

Chapter 2 :

Geophysical data processing.

2.1 Introduction (p. 8)

Geophysical data :

measurement of some physical quantity vs time and/or space

Ex : Gravity measurement
Magnetic
seismic
EM

Geophysical data (GP-data) are the sum of

- a) the response from underlying geology
- b) non-GP contrib
- c) instrument noise

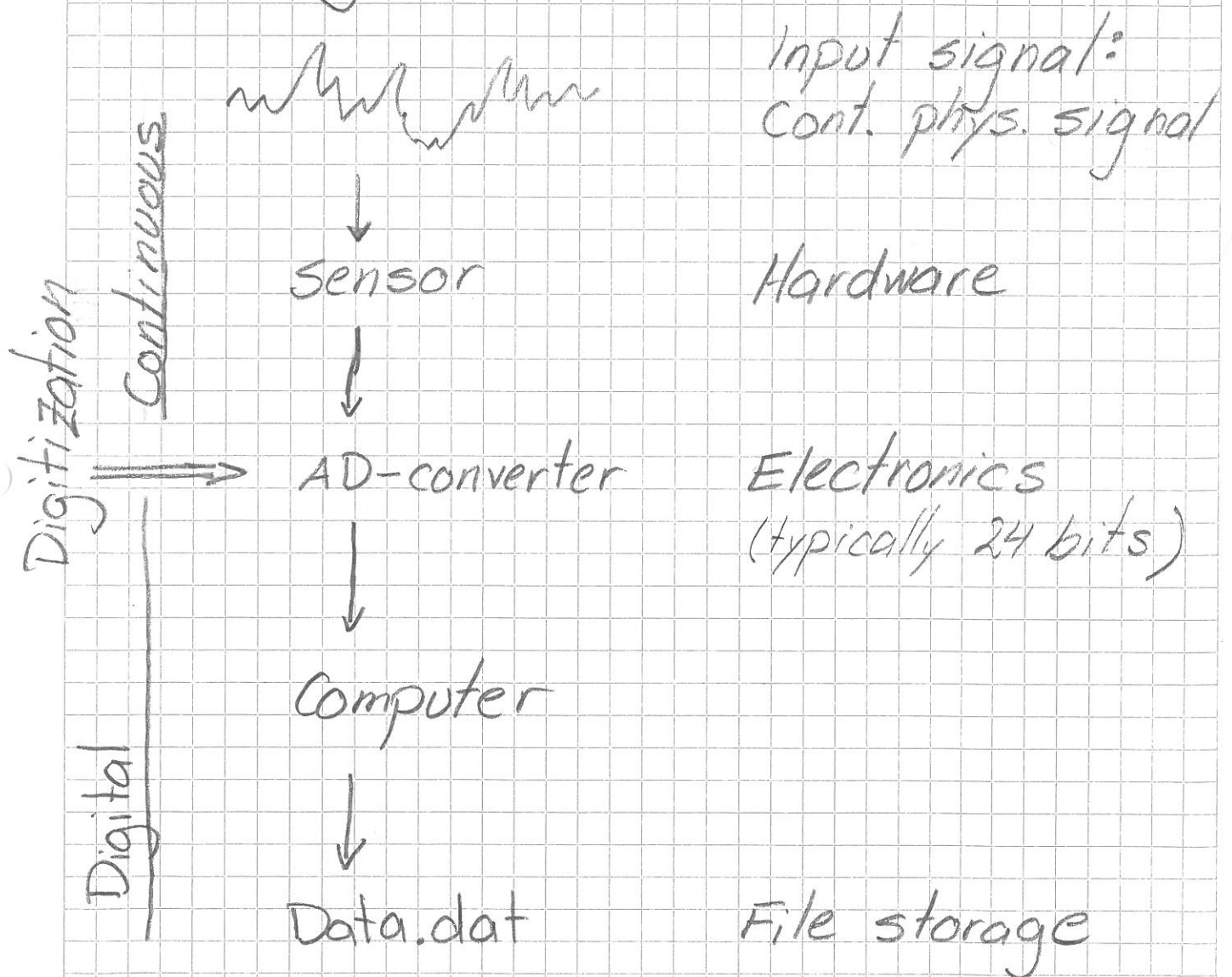
The job of the geophysicist

* Isolate the response from the underlying geology
separate "signal from noise"

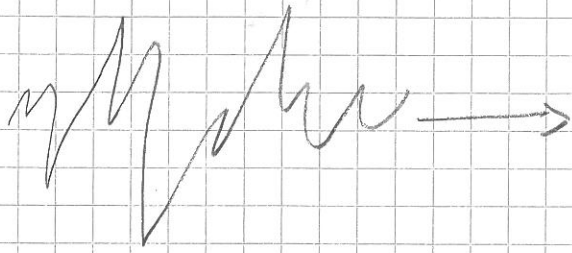
* Process / Interpret the signal
How did the underlying geology look like

* GP-survey planning
"Does one have sufficient signal-to-noise ratio"

2.2 Digitization of GP-data

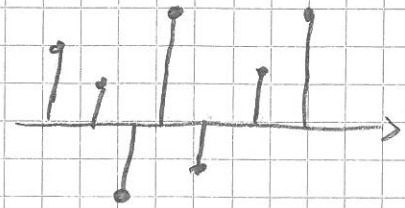


$d(t)$



$d(t_i) = d_i, i \in \mathbb{N}_0$

$t_i = t_0 + i\Delta t$



see Fig 2.2

Q : How accurate are discretely sampled data

A : It depends on the dynamical range and sampling interval.

① Dynamical range :

the ratio between the maximum to the minimum amplitude of a signal that can be rep. accurately.

$$\frac{A_{\max}}{A_{\min}}$$

Typically DR is expressed in decibel (dB)

Power: $P = |A|^2$

$$D_r = 10 \log_{10} \frac{P_2}{P_1}$$

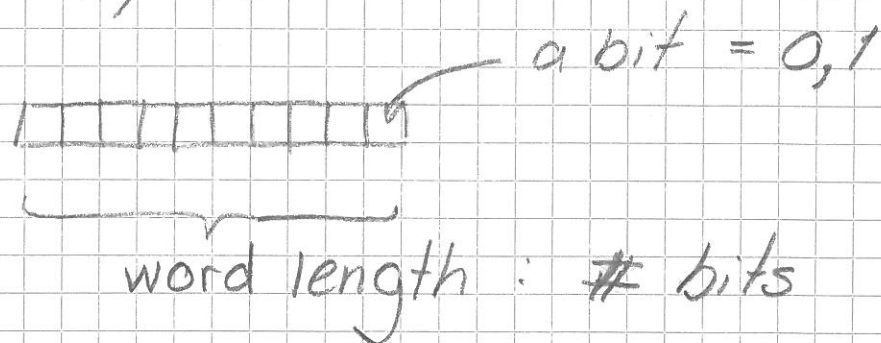
$$= 20 \log_{10} \frac{|A_2|}{|A_1|}$$

Ex: $A_1 = 1$; $A_2 = 1024$

$$D_r = 20 \log_{10} \frac{A_{\max}}{A_{\min}} = 60 \text{ dB}$$

— 11 —

In the computer numbers are represented in binary form:



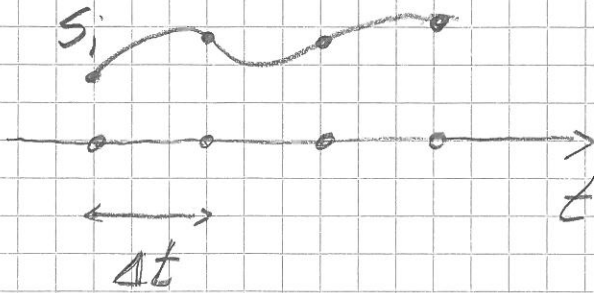
Ex: $D_r = 84 \text{ dB}$; what is the ampl. ratio

$$D_r = 20 \log_{10} \frac{A_{\max}}{A_{\min}}$$

$$\Rightarrow \frac{A_{\max}}{A_{\min}} = 10^{D_r/20} = 2^x$$

$$x = \frac{(D_r/20) \ln 10}{\ln 2} \approx 14 \left| \begin{array}{l} 15\text{-bit} \\ \text{words} \\ \text{needed} \end{array} \right.$$

② Sampling interval / Sampling freq.



Sampling interval (Δt):
time (space) difference between two consecutive samples

Sampling frequency:
of samples per unit time (or distance)

$$f_s = \frac{1}{\Delta t} \quad \left(\text{or } \omega_s = 2\pi f_s \right)$$
$$\left(k_s = 2\pi f_s \right)$$

Q: How can we make sure that our discretely sampled signal d_i represents well the original continuous function

A: We can not in general! ∇

This is only possible under certain requirements.

Nyquist-Shannon sampling theorem

Minimum of two samples per shortest period of the signal

\Rightarrow no loss of information

Nyquist frequency (critical freq.)

$$f_c = \frac{1}{2\Delta t} = \frac{f_s}{2} \quad [f_N \text{ in the book}]$$

or

$$\omega_c = \frac{\pi}{\Delta t} \quad (\omega = 2\pi f)$$

The highest freq component in an analog signal is its bandwidth

Let f_{\max} be the maximum freq contained in an analog signal.

According to the sampling theorem, the analog signal is faithfully represented by its discretely sampled series if

$$f_s \gg 2f_{\max}$$

or

$$2f_c \geq 2f_{max}$$

$$\Rightarrow \boxed{f_c \geq f_{max}} \quad (\omega_c > \omega_{max})$$

$$\text{Ex: } \Delta t = 2 \text{ ms} = 2 \cdot 10^{-3} \text{ s}$$

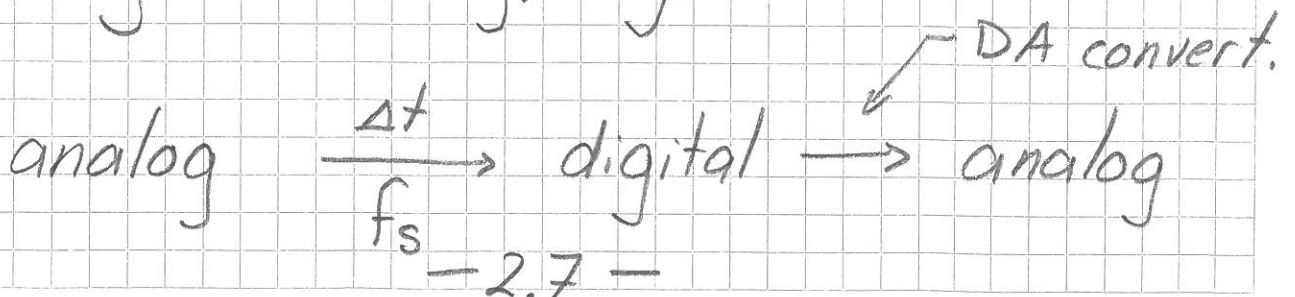
$$f_s = \frac{1}{\Delta t} = 500 \text{ Hz}$$

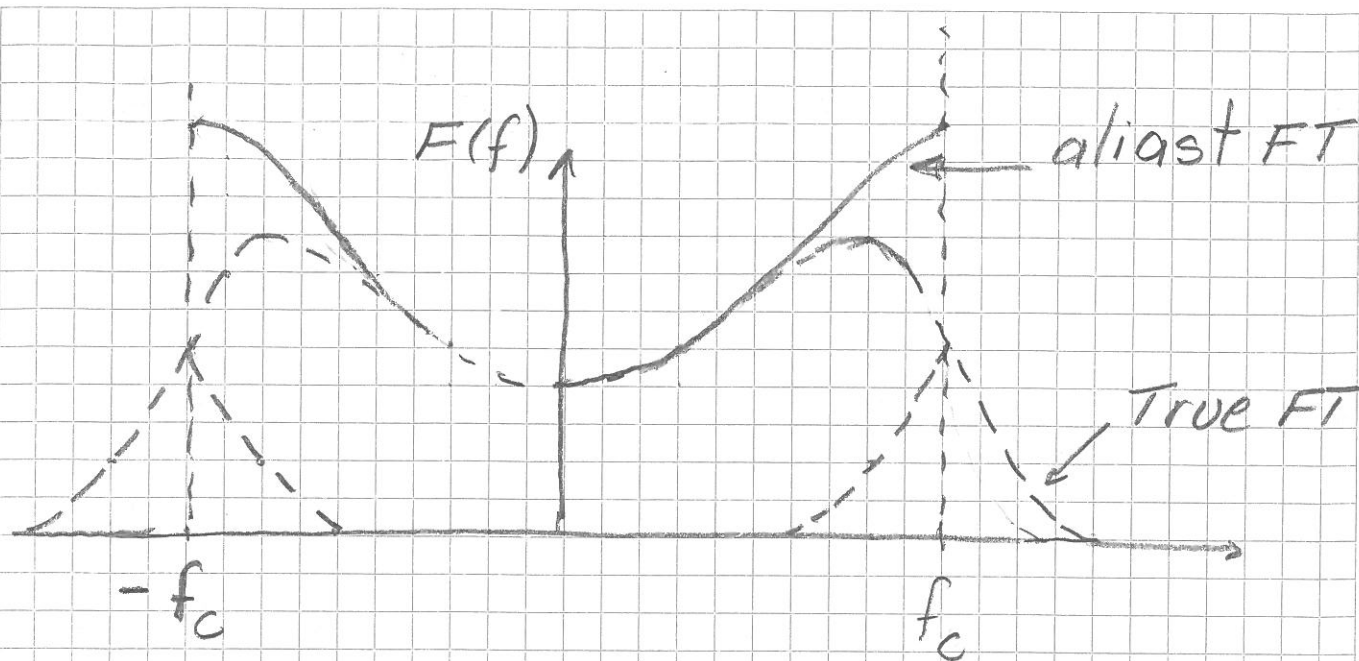
$$f_c = \frac{1}{2\Delta t} = \frac{f_s}{2} = 250 \text{ Hz}$$

Aliasing: a phenomenon that happens if the analog signal has freq. components above the Nyquist freq. i.e., it happens if $f_c < f_{max}$.

see Fig 2.3

Aliasing causes false freq. components to appear that were not in the original (analog) signal



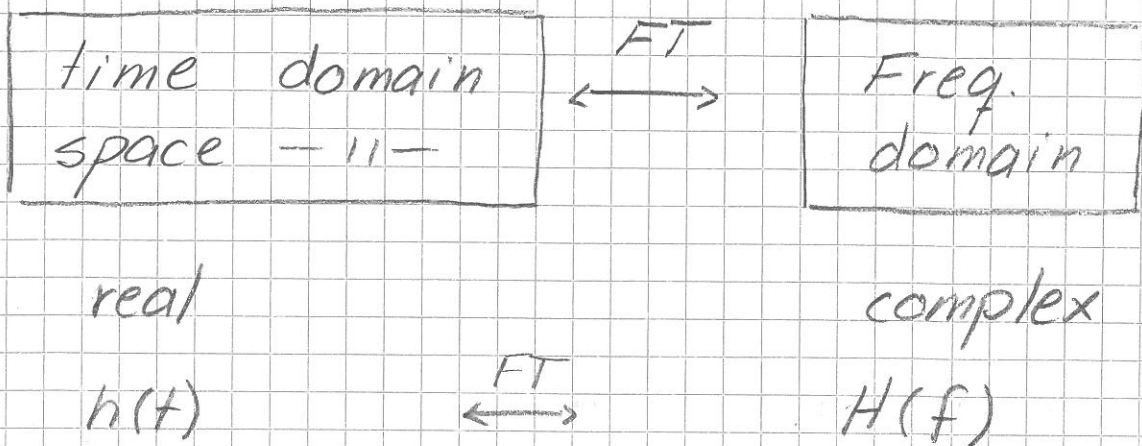


Aliasing is causing a distortion to occur for which $f_{\max} > f_N$ are folded back into the Nyquist interval.

Q : What to do if your signal has $f_{\max} > f_c$

A : Apply a low-pass frequency filter before digitizing so that the input signal has $f_{\max} < f_c$

2.3 Spectral Analysis (p. 10)



Def: Fourier transform:

$$H(f) = \int_{-\infty}^{\infty} dt h(t) e^{i2\pi ft}$$

$$h(t) = \int_{-\infty}^{\infty} df H(f) e^{-i2\pi ft}$$

- || -

$$\omega = 2\pi f \quad d\omega = 2\pi df$$

$$H(\omega) = \int_{-\infty}^{\infty} dt h(t) e^{i\omega t}$$

$$h(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega H(\omega) e^{-i\omega t}$$

$$H(\omega) = [H(f)] \Big|_{f = \omega/2\pi}$$

Typically $h(t)$ is real, so that it follows

$$H(-f) = [H(f)]^*$$

show Fig 2.4

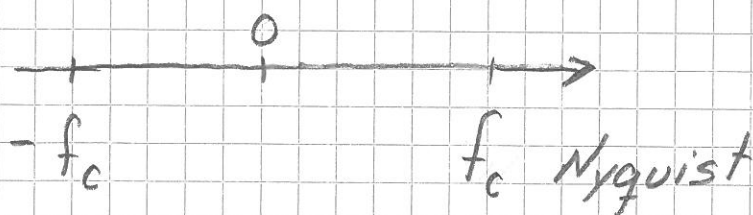
If the function $h(t)$ is periodic of period T

$$h_N(t) = \sum_{n=-N}^N C_n e^{-i \frac{2\pi n t}{T}}$$

$$C_n = \frac{1}{T} \int_{-T/2}^{T/2} dt h(t) e^{-i \frac{2\pi n t}{T}}$$

If one deals with a discretized signal of sampling interval Δt

$$h_k = h(t_k) \quad t_k = k \Delta t \quad k = 1, 2, \dots, N$$



$$\begin{aligned}
 \underbrace{h(t_k)}_{h_k} &= \int_{-\infty}^{\infty} df H(f) e^{-i2\pi f t_k} \\
 &\approx \Delta f \sum_{n=1}^N H(f_n) e^{-i2\pi f_n t_k} \\
 &= \frac{1}{N\Delta t} \sum_{n=1}^N \Delta t H_n e^{-i2\pi f_n t_k} \\
 &= \frac{1}{N} \sum_{n=1}^N H_n e^{-i2\pi f_n t_k}
 \end{aligned}$$

Note :

$$f_n t_k = \frac{n}{N\Delta t} k \Delta t = \frac{nk}{N}$$

— 11 —

Since

$$e^{i\theta} = \cos\theta + i\sin\theta$$

$h_k = h(t_k)$ is decomposed into a sum of sine and cosines.

This is at the heart of Fourier anal.

$$\Delta f = \frac{2f_c}{N} = \frac{2}{N2\Delta t} = \frac{1}{N\Delta t}$$

Therefore, we have discrete freq.

$$f_n = n\Delta f = \frac{n}{N\Delta t}, \quad n = -\frac{N}{2}, \dots, \frac{N}{2}$$

$$H(f_n) = \int_{-\infty}^{\infty} dt h(t) e^{i2\pi f_n t}$$

$$\approx \Delta t \sum_{k=1}^N h_k e^{i2\pi f_n t_k}$$

$$\equiv \Delta t H_n$$

The discrete FT maps N numbers h_k in time space into N numbers H_n in Fourier space.

The inverse transform:

$$h_k = \frac{1}{N} \sum_{n=1}^N H_n e^{-i2\pi f_n t_k}$$

which follows from $h(t_k)$

Fourier transforms converts

$$h(t) \xleftrightarrow[ET]{} H(f)$$

↑
Fourier spectrum

$$H(f) = A(f) e^{i\phi(f)}$$

$A(f)$: real amplitude

$\phi(f)$: real phase

Fourier pairs

$$h(t) \leftrightarrow H(f)$$

Ex: Fig 2.8

Geophysical use : emgs

Fast Fourier Transform (FFT)

FFT algorithm: Cooley-Tukey (1974)

If

$$\vec{T} = [t_k] \quad (\text{time-domain})$$

$$\vec{H} = [H_n] \quad (\text{fourier-domain})$$

then

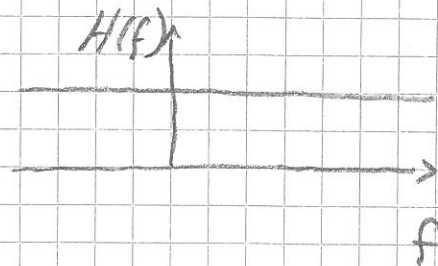
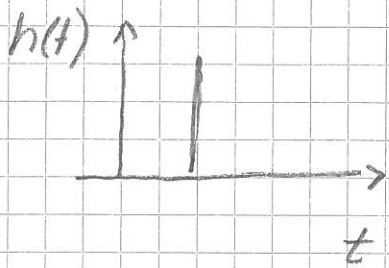
$$\text{FFT}(\vec{T}) = \vec{H}$$

— " —

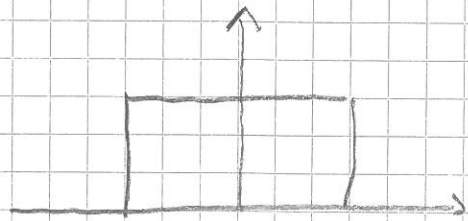
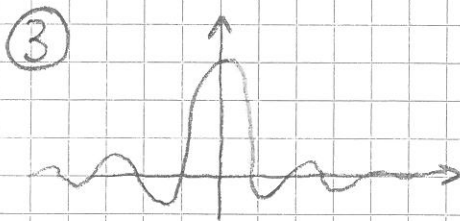
The FT can be generalized to higher dimensions

Note the following Fourier pairs

① $\delta(t) \longleftrightarrow 1$



② $1 \longleftrightarrow \delta(f)$



③ $\text{sinc}(\pi t) \longleftrightarrow \text{rect}(f)$
 $= \frac{\sin(\pi t)}{\pi t}$ (rectangular func)
 (puls-function)

④ $\text{rect}(t) \longleftrightarrow \text{sinc}(\pi f)$

$$\text{rect}(t) = \begin{cases} 1 & |t| < 1/2 \\ 1/2 & |t| = 1/2 \\ 0 & |t| > 1/2 \end{cases}$$

2.4 Convolution (p. 13)

Def: Convolution of two (continuous) functions $g(t)$ and $h(t)$ is

$$g(t) * h(t) = \int_{-\infty}^{\infty} d\tau g(\tau) h(t-\tau)$$

— " —

You will also see the notation $(g * h)(t)$.

It can be shown (by a change of var.) that

$$g(t) * h(t) = h(t) * g(t).$$

— " —

Lets calculate the conv at $t_k = k\Delta t$, $k=1,2,3,\dots$

$$y(t_k) = g(t_k) * h(t_k)$$

$$= \int_{-\infty}^{\infty} d\tau g(\tau) h(t_k - \tau)$$

$$\approx \Delta t \sum_{i=1} g(\tau_i) h(t_k - \tau_i), \quad \tau_i = i\Delta t$$

— 2.16 —

$$= \Delta t \sum_{i=1} g_i h_{k-i}$$

In a discrete convolution one usually neglects the Δt -term and define for

$$\{g_i\}_{i=1}^m \quad \text{and} \quad \{h_j\}_{j=1}^n$$

$$y_k = \sum_{i=1}^m g_i h_{k-i} \quad k = 1, 2, \dots, m+n-1$$

Note: there are $m+n-1$ y_k values

Since $g * h = h * g$ it follows that

$$y_k = \sum_{i=1}^n h_i g_{k-i}$$

can alternatively used to calculate the convolution.

Ex: Calculate the convolution of the two vectors

$$\vec{g} = [2, 0, 1] \quad m = 3$$

$$\vec{h} = [4, 3, 2, 1] \quad n = 4$$

See Fig. 2.12

$$\vec{y} = \text{conv}(\vec{g}, \vec{h}) \quad [\text{matlab}]$$
$$= [8, 6, 8, 5, 2, 1] \quad 6\text{-elements}$$

— 11 —

The Convolution theorem

$$\mathcal{F}(g(t) * h(t)) = \mathcal{F}(g(t)) \cdot \mathcal{F}(h(t))$$
$$= G(f) \cdot H(f)$$

\mathcal{F} : The Fourier Transform

Octave / Matlab

$$\text{conv}(\vec{g}, \vec{h})$$

$$\text{fftconv}(\vec{g}, \vec{h}) \quad - \text{uses FFT}$$

— 2.18 —

Q : what is the meaning of the convolution operator

A : It describes the effect on an input signal after passing through a filter



Ex : What do you hear when the guy next-door is playing loud music?

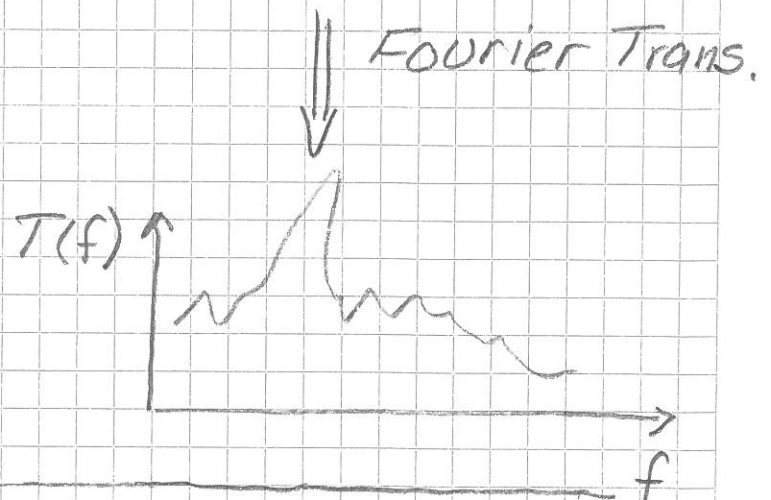
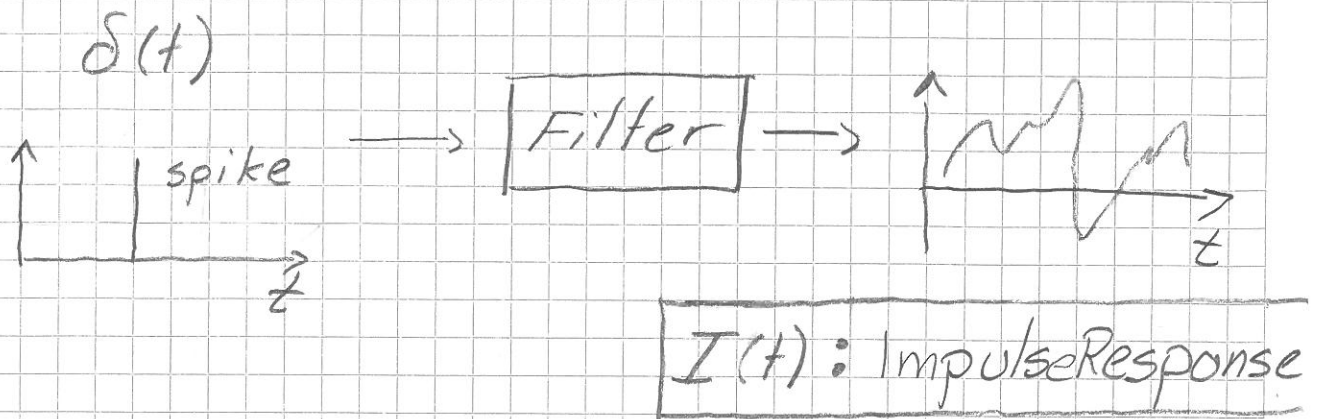
You hear too much bass since the wall acts as a low-pass frequency filter.

Also see the animation on the wikipedia page dedicated to conv.

2.4.1 Impulse response / Transfer funct.

Input signal
(time-domain)

Output signal
(time-domain)



$T(f) : \text{Transfer Function}$

Both the Impulseresponce, $I(t)$, and the Transfer function, $T(f)$, characterize the filter and they contain the same amount of information.

Ex : Fig. 2.9

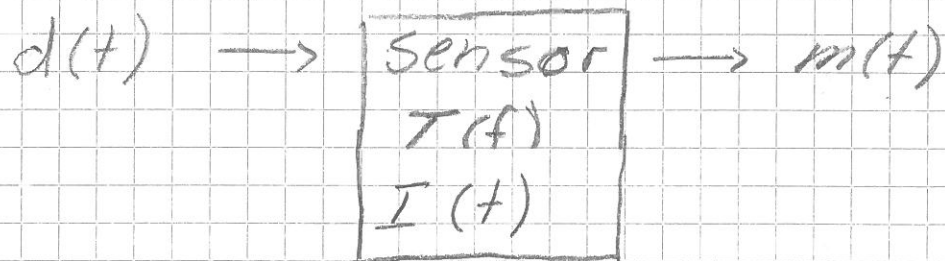
Note :

$I(t)$: time-domain

$T(f)$: frequency domain

$$\boxed{T(f) = \mathcal{F}(I(t))} \quad I(t) = \mathcal{F}^{-1}(T(f))$$

Ex : We measure an input signal, $d(t)$, using a sensor of transfer function, $T(f)$, and obtain the output (measured) signal, $m(t)$.



In the Fourier domain we have :

$$M(f) = T(f) \cdot D(f)$$

which in the time-domain becomes after using the convolution theorem :

$$m(t) = I(t) * d(t).$$

We observe that in general $m(t) \neq d(t)$.

2.4.2 Deconvolution (p. 16)

Q: From the measured signal $m(t)$, how can one obtain the input sig. $d(t)$?

A: Deconvolution (also called inv. filtering) solves this issue and works like this

$$m(t) = I(t) * d(t)$$

↓ FT

$$M(f) = T(f) D(f)$$

$$\Rightarrow D(f) = \frac{1}{T(f)} M(f) = T^{-1}(f) M(f)$$

$$d(t) = \mathcal{F}^{-1}(T^{-1}(f) M(f))$$

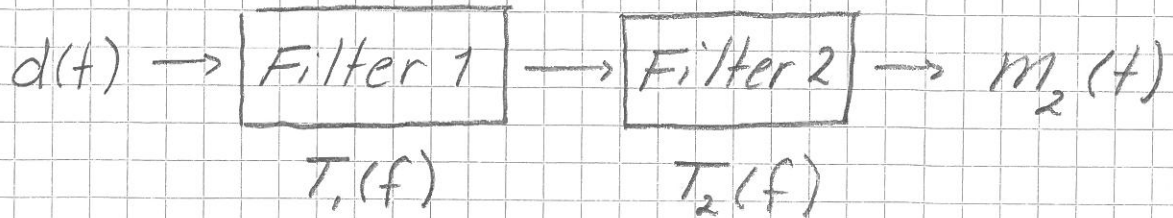
Alternatively, one may ask for an impulse response $J(t)$ for deconvolution process that take $m(t)$ into $d(t)$, i.e

$$d(t) = J(t) * m(t)$$

or after using $m(t) = I(t) * d(t)$ it follows

$$d(t) = \underbrace{J(t) * I(t)}_{\delta(t)} * d(t)$$

Note: Applying two filters



$$M_2(f) = T_2(f) \cdot T_1(f) \cdot D(f)$$

Still $m_2(t) = \mathcal{F}^{-1}(M_2(f))$, but the inv. Fourier transform is more complicated to obtain.

The total transfer function of combination of filter (or sensor) 1 and 2 is

$$T(f) = T_2(f) T_1(f)$$

so that the corresponding total impulse response is

$$\begin{aligned} I(t) &= \mathcal{F}^{-1}(T_2(f) \cdot T_1(f)) \\ &= \underline{T_2(t)} * I_1(t) \end{aligned}$$

2.4.3 Application of convolutions in GP

- We have already seen the use conv. due to a non-trivial impulse response of a sensor. (instrument response)
- However, for instance, the change in a seismic wave as it passes through the ground can be viewed as a filtering process.
 - with the ground we associate an impulse response
 - deconvolution is an essential part of seismic data processing (improve data quality)

Q : Why is deconvolution of GP data challenging

A : The impulse response (or transfer function) of the Earth is typically not known.

2.5 Correlation (p. 16)

Definition:

Given two functions $g(t)$ and $h(t)$, then their correlation function is defined as

$$\phi_{gh}(t) = \int_{-\infty}^{\infty} d\tau g(\tau+t) h(\tau)$$

where t is known as the lag.

— 11 —

If $g = h$ we talk about the auto-correlation function.

In the discrete case it follows

$$\begin{aligned}\phi_{gh}(t_k) &= \int_{-\infty}^{\infty} d\tau g(\tau+t_k) h(\tau) \\ &\approx \Delta t \sum_i g(\tau_i+t_k) h(\tau_i) \\ &= \Delta t \sum_i g_{i+k} h_i\end{aligned}$$

As for the convolution, one normally suppress Δt and under the assumption

— 2.25 —

$$\{g_i\}_{i=1}^n \quad \{h_i\}_{i=1}^n$$

write

$$\phi_{gh}(\tau_k) = \sum_{i=1}^{n-k} g_{i+k} h_i, \quad -m \leq k \leq m$$

— 11 —

The correlation function $\phi_{gh}(\tau)$ quantify how similar the functions $g(t)$ and $h(t)$ are.

see Fig. 2.13

where $\phi_{gh}(\tau) = 0$ (for given τ) we say that g and h are uncorrelated.

The auto-correlation function has the property

$$\phi_{xx}(\tau) = \phi_{xx}(-\tau)$$

The auto-correlation function ϕ_{xx} contains all the amplitude information of $x(t)$ but not any phase information.

Mathematically this is expressed by the so-called Wiener-Khinchin theorem stating that

$$\phi_{xx}(t) \xleftrightarrow{FT} |\hat{\phi}_{xx}(f)|^2$$

where $\hat{\phi}(f) = \mathcal{F}(\phi_{xx}(t))$.

The quantity

$$P(f) = |G(f)|^2$$

is known as the power spectrum of the function $g(t)$.

The correlation theorem

$$\text{corr}(g, h)_k \longleftrightarrow G_i(f) H_i^*(f)$$

For a real signal $x(t)$ it follows from the properties of the Fourier transform that

$$\hat{x}(f) = \mathcal{F}[x(t)]$$

$$\hat{x}^*(f) = \hat{x}(-f)$$

— " —

The power spectrum in more detail:

$$\hat{x}(f) = \mathcal{F}[x(t)]$$

$$P(f) \stackrel{\text{Def.}}{=} |\hat{x}(f)|^2$$

$$= \hat{x}^*(f) \hat{x}(f)$$

$$= \mathcal{F}[\mathcal{F}^{-1}[\hat{x}^*(f) \hat{x}(f)]]$$

$$= \mathcal{F}[\phi_{xx}(t)]$$

2.5.1 Applications of correlations in GP

- Detection of weak signals embedded in noise
- Detecting periodicity (often via the power spectrum)
see Fig 2.15
- Detection of multiple reflections in seismic reflection data.

2.6 Digital Filters (p. 17)

$$d(t) = s(t) + n(t)$$

$$\text{data} = \text{signal} + \text{noise}$$

signal: the part of the waveform data related to the geological structure under investigation

noise: all other components of the waveform

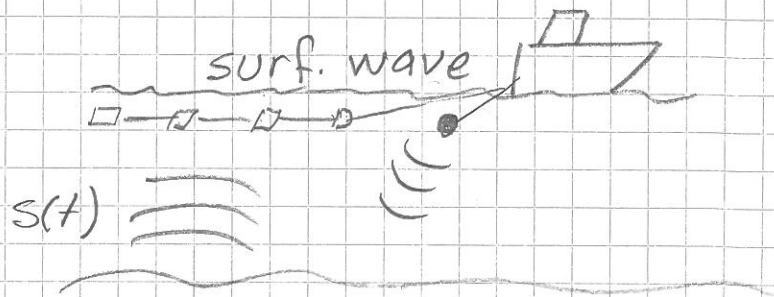
Two types of noise:

- random noise:
statistically random signal typically unrelated to the GP survey

Ex: colored instrument noise
background vibration
(due to wind, waves etc)

- coherent noise:
components of the waveform generated by the GP experim. but not of particular interest for the geological interpretation

Ex.



Signal-to-noise ratio (SNR)

$$\frac{s(t)}{n(t)}$$

One want this signal to have a large amplitude.

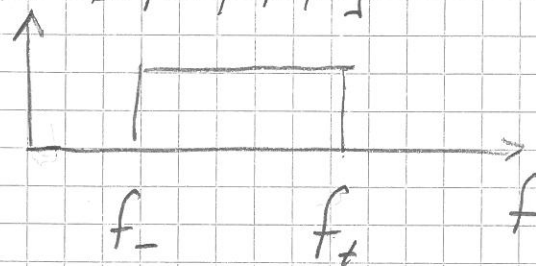
For low SNR one tries to increase the signal by processing to filter out eg coherent noise.

Two types of filters:

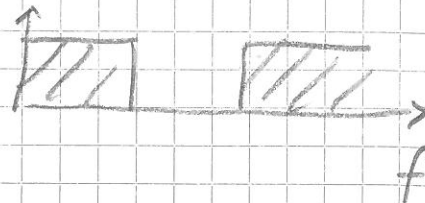
- Freq. Filters
- Inverse (deconvolution) filters

Frequency Filters

Lets through freq. components of the input signal only in a finite interval $[f_-, f_+]$



- low-pass (LP) filter : $f_- = 0$
- high-pass (HP) filter : $f_+ = \infty$
- band-pass (BP) filter
- band-reject



Inverse Filters:

Remove the effect of filters applied previously
(more in Ch 4)

2.7 Imaging and Modelling (Inversion) p 19.

Geophysical waveform



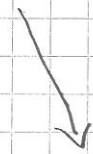
Processing



Data with optimal SNR



Imaging
(data presented)



Modelling/Inversion



Geological Interpretation