

KLASSISK MEKANIK

Løsning Føring 2

a) L' og L ekvivalente dersom

$$\frac{d}{dt} \frac{\partial}{\partial \dot{q}^i} \frac{d}{dt} F(q, t) - \frac{\partial}{\partial q^i} \frac{d}{dt} F(q, t) = 0. \quad (1)$$

Deriverer: $\frac{dF}{dt} = \frac{\partial F}{\partial t} + \frac{\partial F}{\partial q^i} \dot{q}^i$

Venstre side av (1): $\frac{d}{dt} \frac{\partial}{\partial \dot{q}^i} \left(\frac{\partial F}{\partial t} + \frac{\partial F}{\partial q^j} \dot{q}^j \right) - \frac{\partial}{\partial q^i} \left(\frac{\partial F}{\partial t} + \frac{\partial F}{\partial q^j} \dot{q}^j \right) =$

$$= \frac{d}{dt} \frac{\partial F}{\partial \dot{q}^i} - \frac{\partial^2 F}{\partial q^i \partial t} - \frac{\partial^2 F}{\partial q^i \partial q^j} \dot{q}^j =$$

$$= \left(\frac{\partial^2 F}{\partial q^i \partial t} + \frac{\partial^2 F}{\partial q^i \partial q^j} \dot{q}^j \right) - \frac{\partial^2 F}{\partial q^i \partial t} - \frac{\partial^2 F}{\partial q^i \partial q^j} \dot{q}^j = 0, \text{ som stemmer.}$$

b) $[\nabla \times (\nabla \times \vec{A})]_i = \epsilon_{ijk} \partial_j (\nabla \times \vec{A})_k = \epsilon_{ijk} \partial_j \epsilon_{klm} \partial_l A_m$
 $= \epsilon_{kij} \epsilon_{klm} \partial_j \partial_l A_m = (\delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}) \partial_j \partial_l A_m =$
 $= \partial_i \partial_l A_l - \partial_j \partial_j A_i = \partial_i (\nabla \cdot \vec{A}) - \nabla^2 A_i, \Rightarrow$
 $\nabla \times (\nabla \times \vec{A}) = \nabla (\nabla \cdot \vec{A}) - \nabla^2 \vec{A}$

$$[\nabla \times (\nabla \times \vec{V})]_i = \epsilon_{ijk} V_j (\nabla \times \vec{V})_k = \epsilon_{ijk} V_j \epsilon_{klm} \partial_l V_m$$

 $= \epsilon_{kij} \epsilon_{klm} V_j \partial_l V_m = (\delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}) V_j \partial_l V_m =$
 $= V_j \partial_i V_j - V_j \partial_j V_i = \frac{1}{2} \partial_i V^2 - (\vec{V} \cdot \nabla) V_i \Rightarrow$
 $\nabla \times (\nabla \times \vec{V}) = \frac{1}{2} \nabla V^2 - (\vec{V} \cdot \nabla) \vec{V}$

$$\nabla \cdot (\vec{A} \times \vec{B}) = \partial_i \epsilon_{ijk} A_j B_k = \epsilon_{ijk} (A_{j,i} B_k + A_j B_{k,i})$$

$$= B_k \epsilon_{kij} A_{j,i} + A_j \epsilon_{jki} B_{k,i}$$

$$= B_k (\nabla \times \vec{A})_k - A_j \underbrace{\epsilon_{jki} B_{k,i}}_{(\nabla \times \vec{B})_j} = \vec{B} \cdot (\nabla \times \vec{A}) - \vec{A} \cdot (\nabla \times \vec{B}).$$

$$[\text{Her er } \frac{\partial A_i}{\partial x_j} \equiv \partial_j A_i \equiv A_{i,j}]$$

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Øving 2, forts.

c) Friksjonskraft $F_f = -\partial\mathcal{F}/\partial v$.

Systemets arbeid mot friksjonen, per tidsenhet, er

$$\dot{W}_b = -F_f \cdot v = \frac{\partial\mathcal{F}}{\partial v} \cdot v. \quad \text{Da } \mathcal{F} = C \cdot v^2 \text{ blir}$$

$$\underline{\dot{W}_b} = 2Cv^2 = \underline{2\mathcal{F}}$$

Lagrange: $\frac{d}{dt} \frac{\partial L}{\partial v} - \frac{\partial L}{\partial x} + \frac{\partial \mathcal{F}}{\partial v} = 0$ gir, med

$$L = T - V = \frac{1}{2}mv^2 - \frac{1}{2}kx^2 \text{ og } \mathcal{F} = 3\pi\mu a v^2,$$

$$m\ddot{x} + kx + 6\pi\mu a v = 0$$

∴

$$\underline{\ddot{x} + 2\lambda\dot{x} + \omega_0^2 x = 0}, \text{ med } \lambda = \frac{3\pi\mu a}{m}, \omega_0 = \sqrt{\frac{k}{m}}$$

Løsning er $x(t) = x_0 e^{-\lambda t} \cos \omega_0 t$ når $\lambda/\omega_0 \ll 1$.

(Derivasjon $\dot{x}(t) \approx -\omega_0 x_0 e^{-\lambda t} \sin \omega_0 t = 0$ når $t=0$.)

Midler energitap $\overline{\dot{W}_b}$ over én periode $2\pi/\omega_0$:

$$\overline{\dot{W}_b} = 2\overline{\mathcal{F}}$$

Da $\mathcal{F} = 3\pi\mu a v^2 = m\lambda v^2$ blir $\overline{\dot{W}_b} = 2m\lambda \overline{v^2}$.

Her er $\overline{v^2} = \overline{(\omega_0 x_0)^2 e^{-2\lambda t} \sin^2 \omega_0 t} \approx \frac{1}{2}(\omega_0 x_0)^2 e^{-2\lambda t}$.

$$\underline{\overline{\dot{W}_b}} = 2m\lambda \cdot \frac{1}{2}(\omega_0 x_0)^2 e^{-2\lambda t} = \underline{m\lambda(\omega_0 x_0)^2 e^{-2\lambda t}}$$