

### Exercise sheet 3

#### 1. Transformation between inertial frames.

Consider two inertial frames  $K$  and  $K'$  with parallel axes at  $t = t' = 0$  that are moving with the relative velocity  $v$  in the  $x$  direction.

a.) Show that the linear transformation between the coordinates in  $K$  and  $K'$  is given by

$$\begin{pmatrix} t' \\ x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} At + Bx \\ Dt + Ex \\ y \\ z \end{pmatrix} = \begin{pmatrix} At + Bx \\ A(x - vt) \\ y \\ z \end{pmatrix} \quad (1)$$

b.) Show that requiring the invariance of

$$\Delta s^2 \equiv c^2 t^2 - x^2 - y^2 - z^2 = c^2 t'^2 - x'^2 - y'^2 - z'^2 \quad (2)$$

leads to Lorentz transformations.

c.) What is the condition leading to Galilean transformations?

a.) The origin  $x' = 0$  of  $K'$  correspond in  $K$  to  $x = vt$ . Then

$$0 = Dt + Evt \Rightarrow D = -Ev$$

The origin  $x = 0$  of  $K$  correspond in  $K'$  to  $x' = -vt'$ . Then

$$t' = At \quad (3)$$

$$-vt' = Dt \Rightarrow t' = -\frac{D}{v}t = At \quad (4)$$

and thus  $A = -D/v$  and hence  $A = E$ .

Remark: The transformation has to be linear, if we assume that space and time are uniform.

b.) Square the terms on the RHS of

$$t^2 - x^2 = (At - Bx)^2 - A^2(x - vt)^2,$$

order them as  $t^2(\dots) - x^2(\dots) + 2tx(\dots)$  and compare coefficients to the LHS. This gives  $A = \gamma$  and  $B = -\beta\gamma$ .

c.) Absolute time  $t = t'$  requires  $A = 1$  and  $B = 0$ .

#### 2. Index gymnastics.

Split the tensor  $T^{\mu\nu}$  into its symmetric part  $S^{\mu\nu} = S^{\nu\mu}$  and its anti-symmetric part  $A^{\mu\nu} = -A^{\nu\mu}$ .

a.) Show that this splitting is invariant under Lorentz transformations.

b.) Show that  $S^{\mu\nu} A_{\mu\nu} = 0$ .

Solutions are discussed Thursday, 11.09.19

a.) The antisymmetric tensor components are given by

$$A_{\mu\nu} = \frac{1}{2} (T_{\mu\nu} - T_{\nu\mu}) .$$

The definition is invariant under a Lorentz transformation  $\Lambda^\mu{}_\nu$ , since

$$2\tilde{A}_{\mu\nu} = \tilde{T}_{\mu\nu} - \tilde{T}_{\nu\mu} = \Lambda^\rho{}_\mu \Lambda^\sigma{}_\nu T_{\rho\sigma} - \Lambda^\rho{}_\nu \Lambda^\sigma{}_\mu T_{\rho\sigma} \quad (5)$$

$$= \Lambda^\rho{}_\mu \Lambda^\sigma{}_\nu T_{\rho\sigma} - \Lambda^\sigma{}_\nu \Lambda^\rho{}_\mu T_{\sigma\rho} = \Lambda^\rho{}_\mu \Lambda^\sigma{}_\nu (T_{\rho\sigma} - T_{\sigma\rho}) = \Lambda^\rho{}_\mu \Lambda^\sigma{}_\nu 2A_{\rho\sigma} \quad (6)$$

Here we used first the transformation law for a tensor of rank 2, and exchanged then dummy indices in the second term.

Note that this holds also for a general coordinate transformation,  $\Lambda^\mu{}_\nu \rightarrow \frac{\partial \tilde{x}^\nu}{\partial x^\mu}$ .

b.) To see that the contraction of a symmetric tensor  $S_{\mu\nu}$  with an antisymmetric tensor  $A_{\mu\nu}$  gives zero, consider

$$S_{\mu\nu} A^{\mu\nu} = -S_{\mu\nu} A^{\nu\mu} = -S_{\nu\mu} A^{\nu\mu} = -S_{\mu\nu} A^{\mu\nu} . \quad (7)$$

Here we used first the antisymmetry of  $A^{\mu\nu}$ , then the symmetry of  $S_{\mu\nu}$ , and finally exchanged the dummy summation indices. Note that this remains true, if the tensor expression contains additional indices.

### 3. Two particle decay.

A particle at rest decays into two particles,  $A \rightarrow B + C$ .

a.) Find the energy and momenta of  $B$  and  $C$ .

b.) What happens for  $m_A < m_B + m_C$ ?

a.) Squaring  $p_c^\mu = p_a^\mu - p_b^\mu$  and using  $\mathbf{p}_a = 0$  gives

$$m_c^2 = m_a^2 - 2m_a E_b + m_b^2 \quad (8)$$

Solving for the energies of the decay products,

$$E_b = \frac{m_a^2 + m_b^2 - m_c^2}{2m_a} \quad \text{and} \quad E_c = \frac{m_a^2 + m_c^2 - m_b^2}{2m_a} . \quad (9)$$

Then the momenta follow as

$$\mathbf{p}_b^2 = E_b^2 - m_b^2 = \left( \frac{m_a^2 + m_b^2 - m_c^2}{2m_a} \right)^2 - \frac{4m_a^2 m_b^2}{4m_a^2} \quad (10)$$

or

$$|\mathbf{p}_b| = |\mathbf{p}_c| = \frac{1}{2m_a} [m_a^4 + m_b^4 + m_c^4 - 2m_a m_b - 2m_a m_c - 2m_b m_c]^{1/2} \quad (11)$$

$$\equiv \frac{1}{2m_a} \lambda^{1/2}(m_a^2, m_b^2, m_c^2) . \quad (12)$$

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The Kibble or triangle function  $\lambda(x, y, z)$  satisfies

$$\lambda(x, y, z) = x^2 + y^2 + z^2 - 2xy - 2yz - 2xz \quad (13a)$$

$$= [x^2 - (y + z)^2] [x^2 - (y - z)^2] . \quad (13b)$$

b.) The decay is kinematically forbidden, since the momenta  $|\mathbf{p}_b|$  and  $|\mathbf{p}_c|$  become imaginary, cf. with Eq. (13b).

#### 4. Threshold energy.

Particle A hits particle B at rest, producing the particles  $C_1, \dots, C_n$ .

a.) Calculate the threshold (i.e. the minimal energy of A) for this reaction.

b.) What is the threshold energy for  $pp \rightarrow X\pi^0$ , where  $X$  is the minimal set of particles such that the reaction is allowed?

c.) What is the threshold energy for  $pp \rightarrow X\bar{p}$ ?

a.) Denote the total momentum by  $P^\mu$ . The initial state in the lab frame has  $P^\mu = (E_a + m_b, \mathbf{p}_a)$  and thus

$$P^2 = 2E_a m_b + m_a^2 + m_b^2 .$$

After the collision, in the CMS frame at threshold

$$P^2 = \left( \sum_i m_i, \mathbf{0} \right)^2 \equiv (M, \mathbf{0})^2 = M^2$$

Thus

$$E_a = \frac{M^2 - m_a^2 - m_b^2}{2m_b} .$$

b.) Since baryon number is conserved,  $X = \{pp\}$ , and thus  $M = 2m_p + m_\pi \simeq 2011 \text{ MeV}$  or  $E_{\text{th}} \simeq 1.2 \text{ GeV}$ .

c.) Since baryon number is conserved,  $X = \{ppp\bar{p}\}$ , and thus  $M^2 = 16m_p^2$  or  $E_{\text{th}} = 7m_p \simeq 6.5 \text{ GeV}$ .