

## Exercise sheet 9

## 1. Energy losses.

a.) Find the time evolution of the energy,  $E(t)$ , of a particle suffering quadratic energy losses,  $-dE/dt = bE^2$ .

b.) Compare the energy density of a magnetic field with  $B = 3\mu\text{G}$  and of the CMB. Determine  $b$ .

c.) Make a  $E(t)$  plot for a 0.1, 1, 10, 100 TeV electron. Assuming a diffusion coefficient  $D(E) = D_0(E/E_0)^{1/3}$  with  $D_0 = 5 \times 10^{26} \text{cm}^2/\text{s}$  and  $E_0 = 10 \text{GeV}$ , what do you conclude if 10 TeV electrons are observed at Earth?

a.) Integrating  $-dE/dt = bE^2$  gives

$$E(t) = \frac{E_0}{1 + bE_0 t} = \frac{E_0}{1 + t/\tau_{1/2}}$$

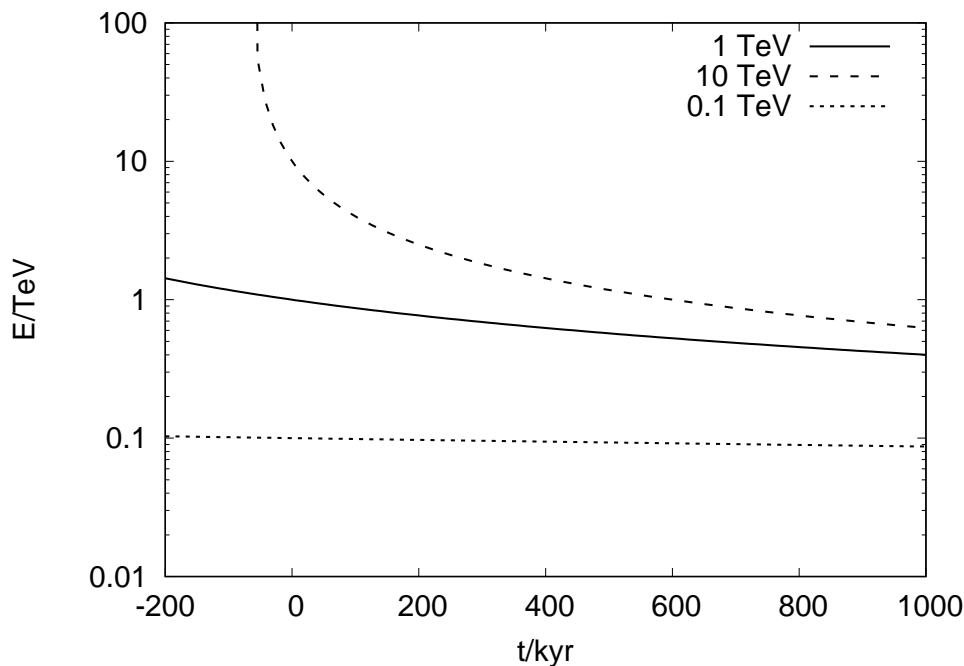
with  $\tau_{1/2} = 1/(bE_0)$ . Looking forward in time,  $\tau_{1/2}$  is the time after which the energy of the particle is halved,  $E(\tau_{1/2}) = E_0/2$ .

b.) The energy density  $u$  of a uniform magnetic field  $B$  is

$$u = \frac{B^2}{8\pi} \simeq 0.45 \frac{\text{eV}}{\text{cm}^3} \left( \frac{B}{3\mu\text{G}} \right)^2 \quad (1)$$

The energy density  $u = aT^4$  of CMB photons with temperature  $T = 2.7 \text{K}$  is  $u \simeq 0.45 \text{eV}/\text{cm}^3$ . Thus

$$b = \frac{3}{4} \sigma_{\text{Th}} \frac{cu}{(m_e c^2)^2} \simeq 5 \times 10^{-17} \frac{1}{\text{GeV} \cdot \text{s}} = 1.5 \times 10^{-3} \frac{1}{\text{TeV} \cdot \text{kyr}}.$$



c.) The interesting point to note is that  $E(t)$  diverges for  $t \rightarrow -\tau_{1/2}$ . Thus we can interpret  $\tau_{1/2}$  also as the earliest creation time of an electron observed with energy  $E_0$  today ( $t = 0$ ). For  $E_0 = 10$  TeV, it is  $\tau_{1/2} \simeq 70$  kyr.

If the particle diffuses (random walk in a turbulent magnetic field), then the distance travelled is  $r = \langle x^2 \rangle^{1/2} = \sqrt{2D\tau_{1/2}}$ . For simplicity, we ignore that  $E = E(t)$  and evaluate  $D(10 \text{ TeV}) \simeq 5 \times 10^{27} \text{ cm}^2/\text{s}$  and thus  $r \simeq 50 \text{ pc}$ . Hence 10 TeV electrons should come from sources in our very local neighbourhood. [The diffusion equation including energy losses can be solved analytically. The result differs not much: the source should be within 100 pc.]

## 2. Charged pion decay.

A charged pion decays mainly via the reaction  $\pi^\pm \rightarrow \mu^\pm + \nu_\mu$ . Discuss this decay analogous to the decay of a neutral pions.

a.) Combining energy conservation  $m_\pi = E_\nu + E_\mu$ , smallness of neutrino masses  $E_\nu = p_\nu$  and the cms condition  $\mathbf{p}_\nu = -\mathbf{p}_\mu$  gives

$$\begin{aligned} (E_\mu - m_\pi)^2 &= |\mathbf{p}_\mu|^2 \\ \underbrace{E_\mu^2 - |\mathbf{p}_\mu|^2}_{=m_\mu^2} - 2E_\mu m_\pi + m_\pi^2 &= 0 \\ m_\mu^2 + m_\pi^2 &= 2E_\mu m_\pi \\ E_\mu &= \frac{m_\pi^2 + m_\mu^2}{2m_\pi} = \frac{m_\pi}{2} \left( 1 + \frac{m_\mu^2}{m_\pi^2} \right) = \frac{m_\pi}{2} (1 + r) \end{aligned}$$

with  $r \equiv m_\mu^2/m_\pi^2 \simeq 0.58$ . The momentum follows as

$$|\mathbf{p}_\mu|^2 = E_\mu^2 - m_\mu^2 = \frac{m_\pi^2}{4} (1 + 2r + r^2) - \frac{4m_\pi^2 m_\mu^2}{4m_\pi^2} = \frac{m_\pi^2}{4} (1 - 2r + r^2)$$

or  $|\mathbf{p}_\mu| = (m_\pi/2)(1 - r)$ .

b.) The energy in the lab system follows from the general Lorentz transformation  $E' = \gamma(E + \beta p \cos \vartheta)$  with  $\vartheta$  as the angle between the velocity  $\beta$  of the pion and the emitted muon. The maximal/minimal values of  $E'$  follow for  $\cos \vartheta = \pm 1$ , i.e. if the muon is emitted parallel and anti-parallel to the direction of flight of the pion.

Inserting  $E = (m_\pi/2)(1 + r)$  and  $p = (m_\pi/2)(1 - r)$  gives

$$E'_{\min} = \frac{\gamma m_\pi}{2} (1 + r \pm \beta(1 - r)).$$

In the ultra-relativistic limit,  $\beta \rightarrow 1$ , and thus  $E_{\max} \simeq \gamma m_\pi = E_\pi$  and  $E_{\min} \simeq r E_\pi \simeq 0.58 E_\pi$ .

c.) In the rest-frame of the pion, the muon is emitted isotropically,  $dN/d\Omega = 1/(4\pi)$ , since there is no preferred direction. In a frame where the pion is moving, we differentiate  $E' = \gamma(E + \beta p \cos \vartheta)$ , obtaining

$$dE' = \gamma \beta p d(\cos \vartheta).$$

Solutions are discussed Monday, 13.11.23

Thus

$$dN = \frac{1}{2} d(\cos \vartheta) = \frac{dE'}{2\gamma\beta p}$$

or

$$\frac{dN}{dE'} = \frac{1}{2\gamma\beta p}.$$

The energy distribution is again a flat box, with the boundaries found in b.)

### 3. Muon decay.

Consider the decay of the muon  $\mu^- \rightarrow e^- + \bar{\nu}_e + \nu_\mu$  (at rest): find the condition that the electron energy is maximal.

Denoting the four-momenta as  $\mu^-(p_1) \rightarrow e^-(p_2) + \bar{\nu}_e(p_3) + \nu_\mu(p_4)$ , it is

$$(p_1 - p_2)^2 = (p_3 + p_4)^2$$

Squaring and solving for  $E_e$ , it follows

$$E_e = \frac{m_\mu^2 + m_e^2 - (p_3 + p_4)^2}{2m_\mu}$$

Thus  $E_e$  is maximal, if

$$(p_3 + p_4)^2 = 2E_3 E_4 (1 - \cos \vartheta)$$

is minimal, or  $\cos \vartheta = 1$ . Hence the neutrino 3-momenta are parallel, opposite to the electron momentum.