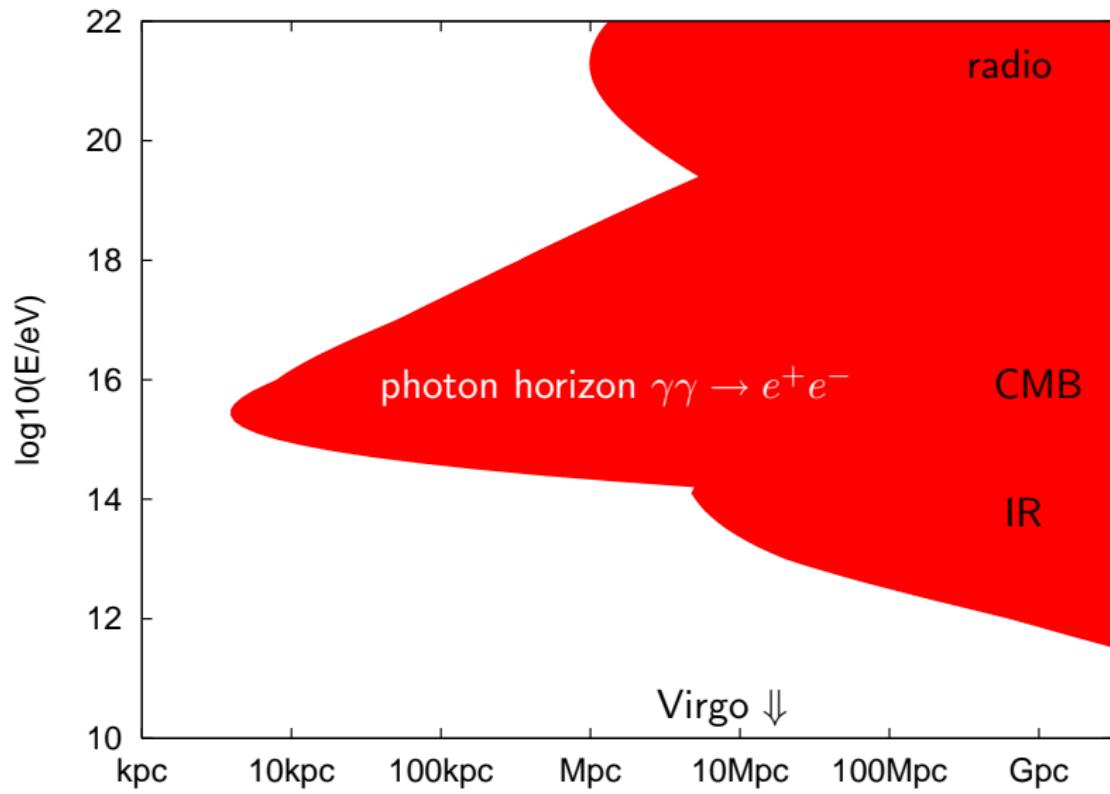
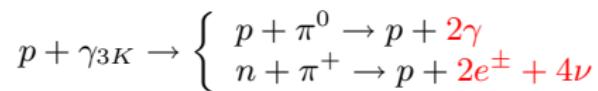


Mean free path of photons



Origin of cascade photons:

- UHECRs:
 - ▶ Photon and neutrino production relatively tight connected:
 - ★ protons:



Origin of cascade photons:

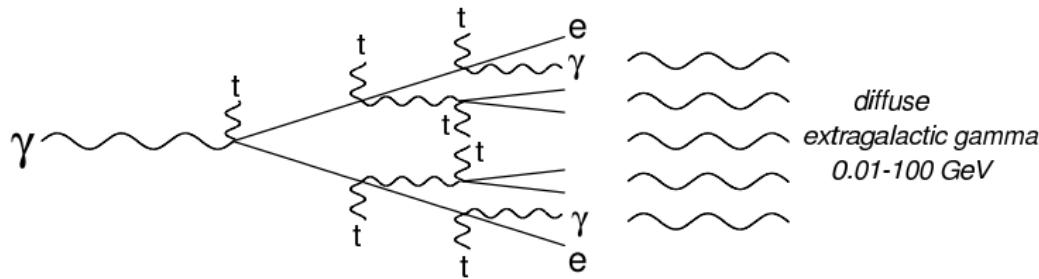
- UHECRs:
 - ▶ Photon and neutrino production relatively tight connected:
 - ★ protons:
$$p + \gamma_{3K} \rightarrow \begin{cases} p + \pi^0 \rightarrow p + 2\gamma \\ n + \pi^+ \rightarrow p + 2e^\pm + 4\nu \end{cases}$$
 - ★ nuclei: $A + \gamma_{3K} \rightarrow (A - 1) + n \rightarrow (A - 1) + p + e^- + \nu_e$
 - ★ connection to UHECRs looser

Origin of cascade photons:

- UHECRs:
 - ▶ Photon and neutrino production relatively tight connected:
 - ★ protons:
$$p + \gamma_{3K} \rightarrow \begin{cases} p + \pi^0 \rightarrow p + 2\gamma \\ n + \pi^+ \rightarrow p + 2e^\pm + 4\nu \end{cases}$$
 - ★ nuclei: $A + \gamma_{3K} \rightarrow (A - 1) + n \rightarrow (A - 1) + p + e^- + \nu_e$
- HE and VHE photons from AGNs

Origin of cascade photons:

- UHECRs:
 - ▶ Photon and neutrino production relatively tight connected:
 - ★ protons:
$$p + \gamma_{3K} \rightarrow \begin{cases} p + \pi^0 \rightarrow p + 2\gamma \\ n + \pi^+ \rightarrow p + 2e^\pm + 4\nu \end{cases}$$
 - ★ nuclei: $A + \gamma_{3K} \rightarrow (A - 1) + n \rightarrow (A - 1) + p + e^- + \nu_e$
- HE and VHE photons from AGNs



Diffuse cascade flux:

- analytical estimate:

[Berezinsky, Smirnov '75]

$$J_\gamma(E) = \begin{cases} K(E/\varepsilon_X)^{-3/2} & \text{at } E \leq \varepsilon_X \\ K(E/\varepsilon_X)^{-2} & \text{at } \varepsilon_X \leq E \leq \varepsilon_a \\ 0 & \text{at } E > \varepsilon_a \end{cases}$$

- three regimes:

Diffuse cascade flux:

- analytical estimate:

[Berezinsky, Smirnov '75]

$$J_\gamma(E) = \begin{cases} K(E/\varepsilon_X)^{-3/2} & \text{at } E \leq \varepsilon_X \\ K(E/\varepsilon_X)^{-2} & \text{at } \varepsilon_X \leq E \leq \varepsilon_a \\ 0 & \text{at } E > \varepsilon_a \end{cases}$$

- three regimes:

- Thomson cooling:

$$E_\gamma = \frac{4}{3} \frac{\varepsilon_{bb} E_e^2}{m_e^2} \approx 100 \text{ MeV} \left(\frac{E_e}{1 \text{ TeV}} \right)^2$$

- plateau region
- above pair-creation threshold $s_{min} = 4E_\gamma\varepsilon_{bb} = 4m_e^2$:
flux exponentially suppressed

Diffuse flux, analytical estimate for low-energy part:

- $q_i(E)$: # particles crossing energy E

Diffuse flux, analytical estimate for low-energy part:

- $q_i(E)$: # particles crossing energy E

- cooling regime:

no generation of electrons for $\varepsilon < \varepsilon_a/2$: $q_e(E_e) = q_0$

$$E_\gamma \propto E_e^2 \Rightarrow dE_\gamma \propto E_e dE_e$$

Diffuse flux, analytical estimate for low-energy part:

- $q_i(E)$: # particles crossing energy E

- cooling regime:

no generation of electrons for $\varepsilon < \varepsilon_a/2$: $q_e(E_e) = q_0$

$$E_\gamma \propto E_e^2 \Rightarrow dE_\gamma \propto E_e dE_e$$

inserting in **energy conservation**,

$$E_\gamma dn_\gamma = q_e(E_e) dE_e ,$$

gives

$$J(E_\gamma) \propto E_\gamma^{-3/2}$$

Diffuse flux, analytical estimate for plateau region:

- energy conservation and $N_e/N_\gamma = \text{const.}$

$$\Rightarrow q_i(E_i)E_i = \text{const} \quad \Rightarrow \quad q_e(E_e) \propto 1/E_e$$

Diffuse flux, analytical estimate for plateau region:

- energy conservation and $N_e/N_\gamma = \text{const.}$

$$\Rightarrow q_i(E_i)E_i = \text{const} \quad \Rightarrow \quad q_e(E_e) \propto 1/E_e$$

- IC regime:

$$E_\gamma = \frac{4E_e}{3 \ln(2E_e \varepsilon_{\text{bb}}/m_e^2)}$$

Diffuse flux, analytical estimate for plateau region:

- energy conservation and $N_e/N_\gamma = \text{const.}$

$$\Rightarrow q_i(E_i)E_i = \text{const} \quad \Rightarrow \quad q_e(E_e) \propto 1/E_e$$

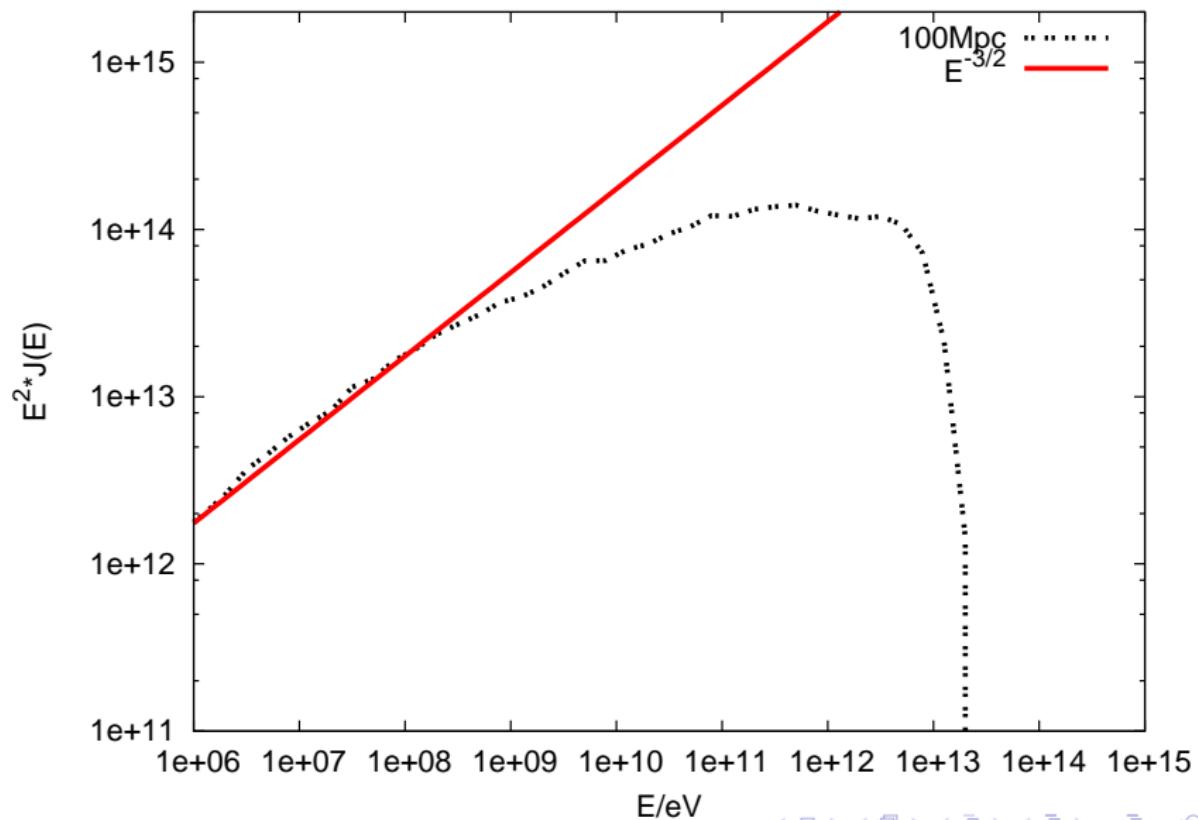
- IC regime:

$$E_\gamma = \frac{4E_e}{3 \ln(2E_e \varepsilon_{\text{bb}}/m_e^2)}$$

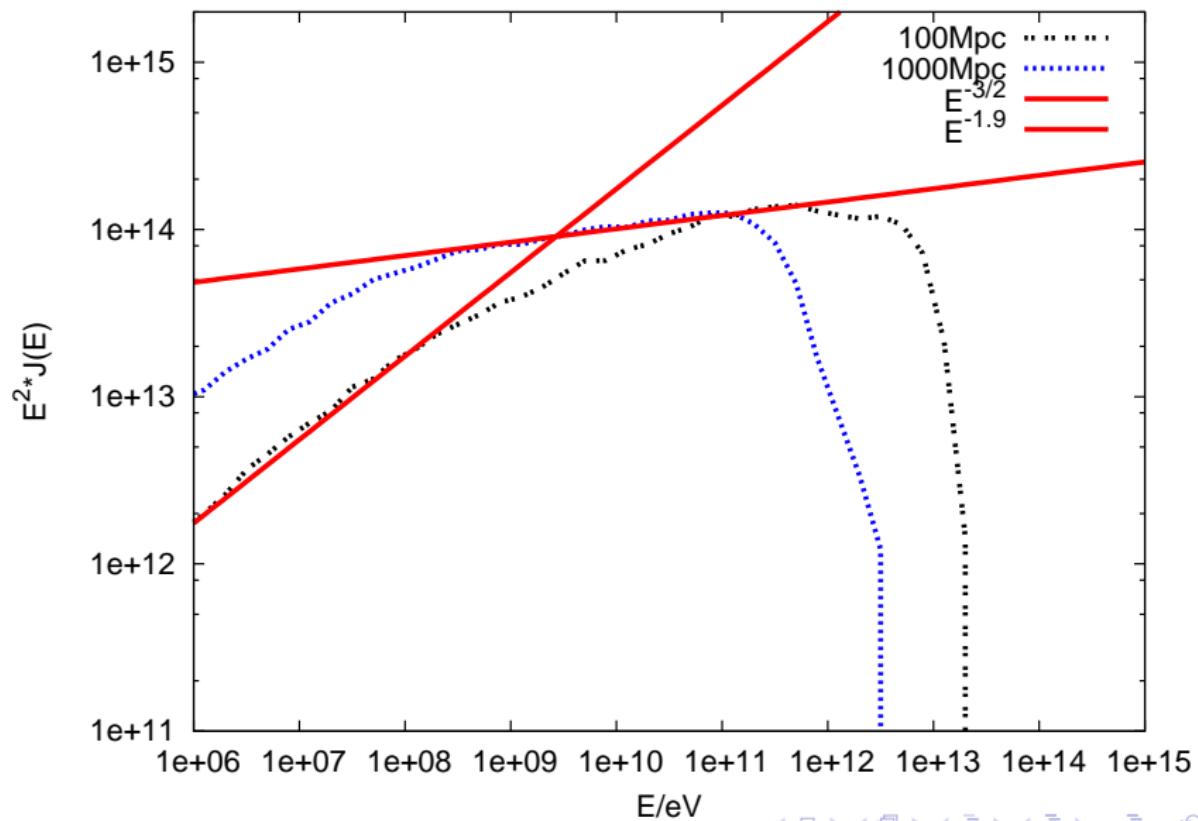
- to log. accuracy

$$J(E_\gamma) \propto E_\gamma^{-2}$$

Monte Carlo vs. analytical estimate: single source

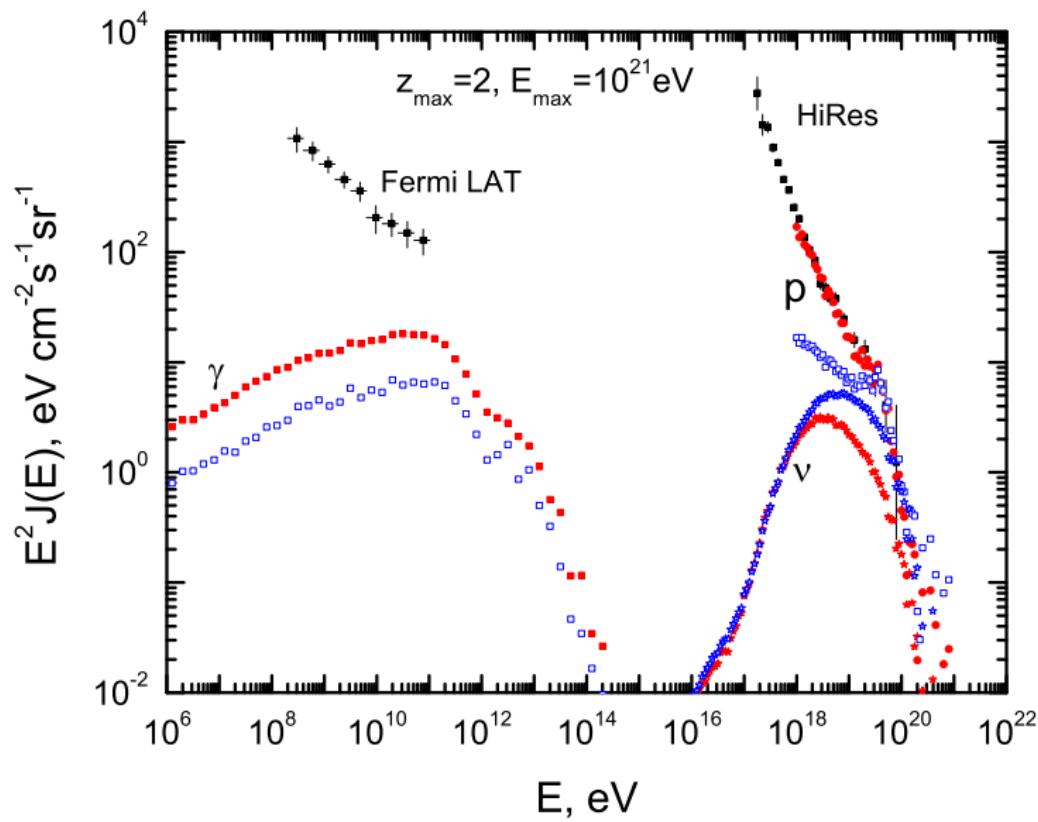


Monte Carlo vs. analytical estimate: single source



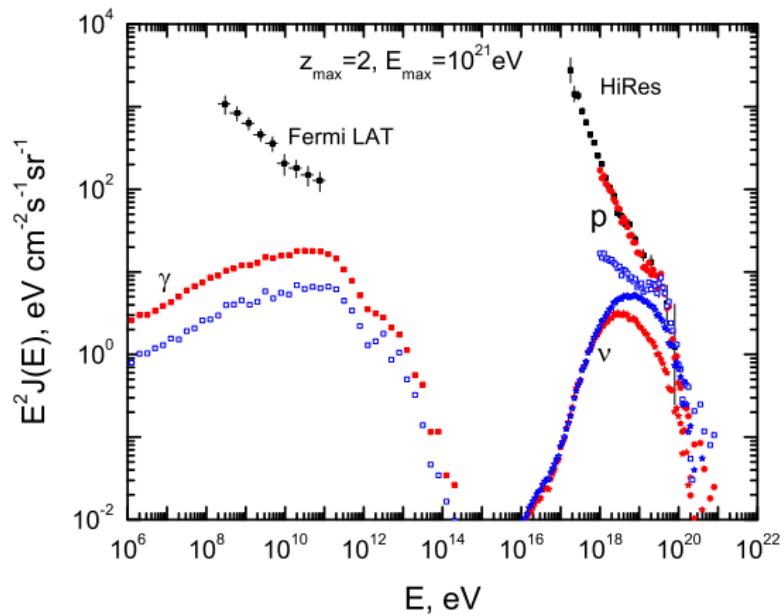
Fermi-LAT vs. UHECR data: no evolution

[Berezinsky et al. '10]



Fermi-LAT vs. UHECR data: no evolution

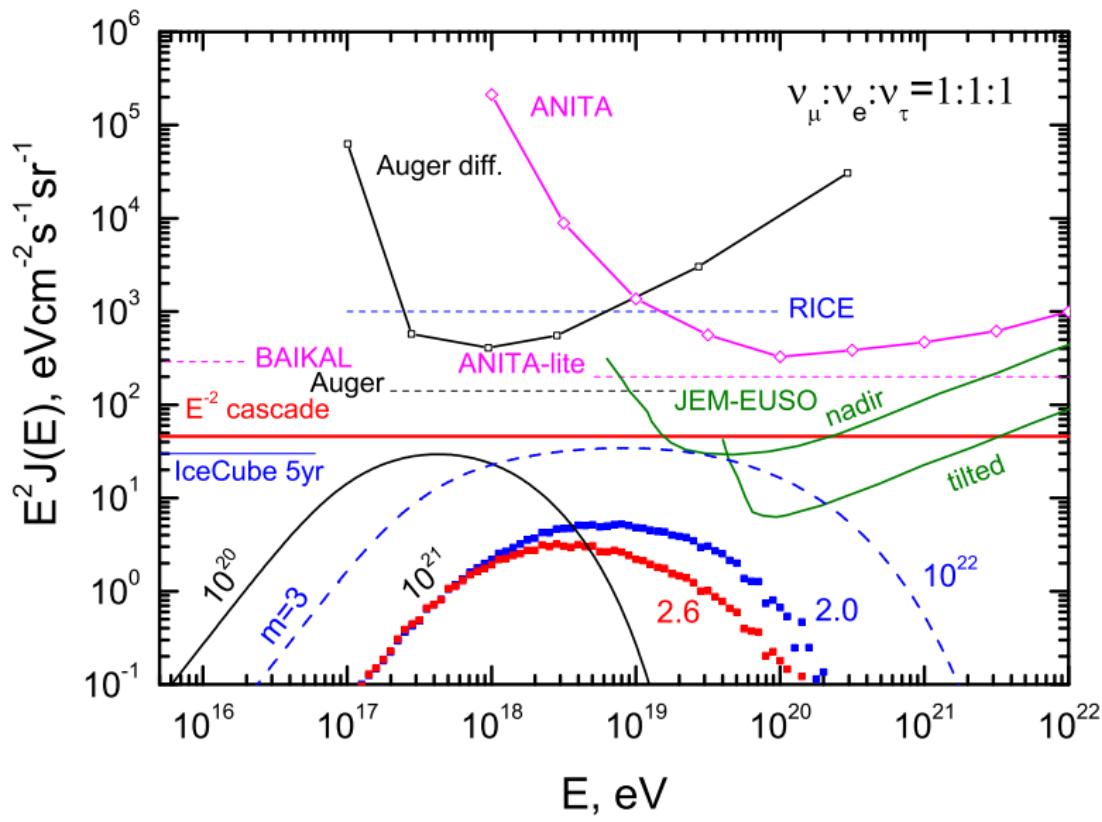
[Berezinsky et al. '10]



integrating $EJ(E)$ gives bound $\omega_{\text{cas}} \lesssim 6 \cdot 10^{-7} \text{ eV/cm}^3$

Cascade limit for cosmogenic neutrinos

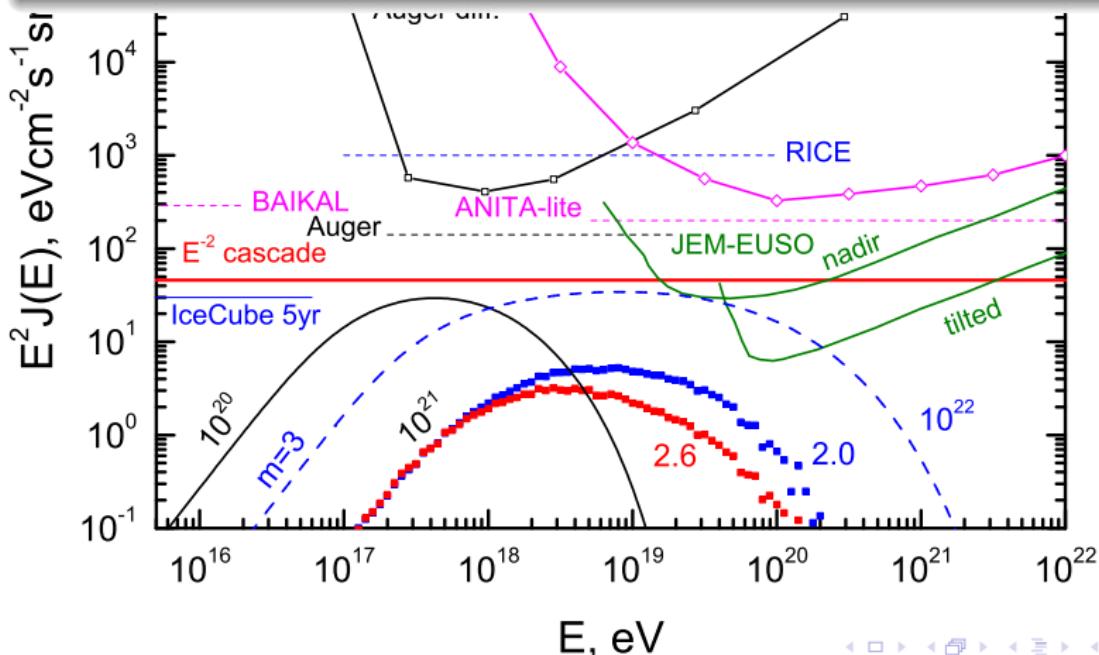
[Berezinsky et al. '10]



Cascade limit for cosmogenic neutrinos

Assumes proton primaries

- for nuclei reduced neutrino fluxes...



Gamma-rays and extragalactic magnetic fields (EGMF)

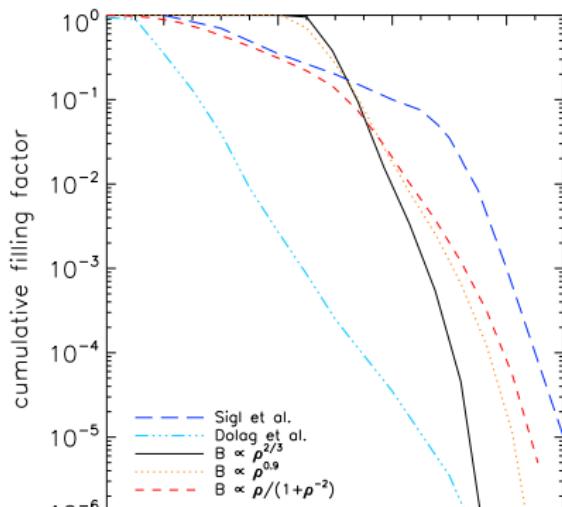
- Observations only in clusters,
 - ▶ synchrotron halo: $\Rightarrow B \sim (0.1 - 1) \mu\text{G}$
 - ▶ Faraday rotation: $\Rightarrow B \sim (1 - 10) \mu\text{G}$

Gamma-rays and extragalactic magnetic fields (EGMF)

- Observations only in clusters,
 - ▶ synchrotron halo: $\Rightarrow B \sim (0.1 - 1) \mu\text{G}$
 - ▶ Faraday rotation: $\Rightarrow B \sim (1 - 10) \mu\text{G}$
- Origin of **seed** for EGMF is **mysterious**

Gamma-rays and extragalactic magnetic fields (EGMF)

- Observations only in clusters,
 - ▶ synchrotron halo: $\Rightarrow B \sim (0.1 - 1) \mu\text{G}$
 - ▶ Faraday rotation: $\Rightarrow B \sim (1 - 10) \mu\text{G}$
- Origin of seed for EGMF is mysterious
- Seed required as **input for EGMF simulations**



Gamma-rays and extragalactic magnetic fields (EGMF)

- Observations only in clusters,
 - ▶ synchrotron halo: $\Rightarrow B \sim (0.1 - 1) \mu\text{G}$
 - ▶ Faraday rotation: $\Rightarrow B \sim (1 - 10) \mu\text{G}$
- Origin of seed for EGMF is mysterious
- Seed required as input for EGMF simulations
- Aharonian, Coppi, Völk '94: **Pair halos** around AGNs

Gamma-rays and extragalactic magnetic fields (EGMF)

- Observations only in clusters,
 - ▶ synchrotron halo: $\Rightarrow B \sim (0.1 - 1) \mu\text{G}$
 - ▶ Faraday rotation: $\Rightarrow B \sim (1 - 10) \mu\text{G}$
- Origin of seed for EGMF is mysterious
- Seed required as input for EGMF simulations
- Aharonian, Coppi, Völk '94: Pair halos around AGNs
- Plaga '95: **EGMFs deflect and delay cascade electrons**
 \Rightarrow search for delayed “echoes” of multi-TeV AGN flares/GRBs

Influence of EGMF on flux from single source: deflections

- deflection of electrons:

$$\vartheta \sim \frac{l_{\text{cool}}}{R_L} \propto E_e^{-2}$$

Influence of EGMF on flux from single source: deflections

- deflection of electrons:

$$\vartheta \sim \frac{l_{\text{cool}}}{R_L} \propto E_e^{-2}$$

\Rightarrow flux within angle ϑ reduced by factor E^2

Influence of EGMF on flux from single source: deflections

- deflection of electrons:

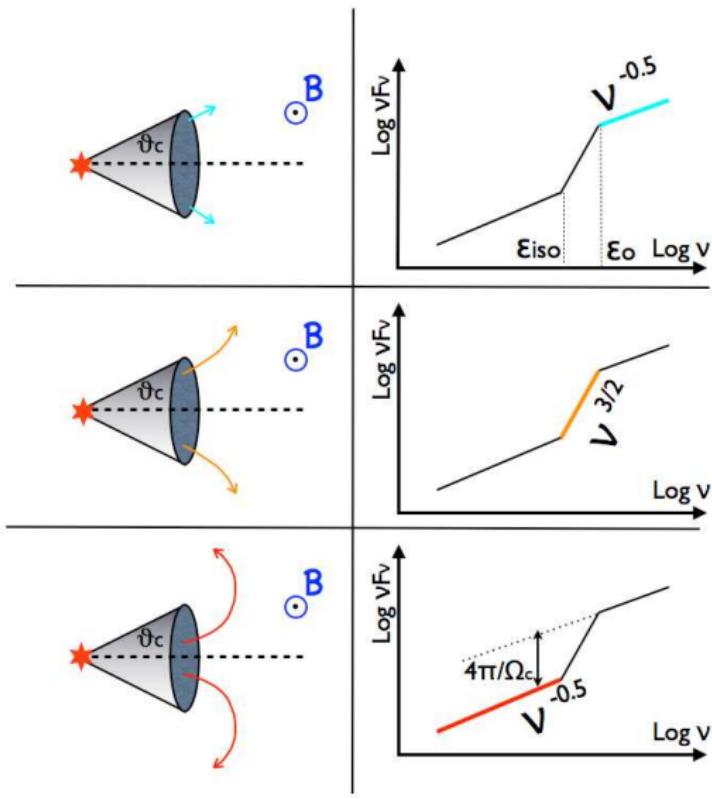
$$\vartheta \sim \frac{l_{\text{cool}}}{R_L} \propto E_e^{-2}$$

⇒ flux within angle ϑ reduced by factor E^2

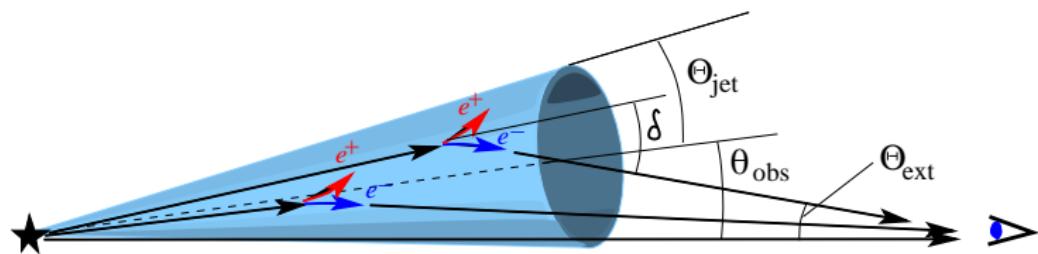
⇒ cooling regime: transition from

$$J(E) \propto E^{-1.5} \rightarrow E^{0.5} \rightarrow E^{-1.5}$$

Influence of EGMF on flux from single source: deflections

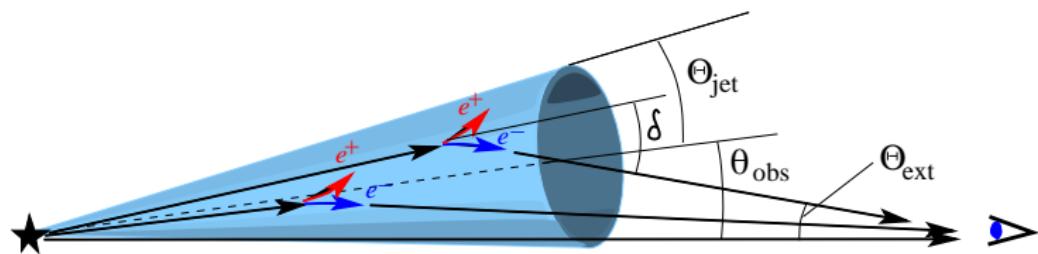


Influence of EGMF on flux from single source: time



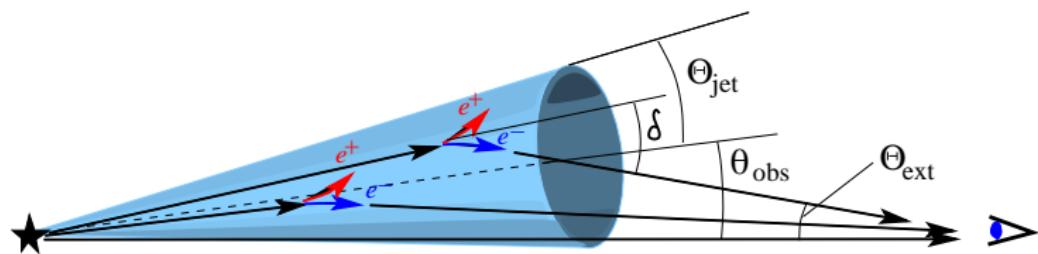
- probability for misalignment $p \propto \vartheta_{\text{obs}}$ \Rightarrow most blazars viewed with $\vartheta_{\text{obs}} \sim \vartheta_{\text{jet}}$

Influence of EGMF on flux from single source: time



- probability for misalignement $p \propto \vartheta_{\text{obs}} \Rightarrow$ most blazars viewed with $\vartheta_{\text{obs}} \sim \vartheta_{\text{jet}}$
- ⇒ halos are **not symmetric**

Influence of EGMF on flux from single source: time

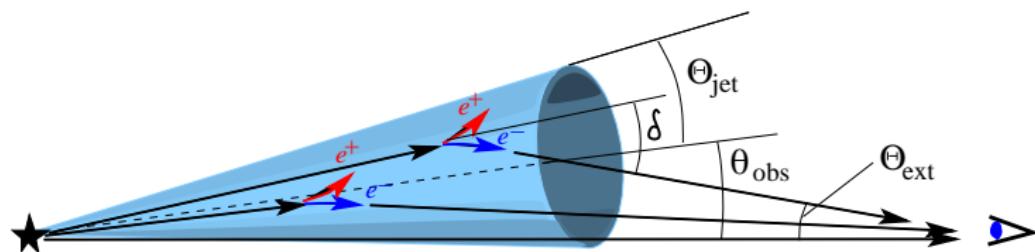


- probability for misalignment $p \propto \vartheta_{\text{obs}} \Rightarrow$ most blazars viewed with $\vartheta_{\text{obs}} \sim \vartheta_{\text{jet}}$
 - ⇒ halos are not symmetric
 - ⇒ **time-delay** is function of ϑ ,

$$T_{\text{delay}}(\vartheta) \sim 3 \times 10^6 \text{ yr} \left[\frac{(\vartheta_{\text{obs}} + \Theta_{\text{jet}})}{5^\circ} \right] \left[\frac{\vartheta}{5^\circ} \right]$$

Observer misaligned with jet:

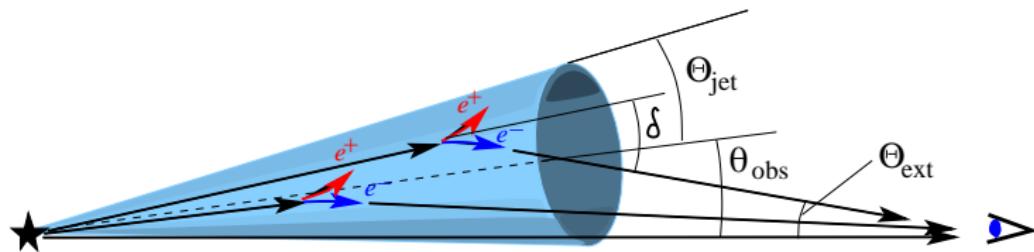
[Neronov et al. '10]



- probability for misalignment $p \propto \vartheta_{\text{obs}} \Rightarrow$ most blazars viewed with $\vartheta_{\text{obs}} \sim \vartheta_{\text{jet}}$

Observer misaligned with jet:

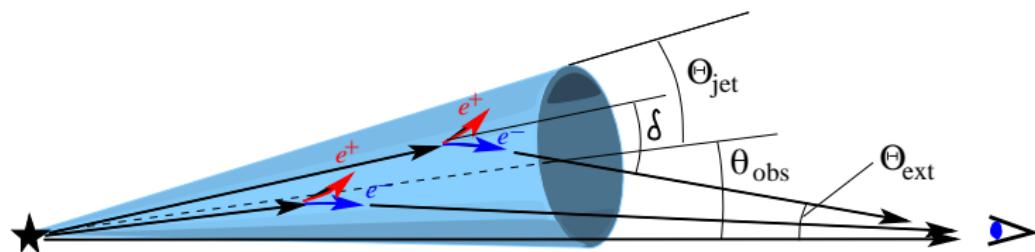
[Neronov et al. '10]



- probability for misalignment $p \propto \vartheta_{\text{obs}} \Rightarrow$ most blazars viewed with $\vartheta_{\text{obs}} \sim \vartheta_{\text{jet}}$
- ⇒ halos are **not symmetric**

Observer misaligned with jet:

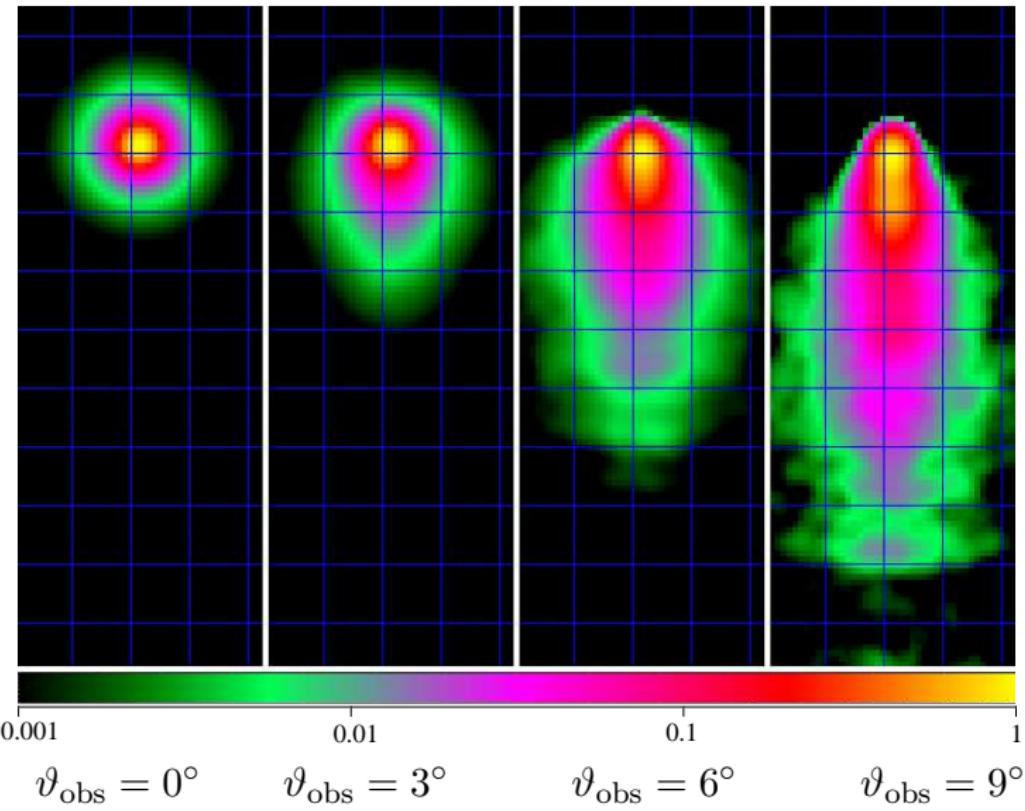
[Neronov et al. '10]



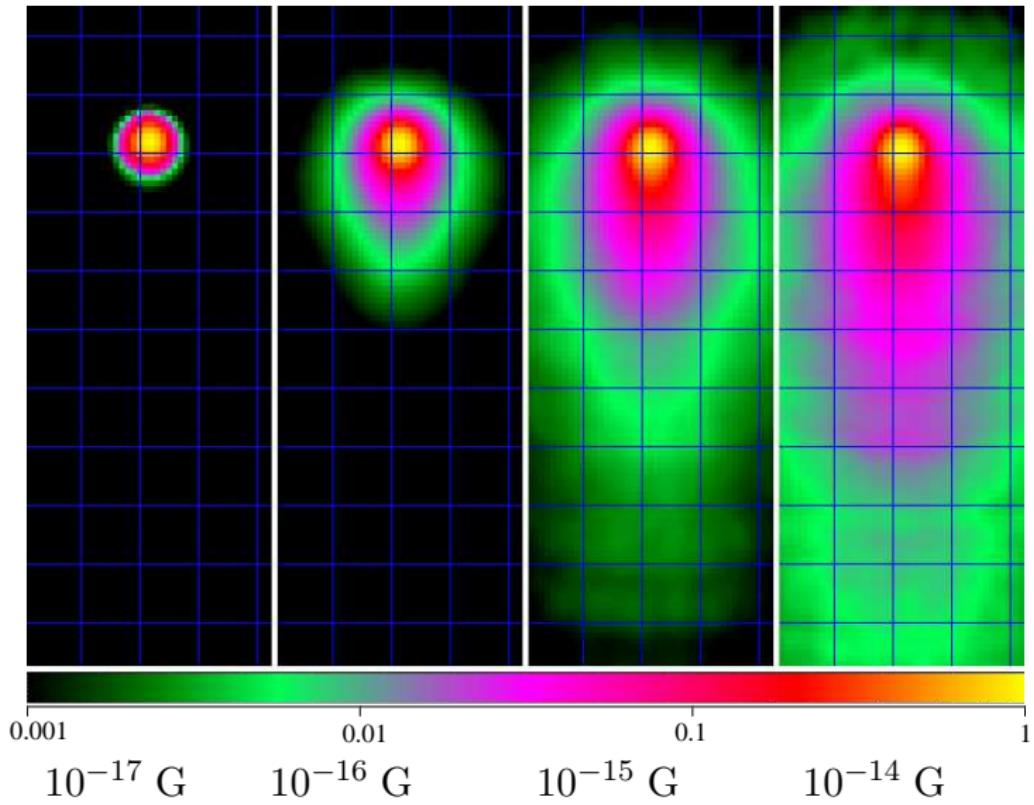
- probability for misalignment $p \propto \vartheta_{\text{obs}} \Rightarrow$ most blazars viewed with $\vartheta_{\text{obs}} \sim \vartheta_{\text{jet}}$
 - ⇒ halos are not symmetric
 - ⇒ time-delay is function of ϑ ,

$$T_{\text{delay}}(\vartheta) \sim 3 \times 10^6 \text{ yr} \left[\frac{(\vartheta_{\text{obs}} + \Theta_{\text{jet}})}{5^\circ} \right] \left[\frac{\vartheta}{5^\circ} \right]$$

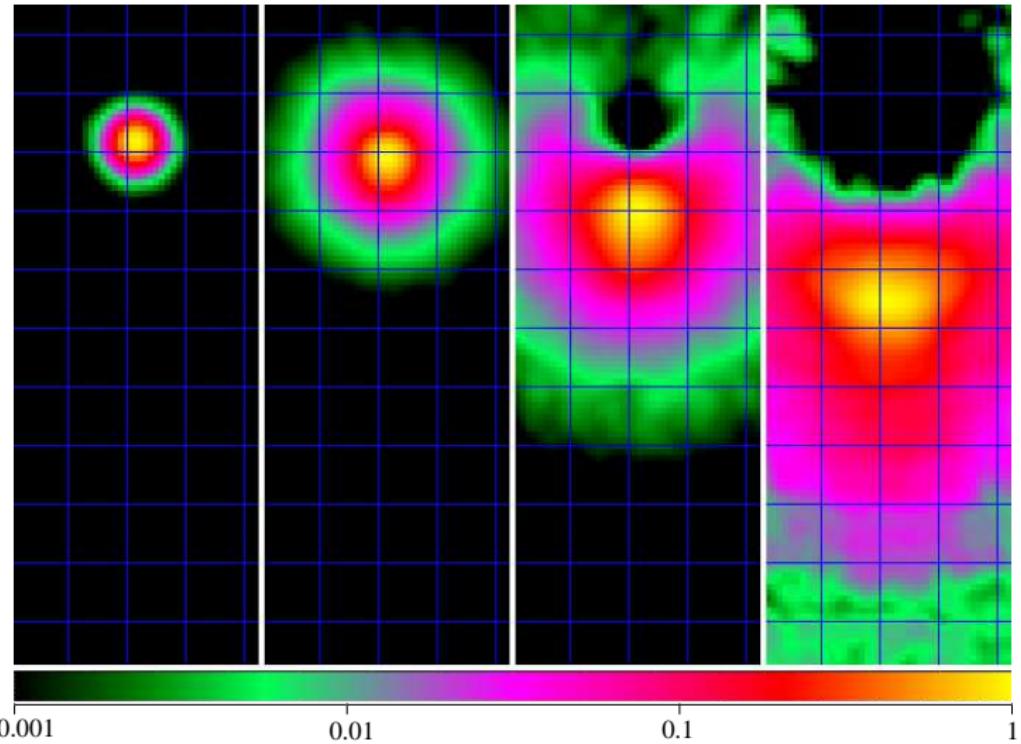
Asymmetric halos around TeV blazars (“GeV jets”):



“GeV jets”: B dependence



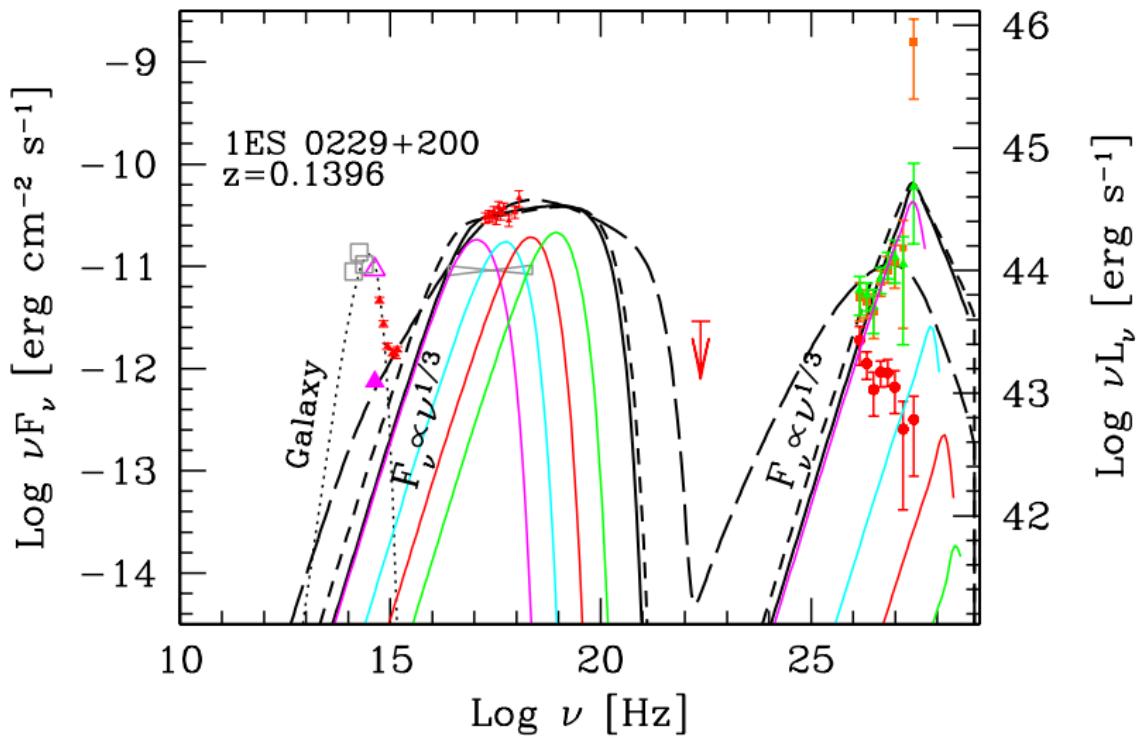
“GeV jets”: time dependence of flares



Lower limit on EGMF:

[A. Neronov, I. Vovk '10, F. Tavecchio et al. '10]

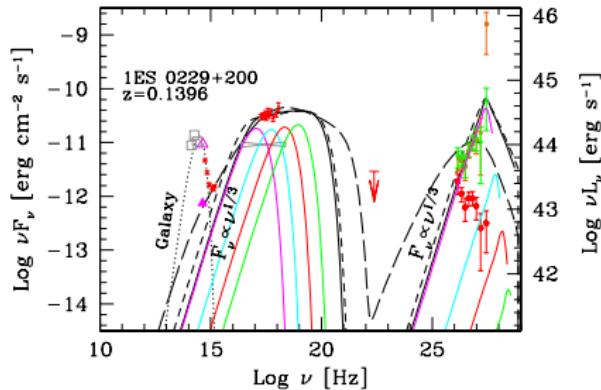
- choose blazar: large z , stationary, low GeV, high multi-TeV emission



Lower limit on EGMF:

[A. Neronov, I. Vovk '10, F. Tavecchio et al. '10]

- choose blazar: large z , stationary, low GeV, high multi-TeV emission

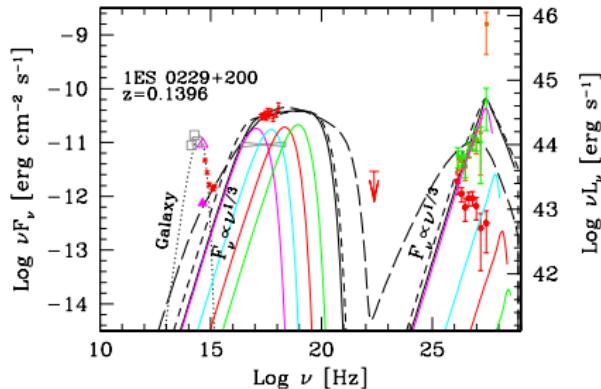


- TeV photons cascade down:
 - small EGMF: fill up GeV range
 - “large” EGMF: deflected outside, isotropized

Lower limit on EGMF:

[A. Neronov, I. Vovk '10, F. Tavecchio et al. '10]

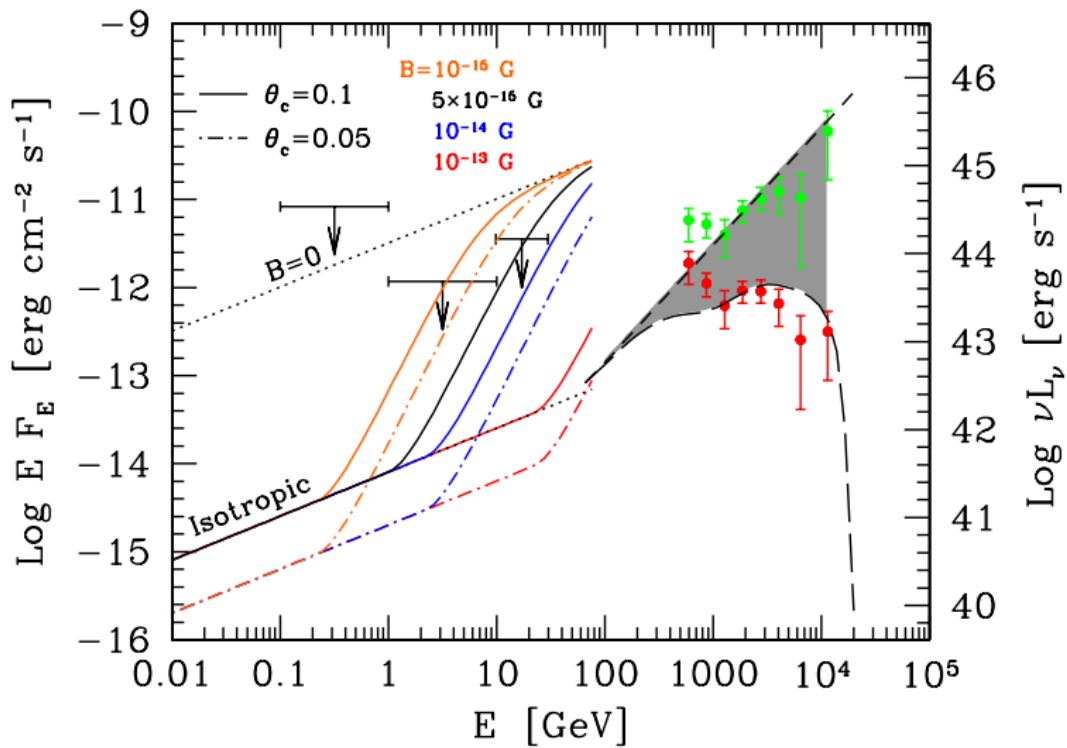
- choose blazar: large z , stationary, low GeV, high multi-TeV emission



- TeV photons cascade down:
 - small EGMF: fill up GeV range
 - “large” EGMF: deflected outside, isotropized
- open questions:
 - influence of **EGMF structure?**
 - time-dependence** for flaring sources?

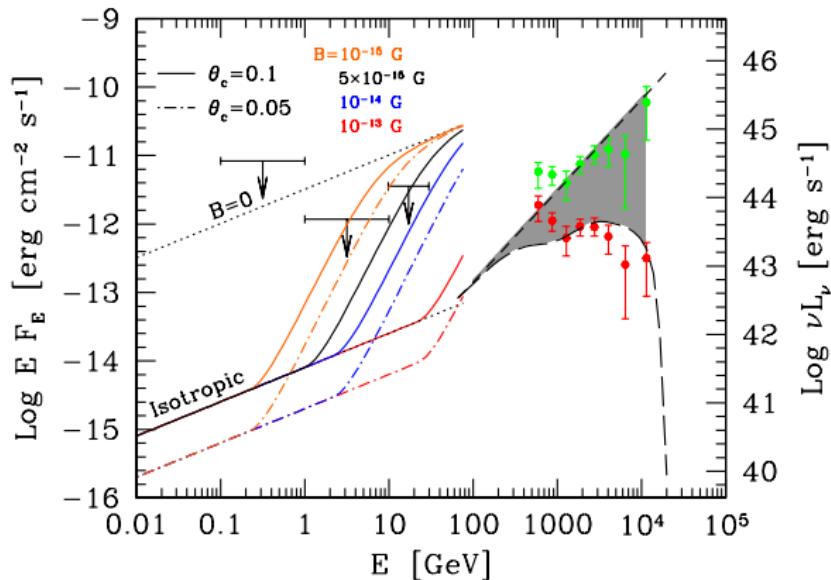
Lower limit on EGMF:

[A. Neronov, I. Vovk '10, F. Tavecchio et al. '10]



Lower limit on EGMF:

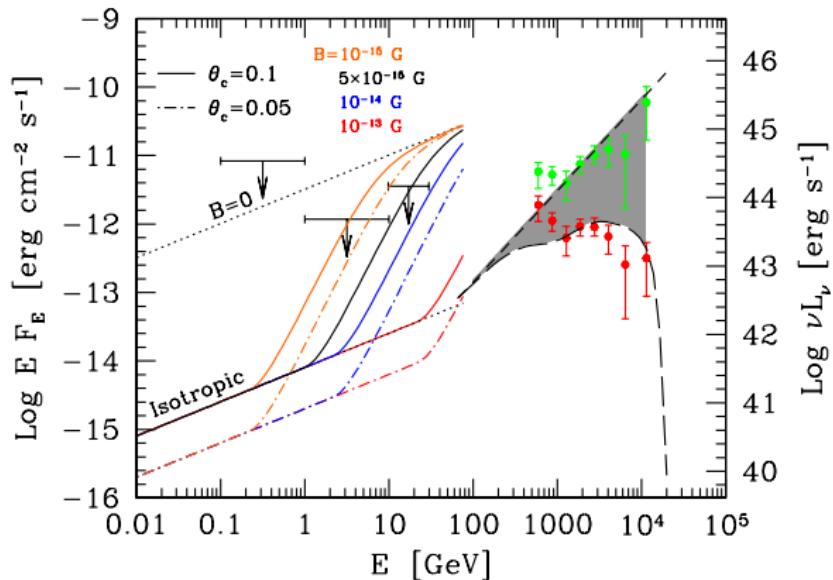
[A. Neronov, I. Vovk '10, F. Tavecchio et al. '10]



- $B \gtrsim 10^{-15} \text{ G}$
- some dependence on v_{jet}
- no simulation of elmag. cascade with B

Lower limit on EGMF:

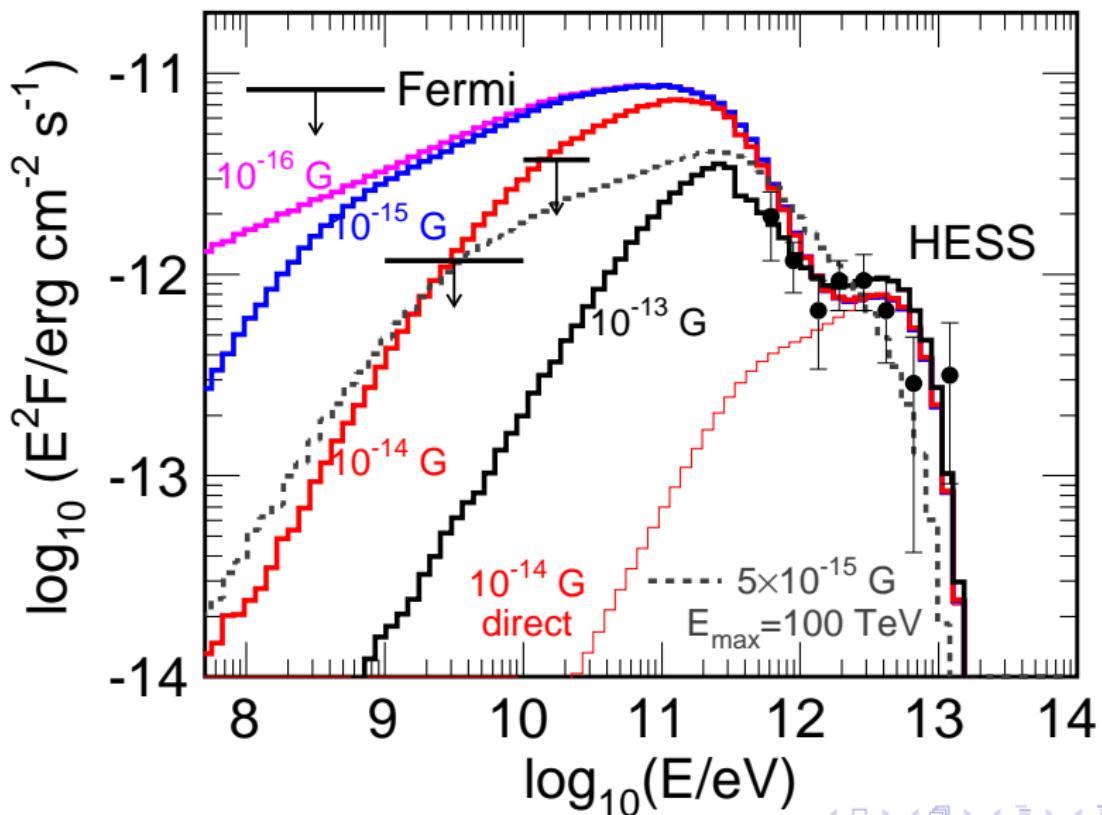
[A. Neronov, I. Vovk '10, F. Tavecchio et al. '10]



- $B \gtrsim 10^{-15} \text{ G}$
- some dependence on ϑ_{jet}
- no simulation of elmag. cascade with B

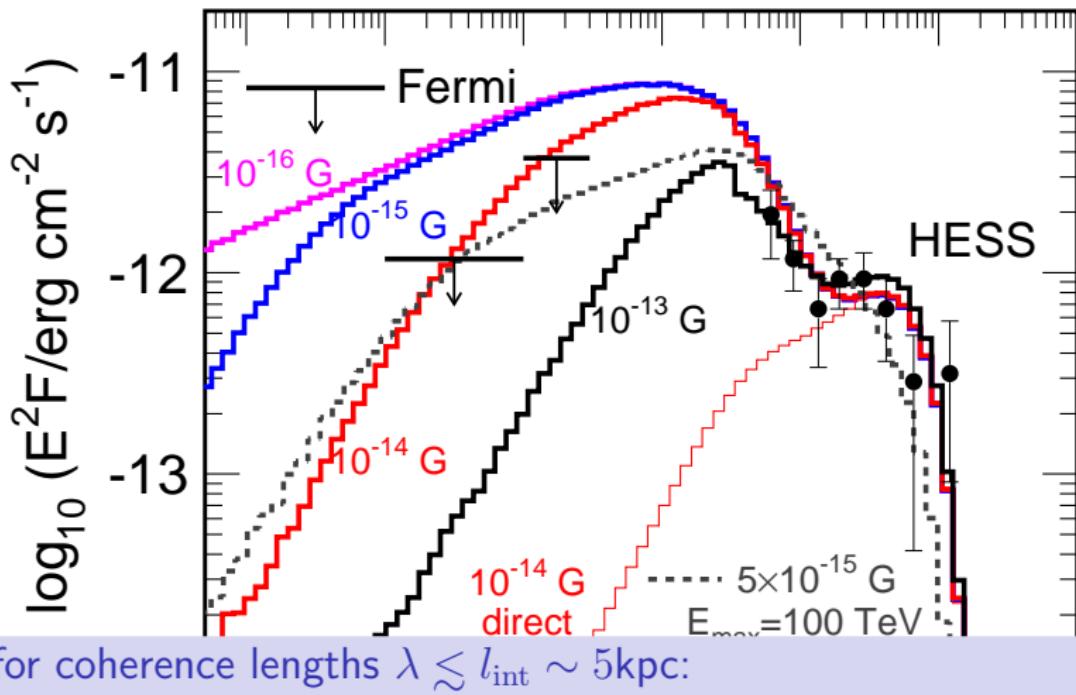
Lower limit on EGMF: uniform field

[Dolag et al. '10]



Lower limit on EGMF: uniform field

[Dolag et al. '10]

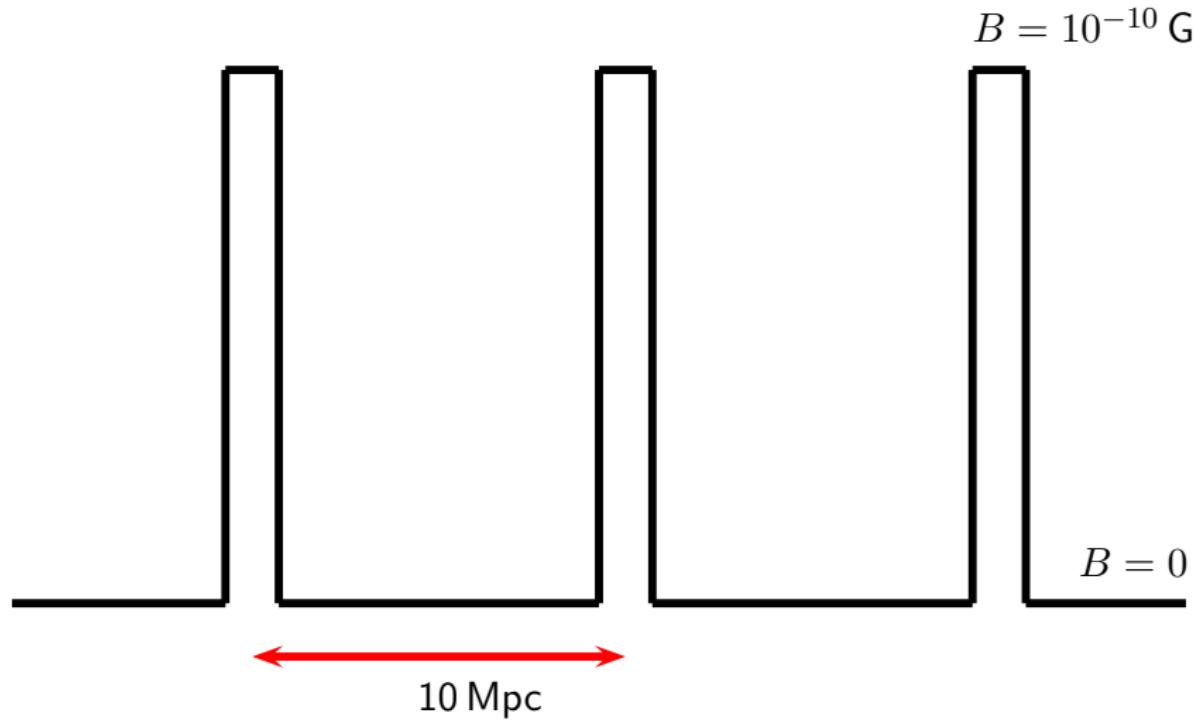


⇒ bound improves as $\lambda^{1/2}$

Lower limit on filling factor:

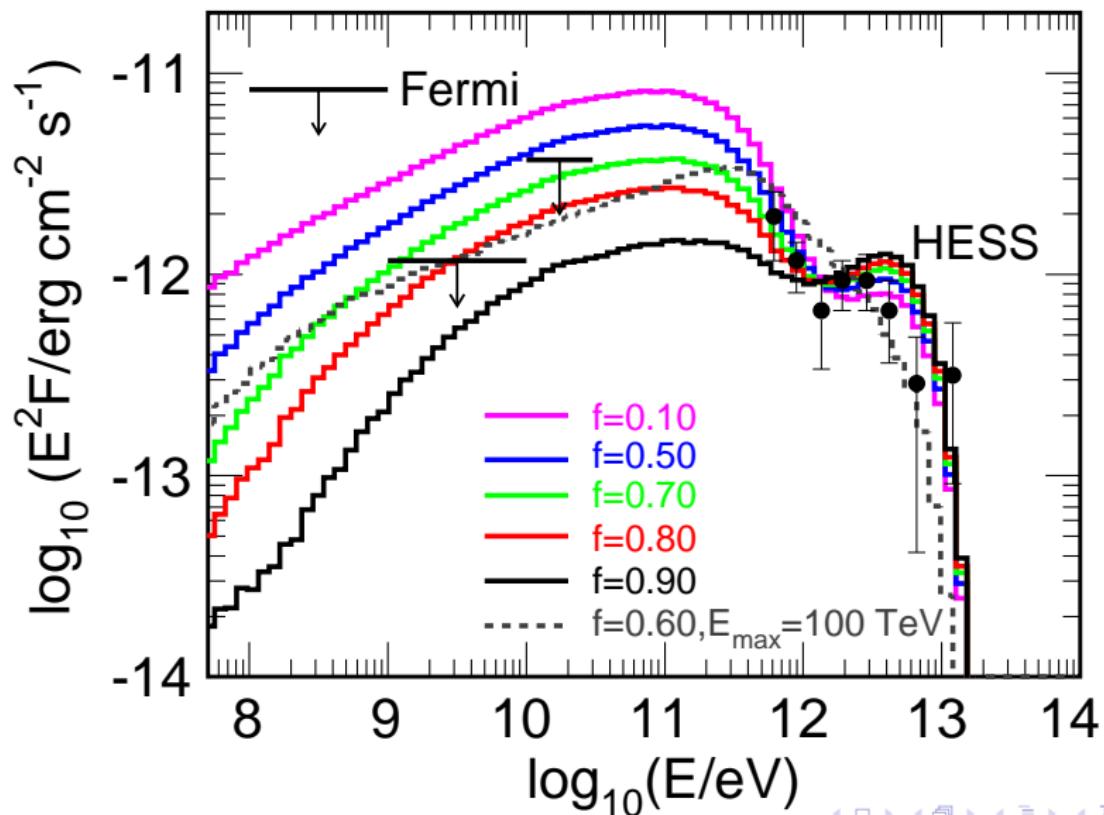
[Dolag et al. '10]

- model filaments by a top-hat:



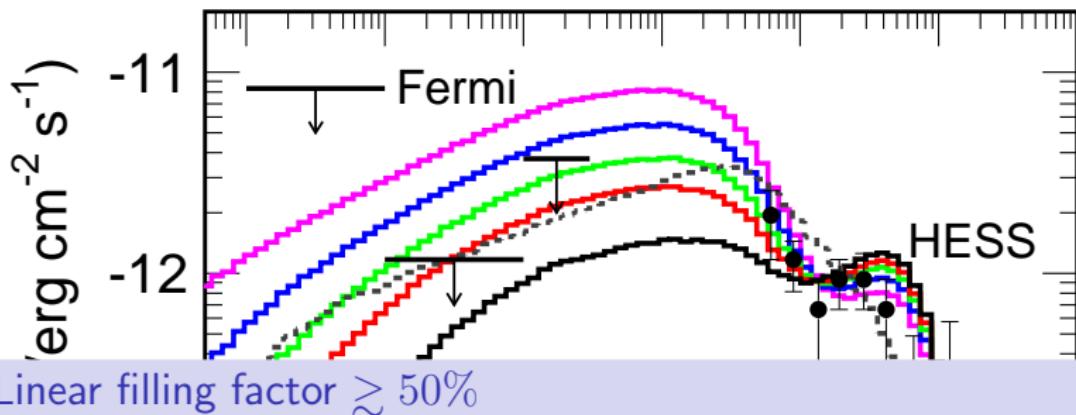
Lower limit on filling factor:

[Dolag et al. '10]



Lower limit on filling factor:

[Dolag et al. '10]

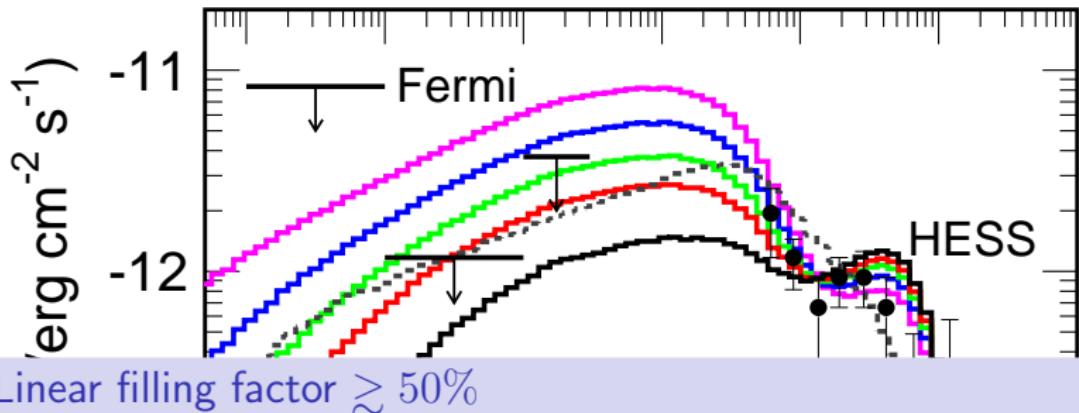


- mainly 3-step cascade: $\gamma \rightarrow e^\pm \rightarrow \gamma$
- photon mean free path $D_\gamma(E) \sim 1000\text{--}50\text{ Mpc}$
- electron mean free path $D_e(E) \sim \text{few kpc}$

$$\log_{10}(E/\text{eV})$$

Lower limit on filling factor:

[Dolag et al. '10]

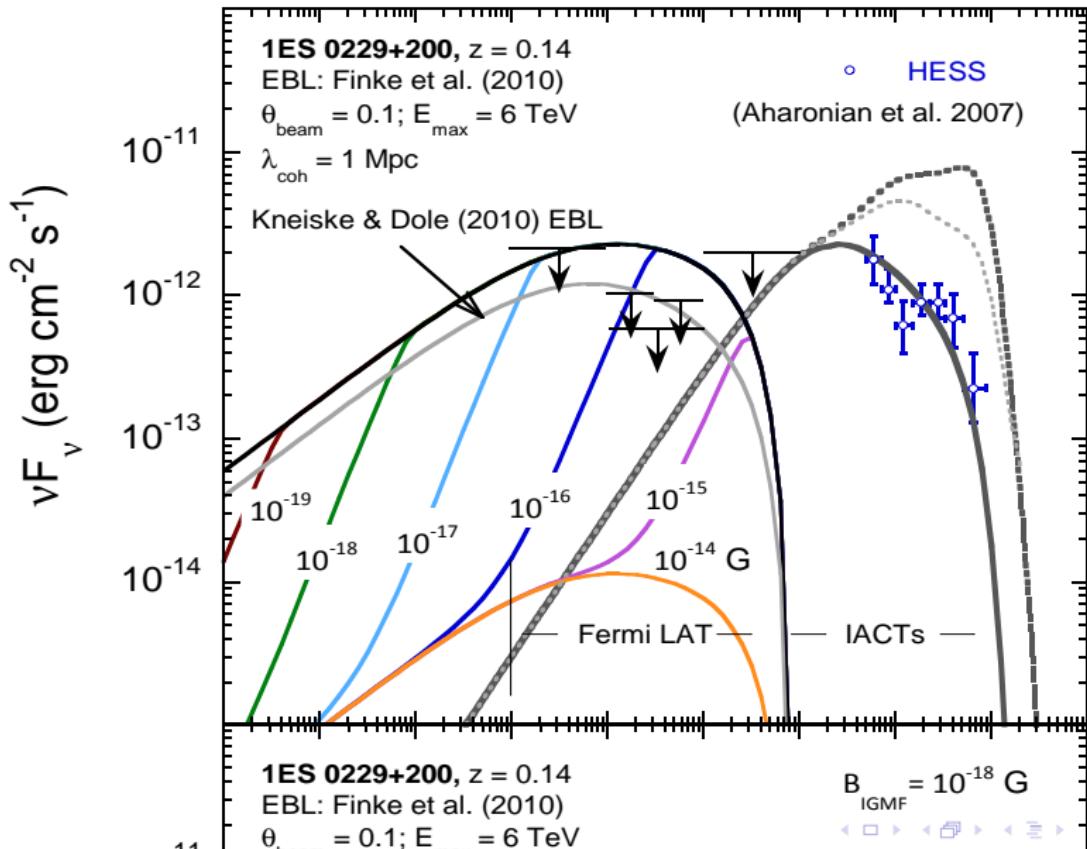


- mainly 3-step cascade: $\gamma \rightarrow e^\pm \rightarrow \gamma$
 - photon mean free path $D_\gamma(E) \sim 1000\text{--}50\text{ Mpc}$
 - electron mean free path $D_e(E) \sim \text{few kpc}$
- ⇒ electrons are created “everywhere” and feel B only close to interaction point

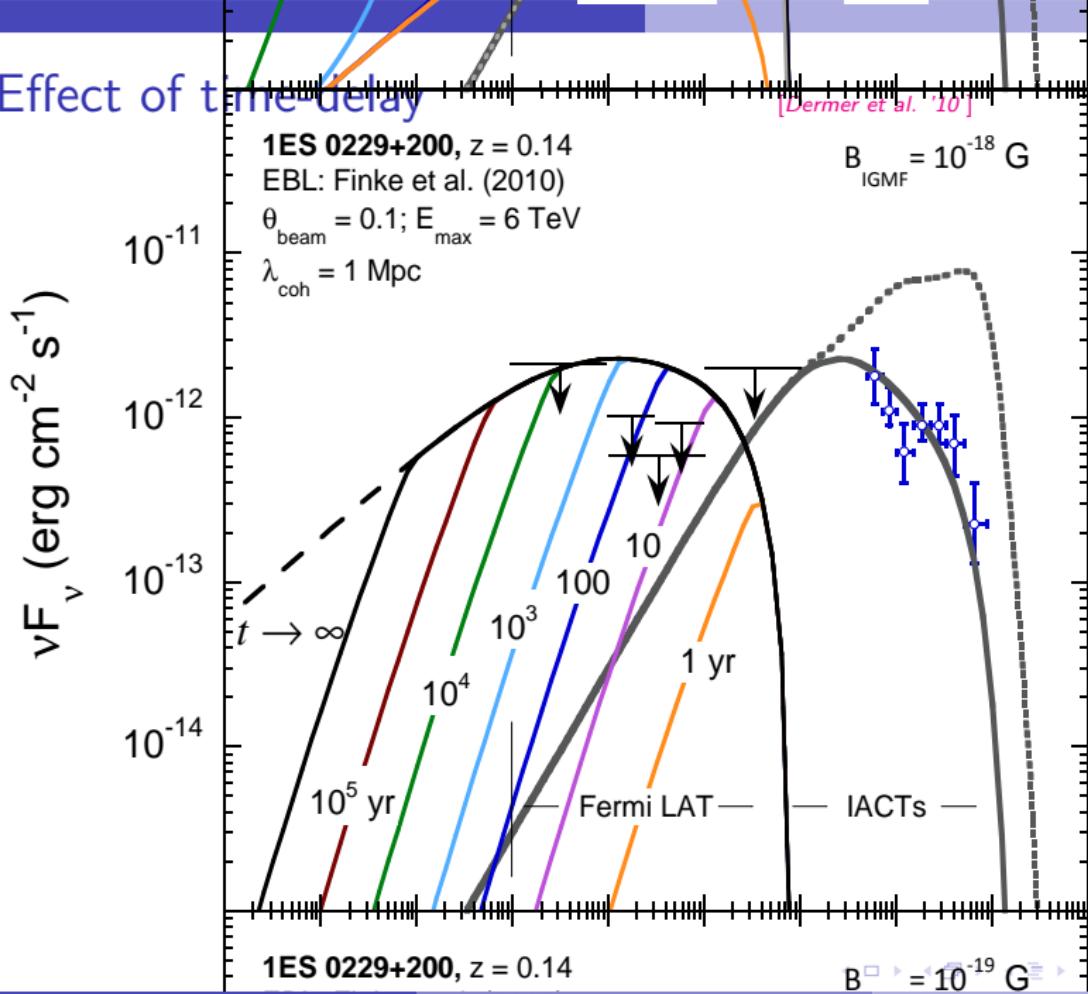
$$\log_{10}(E/\text{eV})$$

Effect of time-delay

[Dermer et al. '10]



Effect of time delay



How to create EGMFs in voids?

- **primordial fields:**
 - ▶ inflation
 - ▶ phase transitions (QCD, electroweak)
 - ▶ reionization

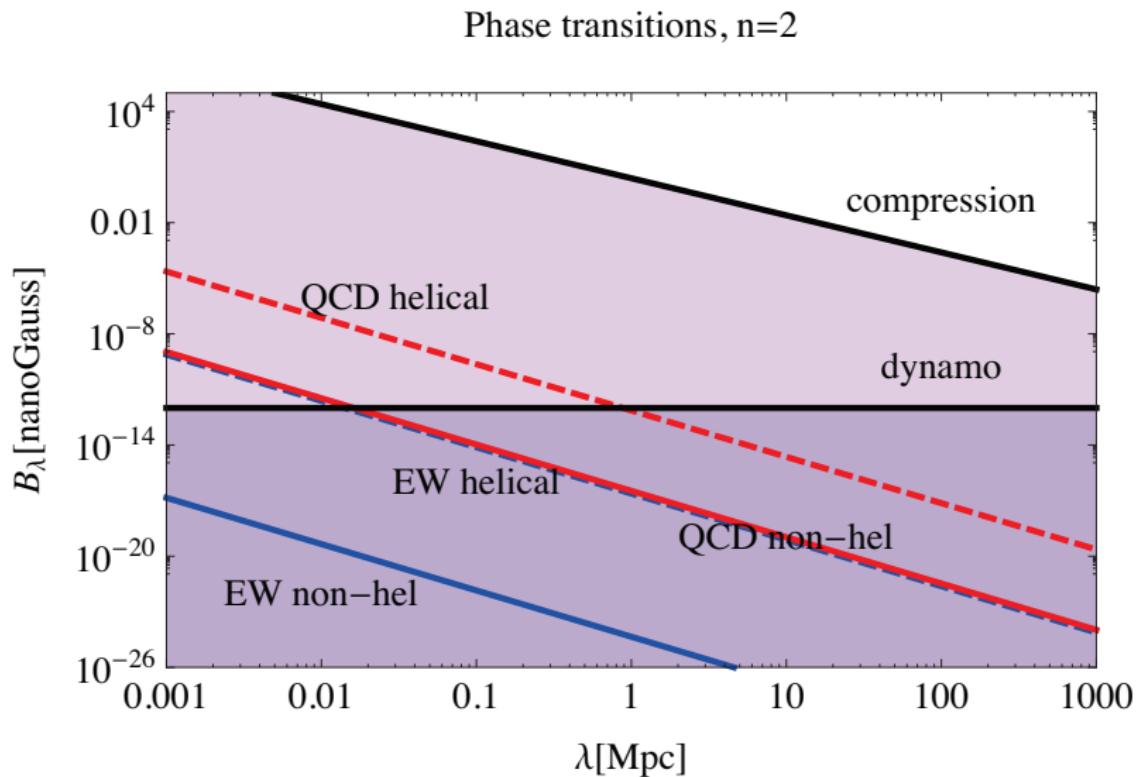
- **astrophysical** (require seed fields):
 - + outflows from AGNs, dwarf galaxies

How to create EGMFs in voids?

- primordial fields:
 - ▶ inflation
 - ▶ phase transitions (QCD, electroweak)
 - ▶ reionization **too weak**
- astrophysical (require seed fields):
 - + outflows from AGNs, dwarf galaxies
 - outflows **collimated**
 - $B > 0$ and $B = 0$ **plasma does not mix**
 - contamination with **heavy elements**

Expectation for “causally” generated fields:

[Caprini, Durer, ...]



Primordial magnetic fields:

- for a Gaussian field

$$\langle B_i(k) B_j^*(k') \rangle = \delta(k - k') \left[(\delta_{ij} - \hat{k}_i \hat{k}_j) S(k) + i \varepsilon_{ijk} k^l H(k) \right]$$

energy density $\rho = 4\pi \int_0^\infty k^2 S(k)$

helicity density $h = 4\pi \int_0^\infty k H(k)$

- characterized by B_λ and coherence length L_c

Primordial magnetic fields:

- for a Gaussian field

$$\langle B_i(k) B_j^*(k') \rangle = \delta(k - k') \left[\left(\delta_{ij} - \hat{k}_i \hat{k}_j \right) S(k) + i \varepsilon_{ijl} k^l H(k) \right]$$

- characterized by B_λ and coherence length L_c
- inflation: “acausal”, $L_c \gg H_0^{-1}$ possible
- phase transitions:
 - ▶ require 1./2.order transition
 - ▶ “causal”, $L_c(t_*) \lesssim H(t_*)^{-1}$

Primordial magnetic fields:

- for a Gaussian field

$$\langle B_i(k)B_j^*(k') \rangle = \delta(k - k') \left[\left(\delta_{ij} - \hat{k}_i \hat{k}_j \right) S(k) + i \varepsilon_{ijl} k^l H(k) \right]$$

- characterized by B_λ and coherence length L_c
 - inflation: “acausal”, $L_c \gg H_0^{-1}$ possible
 - phase transitions:
 - ▶ require 1./2.order transition
 - ▶ “causal”, $L_c(t_*) \lesssim H(t_*)^{-1}$
- $\Rightarrow \langle B_i(x)B_j(y) \rangle$ has compact support $\Rightarrow [\dots]$ is analytic function
- ▶ analyticity & finite $\rho \Rightarrow B_\lambda \sim B_0(L_c/\lambda)^{5/2}$

Primordial magnetic fields:

- for a Gaussian field

$$\langle B_i(k)B_j^*(k') \rangle = \delta(k - k') \left[\left(\delta_{ij} - \hat{k}_i \hat{k}_j \right) S(k) + i \varepsilon_{ijl} k^l H(k) \right]$$

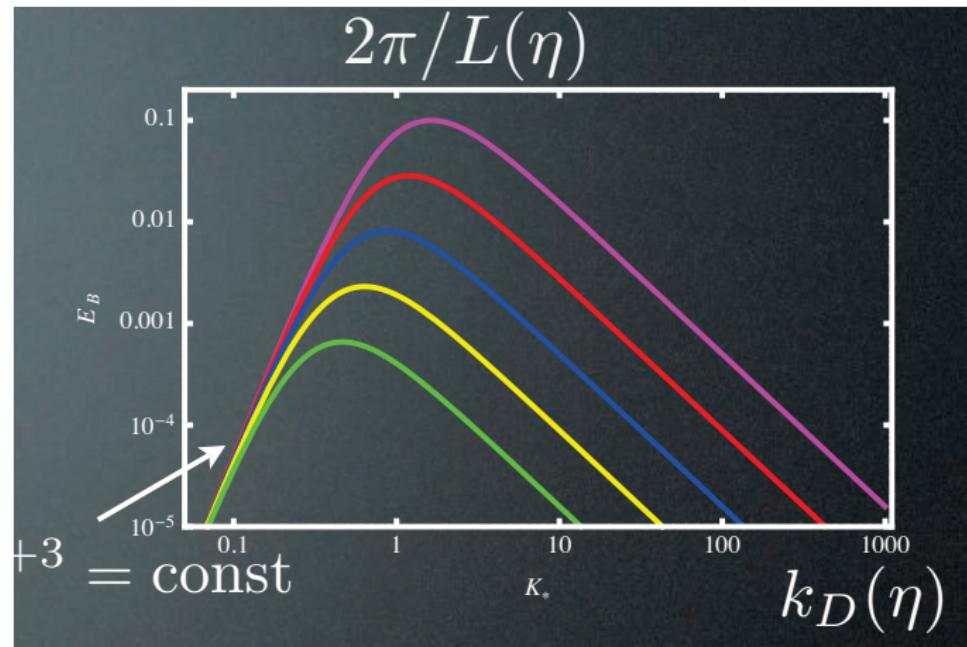
- characterized by B_λ and coherence length L_c
- inflation: “acausal”, $L_c \gg H_0^{-1}$ possible
- phase transitions:
 - ▶ require 1./2.order transition
 - ▶ “causal”, $L_c(t_*) \lesssim H(t_*)^{-1}$

$\Rightarrow \langle B_i(x)B_j(y) \rangle$ has compact support $\Rightarrow [\dots]$ is analytic function

 - ▶ analyticity & finite $\rho \Rightarrow B_\lambda \sim B_0(L_c/\lambda)^{5/2}$
- how are fields evolving after creation?

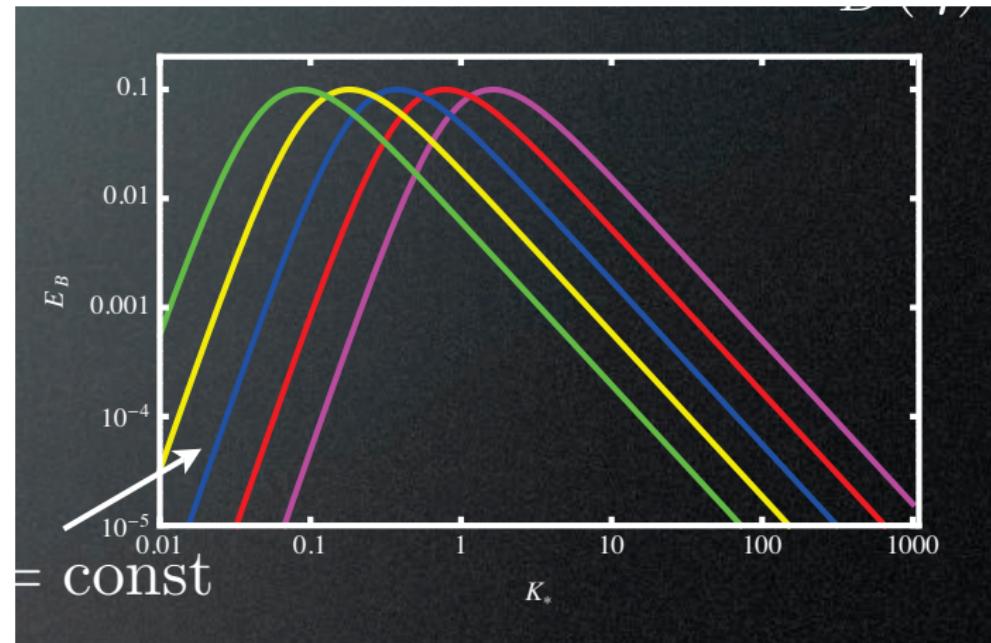
Evolution of primordial magnetic fields:

- non-helical fields: damped above k_D , below $B \propto 1/a^2$



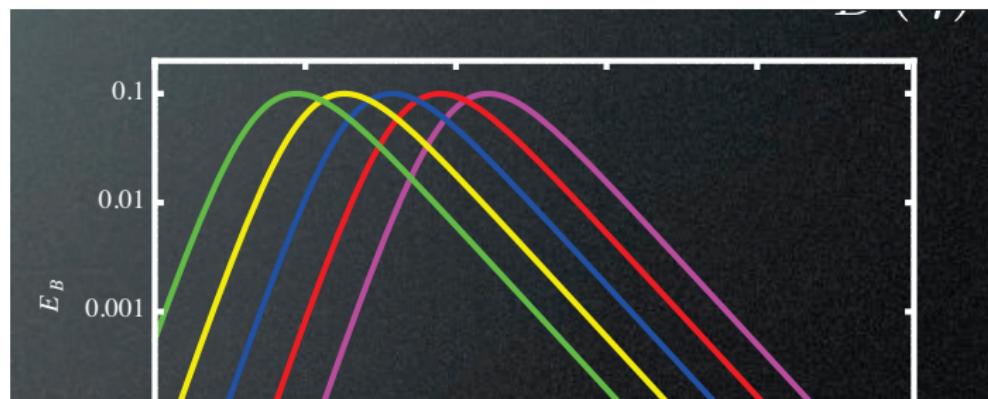
Evolution of primordial magnetic fields:

- helical fields: fluctuations are transferred to larger scales $1/k$



Evolution of primordial magnetic fields:

- helical fields: fluctuations are transferred to larger scales $1/k$



Picture assumes “static” large scale tail $B_\lambda \propto \lambda^{-5/2}$:

- ▶ interaction with turbulent fluid generates $B_\lambda \propto \lambda^{-3/2}$
- ⇒ fields from phase transitions could be strong enough

Summary

- Fermi **non-observation** of TeV blazars requires **EGMF**

⇒ quantitative conclusions:

- ▶ **sure:** large filling factor $f \gtrsim 0.5$
- ▶ **bound on EGMF:** depends on assumed Δt , $B \gtrsim 10^{18} \text{ G}$

- can be improved by more/longer simultaneous observations
 - limit ⇒ detection: CTA?
 - cascade limit from Fermi data reduced by factor ~ 7
- ⇒ km³ neutrino telescope cannot detect cosmogenic neutrinos

Summary

- Fermi non-observation of TeV blazars requires EGMF
- ⇒ quantitative conclusions:
 - ▶ sure: large filling factor $f \gtrsim 0.5$
 - ▶ bound on EGMF: depends on assumed Δt , $B \gtrsim 10^{18} \text{ G}$
- can be improved by more/longer simultaneous observations
- limit ⇒ detection: CTA?
- cascade limit from Fermi data reduced by factor ~ 7
- ⇒ km³ neutrino telescope cannot detect cosmogenic neutrinos