / Z} \quad (1.5)$$

where the variable  $x = p_0 E_{loc} / (k_B T)$ , the integral over  $\theta$  run from 0 to  $\pi$ , the one over  $\varphi$  from 0 to  $2\pi$ , and the normalization  $Z$  is the partition function:

$$Z = \int d\varphi \int d\theta w(\varphi, \theta) \sin(\theta) = \int d\theta [2\pi \sin(\theta) \exp(x \cos(\theta))] \\ = 2\pi [\exp(x) - \exp(-x)] / x \quad (1.6)$$

• Show that for temperatures  $T \gg p_0 E_{loc} / k_B$ , the orientational part of the polarisability in a dielectric can be written as:

$$\beta_{or} \approx (p_0)^2 / (3k_B T) \quad (1.7)$$

• How can this result be used in a practical experiment to measure the distortional polarisability  $\beta_d$ ?

**1c)**

• Using Equation (1.3), show that

$$\epsilon_0 \chi = \rho \beta / (1 - \rho \beta / (3\epsilon_0)) \quad (1.8)$$

• Discuss what happens, and what it means in terms of physics, if  $\beta$  increases with decreasing  $T$  in such a way that the Lorentz-field contribution  $\rho \beta / (3\epsilon_0) = 1$  at some temperature  $T = T_C$ .

The simplest  $T$ -dependence one can think of for this case, is linear in  $T$ , i.e.  $\rho \beta / (3\epsilon_0) = 1 - c(T - T_C)$ , where  $c$  is a constant.

• How does  $\chi$  depend on temperature for this case, for temperatures  $T > T_C$ ?

**1d)**

Landau developed a phenomenological thermodynamic theory applicable to ferroelectrics (as well as other systems displaying phase transitions) describing the observations that the susceptibility diverges on approaching  $T_C$  both from above and below in  $T$ .

Landau theory assumes that near the phase transition temperature  $T_C$ , it is possible to expand the Helmholtz free energy in powers of the order parameter  $\eta$  (in the case of ferroelectrics  $\eta = P$ ), thus

$$f(\eta, T) = f_0 + c_2 \eta^2 + c_3 \eta^3 + c_4 \eta^4 + c_5 \eta^5 + c_6 \eta^6 + \dots + c_n \eta^n + \dots \quad (1.9)$$

We assume that  $c_2(T) = b(T - T_C)$ , where  $b$  is a positive constant, and also assume that  $c_3, c_4, c_5, c_6, \dots, c_n, \dots$  are constants independent of  $T$ .

- Argue that for ferroelectrics (second order phase transitions by observation):  $c_3 = 0$ ,  $c_4$  is positive, and one need not consider any  $c_n$  for  $n > 4$ , i.e.  $c_n = 0$  for  $n > 4$ .
- Sketch and discuss  $f(\eta, T)$  for various  $T$ .
- Derive the temperature dependence of the order parameter for this case.

Assume that there exists an external field,  $X$  (in the case of ferroelectrics  $X = E$ ), which couples linearly to the order parameter,  $\eta$ , i.e. an extra term  $-\eta X$  must be added in the free energy expression (i.e. we use Gibbs free energy rather than Helmholtz free energy).

- Use Landau theory to derive an expression for the temperature dependence of the susceptibility both above and below  $T_C$ , thus also verifying the result obtained in **1c**.

## Problem 2. Diffusion

### 2a)

Diffusive motion of individual particles can formally be described by the time-dependent auto-correlation function  $g_1(\mathbf{r}, t)$ , defined such that the quantity  $g_1(\mathbf{r}, t)d^3\mathbf{r}$  gives the probability that a particle moves during a time  $t$  from its starting point into a volume element  $d^3\mathbf{r}$  which is  $\mathbf{r}$  away from the starting point.  $g_1(\mathbf{r}, t)$  is a probability distribution, and as such is normalized, i.e.

$$\int g_1(\mathbf{r}, t) d^3\mathbf{r} = 1 \quad (2.1)$$

We follow Einstein, and choose to break up  $g_1(\mathbf{r}, t'+t)$  into two steps, and we write

$$g_1(\mathbf{r}, t'+t) = \int g_1(\mathbf{r}-\mathbf{r}', t') g_1(\mathbf{r}', t) d^3\mathbf{r}' \quad (2.2)$$

where the first step, during time  $t'$ , achieves a displacement  $\mathbf{r}-\mathbf{r}'$  of the particle, while the second step, in the remaining time  $t$ , brings the particle to  $\mathbf{r}$ .

- What is the basic assumption underlying that we write the total probability as a product of two individual probabilities representing the two steps?
- Show that Equation (2.2) can be rewritten (assuming  $t \ll t'$  and  $|\mathbf{r}'| \ll |\mathbf{r}|$ ) as a diffusion equation:

$$\partial g_1(\mathbf{r}, t') / \partial t = D_S \nabla^2 g_1(\mathbf{r}, t') \quad (2.3)$$

where the self-diffusion coefficient

$$D_S = (\int g_1(\mathbf{r}', t) |\mathbf{r}'|^2 d^3\mathbf{r}') / (6t) = \langle |\mathbf{r}'|^2 \rangle / (6t) \quad (2.4)$$

represents the mean-squared displacement per unit time (divided by a factor 6).

**2b)**

Consider a suspension of non-interacting colloidal particles, and assume the following:

i) A particle density gradient  $\nabla\rho$  resulting in a diffusive particle current described by Fick's law

$$\mathbf{j}_D = -D_S \nabla \rho \quad (2.5)$$

ii) An external force (for example gravity)  $\mathbf{f} = -\nabla u_{\text{pot}}$ , where  $u_{\text{pot}}$  is the potential energy difference set by the external force, driving particle current

$$\mathbf{j}_f = \rho v \mathbf{f} = (\rho/\zeta) \mathbf{f} \quad (2.6)$$

where  $v$  is the mobility, and  $\zeta$  is the friction coefficient.

iii) At equilibrium, the particle density  $\rho = \rho_{\text{eq}} = \exp(-u_{\text{pot}}/k_B T)$  according to Boltzmann statistics, where  $k_B$  is Boltzmann's constant, and  $T$  is the temperature.

Using these three assumptions:

- Show how one derives the Einstein relation:

$$D_S = k_B T v = k_B T / \zeta \quad (2.7)$$

Considering the result obtained in Equation (2.4):

- Discuss the physics contained in the Einstein relation for the "self-diffusion coefficient".

Finally:

- Discuss the difference between the "self-diffusion coefficient",  $D_S$ , and the "diffusion coefficient",  $D$ .

**2c)**

Equation (2.1), including the boundary condition  $\mathbf{r}(t=0) = 0$ , i.e.  $g_1(\mathbf{r}, t=0) = \delta(\mathbf{r})$ , has the Greens-function solution

$$g_1(\mathbf{r}(t), t) = (1/(4\pi D_S t))^{3/2} \exp(-|\mathbf{r}(t)|^2/(4D_S t)) \quad (2.8)$$

For non-interacting particles the time dependent structure function  $S(\mathbf{q}, t)$  equals the Fourier transform of  $g_1(\mathbf{r}(t), t)$ .

- Discuss (briefly in terms of words and equations) the theoretical basis for a light-scattering experiment which may be used to measure  $D_S$  for non-interacting colloidal particles in solution. The scattering vector  $\mathbf{q}$  is the difference between the outgoing and incoming wave-vectors respectively.

### Problem 3. Polymers

#### 3a)

We define the end-to-end vector of a coiled polymer chain to be  $\mathbf{R}$ .

Assume that the chain can be broken up into  $N_S$  individual uncorrelated segments  $\mathbf{a}_j$  ( $j = 1, 2, 3, \dots, N_S$ ), i.e.

$$\mathbf{R} = \sum_j \mathbf{a}_j \quad (3.1)$$

For uncorrelated unrestricted  $\mathbf{a}_j$  vectors, the chain geometry is equivalent to stepwise independent Brownian motion of a diffusing particle, and we can directly adopt the results obtained in Equations (2.4) and (2.8) for the Gaussian probability distribution  $w(\mathbf{R})$  for the chain end-to-end distance  $\mathbf{R}$ , and its mean-squared end-to-end distance

$$R_0^2 = \langle |\mathbf{R}|^2 \rangle = \int w(\mathbf{R}) R^2 4\pi R^2 dR \quad (3.2)$$

- Show that  $R_0^2 \propto N \propto M$ , where  $N$  is the degree of polymerization, and  $M$  is the chain molecular weight.
- Show that a coiled Gaussian chain is an object with fractal dimension  $D_f = 2$ .

#### 3b)

The "tube model" for polymer dynamics introduced by Sam Edwards, takes entanglement effects among different chains into account in order to describe the chain dynamics in a polymer sample. This is a mean-field type model that focuses on individual chains, that includes interactions with other chains by representing them collectively as a "tube" inside which an individual chain is restricted to move, like a reptile (the de Gennes "reptation model") in linear diffusive motion (Equation (2.7)).

- Show that the disentanglement time,  $\tau_d$ , for entangled polymer chains scales as

$$\tau_d \propto M^3 \quad (3.3)$$

#### 3c)

Internal relaxation times such as  $\tau_d$ , may be measured by means of dielectric relaxation experiments. This can be described by a general set of linear equations valid also for dielectric measurements.

Consider a physical system that is forced out of equilibrium by applying a time ( $t$ ) dependent external force,  $\sigma(t)$ , to it. For linear response, the response field perturbation of the system is:

$$\gamma(t) = \alpha \sigma(t) \quad (3.4)$$

where  $\alpha$  is a linear response function which is characteristic of the specific system under consideration.

- What do  $\gamma$ ,  $\alpha$  and  $\sigma$  represent for the dielectric case?

Assuming that after “turning off  $\sigma$ ”, the system returns to equilibrium,  $\gamma(t = \infty)$ , at a rate,  $d\gamma(t)/dt$ , which increases proportionally with the magnitude of the perturbation,  $\gamma(t)$ , we may write:

$$d\gamma(t)/dt = (\gamma(t = \infty) - \gamma(t))/\tau_d \quad (3.5)$$

where the system characteristic time,  $\tau_d$ , is the constant of proportionality between the rate and the perturbation. Equation (3.5) represents a simple relaxation process (Debye relaxation).

Assume, that we apply a time dependent periodic force to the system, such that:

$$\sigma(t) = \sigma_0 \exp(-i\omega t) \quad (3.6)$$

where  $\omega$  is the applied frequency, and  $\sigma_0$  is the force amplitude. ( $i = \sqrt{-1}$ )  
The resulting linear response can be written:

$$\gamma(t) = \alpha(\omega) \sigma(t) = \gamma_0 \exp(-i\omega t + i\delta(\omega)) \quad (3.7)$$

where  $\delta(\omega)$  is the phase difference between the force and the response, and  $\gamma_0$  is the response amplitude.

- Use the above Equations in order to derive an expression for the complex frequency dependent linear response function

$$\alpha(\omega) = \gamma(t)/\sigma(t) = \alpha'(\omega) + i\alpha''(\omega) \quad (3.8)$$

for the case of a simple relaxation process like the one in Equation (3.5).

- Sketch both  $\alpha'(\omega) = \text{Re}(\alpha(\omega))$  and  $\alpha''(\omega) = \text{Im}(\alpha(\omega))$  for this case.
- The real and imaginary parts of the compliance are often referred to as the storage part and the loss part respectively. Discuss the physics behind this distinction.

### Problem 4. Multiple choice questions

1.

The static structure factor  $S(\mathbf{q})$  is a direct measure of:

- a) The positions of the atoms within the unit cell?
- b) The Fourier transform of the order parameter?
- c) The Fourier transform of time-correlations between different particles?
- d) None of the above?

2.

Which of the following types of radiation has the longest penetration depth into materials:

- a) X-rays?
- b) Electrons?
- c) Neutrons?
- d) Visible light?

3.

The order parameter for nematic order of quadrupoles is ( $\langle \rangle$  denotes an average using the appropriate orientational distribution function, and  $\theta = 0$  is the director orientation):

- a)  $\langle (3\sin^2\theta - 1)/2 \rangle$ ?
- b)  $\langle 1 - (3/2)\sin^2\theta \rangle$ ?
- c)  $\langle \cos^2\theta \rangle$ ?
- d) None of the above?

4.

The order parameter for nematic order of dipoles is ( $\langle \rangle$  denotes an average using the appropriate orientational distribution function, and  $\theta = 0$  is the director orientation):

- a)  $\langle (3\cos 2\theta - 1)/2 \rangle$ ?
- b)  $\langle 1 - (3/2)\sin^2\theta \rangle$ ?
- c)  $\langle \cos^2\theta \rangle$ ?
- d) None of the above?

5.

In general, the isotropic to nematic phase transition is:

- a) First order?
- b) Second order?
- c) Without order?
- d) None of the above?

6.

A pure elastic solid is described by:

- a) The Newtonian flow law: stress = viscosity\*shear-rate
- b) Hooke's law: stress = modulus\*strain
- c) A loss modulus that has a peak at the time characteristic for plastic flow
- d) None of the above?



**7.**

Piezoelectricity in a material means that:

- a) There is a linear relationship between elastic and magnetic fields in a material?
- b) There is a non-linear relationship between the elastic and electric fields in a material?
- c) There is no coupling between elastic and electro-magnetic fields in a material?
- d) None of the above?

**8.**

Diamagnetism:

- a) Diamagnetism is the term used for magnetism in materials with permanent magnetic dipoles?
- b) Diamagnetism is present only ferromagnetic materials?
- c) Diamagnetism is present in only in paramagnetic materials?
- d) None of the above?

**9.**

Paramagnetism:

- a) Paramagnetism is present in all materials, except in diamagnetic materials?
- b) The paramagnetic susceptibility is in general very small relative to the diamagnetic susceptibility?
- c) The paramagnetic susceptibility is in general negative?
- d) None of the above?

**10.**

In the ferromagnetic state (i.e.  $T < T_C$ ), the magnetic susceptibility,  $\kappa$ , generally is found to obey a power law of the form  $\kappa \propto (T - T_C)^{-\gamma}$ , where  $\gamma$  is a constant.

It is possible to show that:

- a)  $\gamma = 1$  is always valid for ferromagnets?
- b)  $\gamma$  is system (i.e material) dependent for ferromagnetic transitions?
- c)  $\gamma = 1/2$  is always valid for ferromagnets?
- d) None of above?

**11.**

The superconducting Cooper pairs:

- a) Are fermions?
- b) Are formed by pairs of electrons bound together by exchanging phonons?
- c) Are localized to within typical unit cell dimensions ( $\sim \text{\AA}$ )?
- d) None of the above?

**12.**

The coherence length in a superconductor is defined as:

- a) The decay length for superconducting currents of Cooper pairs into the superconductor?
- b) The length over which a Cooper pair is localized?
- c) The decay length of an external applied magnetic field into a superconductor?
- d) None of the above?

**13.**

Applied magnetic fields are displaced from inside type-I superconducting samples (Meissner effect). The exponential decay length (The London penetration depth) of an applied magnetic field into a type-I superconductor beneath the surface, also describes:

- a) The decay length for superconducting currents of Cooper pairs in the superconductor?
- b) The length over which Cooper pairs are localized?
- c) The decay length for the density of Cooper pairs?
- d) None of the above?

**14.**

The lower critical field  $H_{C1}$  in type-II superconductors describes:

- a) The field for which there is complete penetration of the externally applied field in all the sample?
- b) The field for which the first single flux-line is created?
- c) The field for which the number of flux-lines equals the number of Cooper-pairs?
- d) None of the above?

**15.**

The upper critical field  $H_{C2}$  in type-II superconductors describes:

- a) The field for which there is complete penetration of the externally applied field in all the sample?
- b) The field for which the first single flux-line is created?
- c) The field for which the number of flux-lines equals the number of Cooper-pairs?
- d) None of the above?

**16.**

The frequency positions of the Brillouin doublet peaks measured in Rayleigh-Brillouin experiments provide information about:

- a) Lifetimes of pressure fluctuations?
- c) Lifetimes of heat fluctuations?
- b) Propagation velocities of sound waves?
- d) None of the above?

**17.**

The frequency widths of the Brillouin doublet peaks measured in Rayleigh-Brillouin experiments provide information about:

- a) Lifetimes of pressure fluctuations?
- c) Lifetimes of heat fluctuations?
- b) Propagation velocities of sound waves?
- d) None of the above?

**18.**

Landau theory:

Phase transitions are often described by Landau theory starting from a free energy expansion of the form  $g = g_0 + b(T - T_C)P^2 + C_3P^3 + C_4P^4 + C_6P^6 - EP$ , where  $T$  is the temperature and  $P$  is the order-parameter, and  $T_C$ ,  $b$ ,  $C_2$ ,  $C_3$ ,  $C_4$ ,  $C_6$  are constants.

- a) Landau theory in this form is a general molecular field theory valid only for second order phase transitions independent of the values and signs of the constants  $T_C$ ,  $b$ ,  $C_2$ ,  $C_3$ ,  $C_4$ ,  $C_6$ ?
- b) Landau theory in this form describes second order phase transitions only when  $C_3 < 0$ ,  $C_6 = 0$  and  $C_4 > 0$ ?
- c) Landau theory in this form always describes second order phase transitions provided  $C_3 = 0$ ?
- d) None of the above?

**19.**

For second order phase transition, near the critical temperature  $T_C$ , critical fluctuations may become important for the observed behavior.

Which of the following statements is true?

- a) Critical fluctuations are formally included in mean-field theories?
- b) Critical fluctuations are caused by infinitely large restoring forces giving emergent dynamic and spontaneous regions of finite order parameter near  $T_C$ ?
- c) Critical fluctuations are typically restricted in size to unit cell dimensions ( $\sim \text{\AA}$ )?
- d) None of the above?

**20.**

The Kramers-Kronig (K-K) formulas give relations between the real and imaginary parts of generalized linear response functions, i.e. between the storage and the loss moduli respectively. The physics behind the K-K relations is:

- a) That the system obeys simple Debye relaxation behaviour?
- b) There is no physics behind the K-K relations, they are purely mathematical?
- c) The superposition principle, i.e. the system behavior may be described as a linear combination of fundamental modes?
- d) None of the above?