

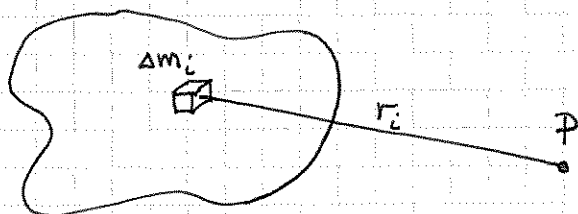
# Potensial og felt fra massefordelinger [YF 13.6, LL 11.2]

Potensial fra liten masse  $\Delta m$ :



$$\Delta V(r) = - \frac{G \cdot \Delta m}{r}$$

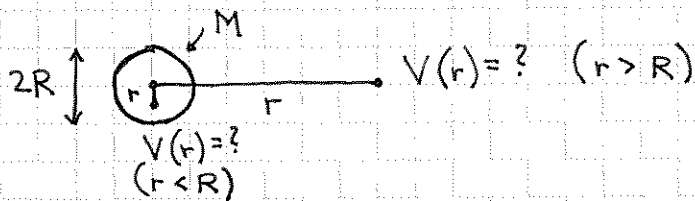
Potensial fra massefordeling:



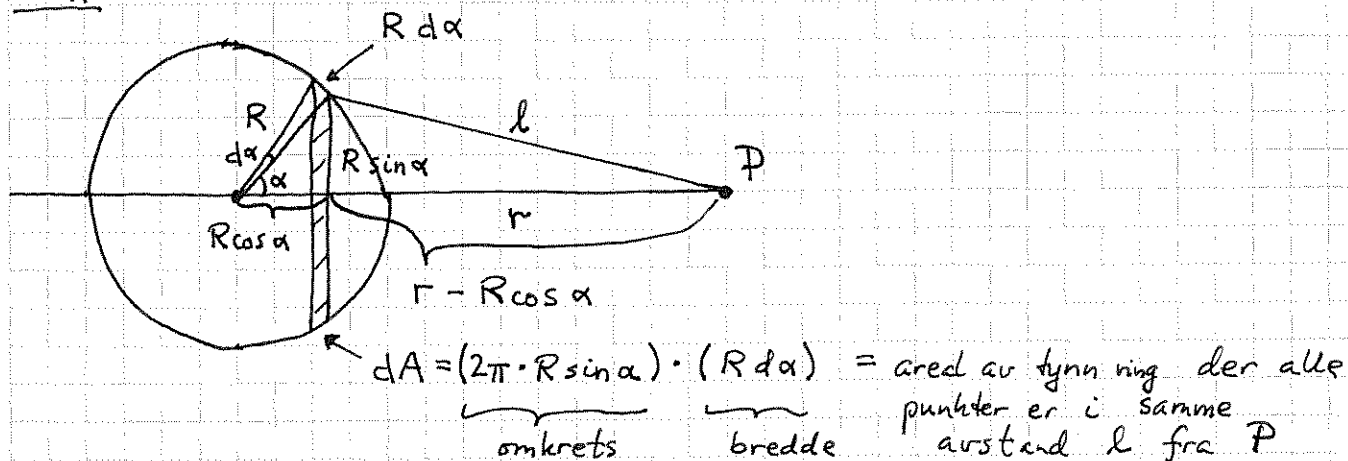
$V_P$  = potensial i punkt P

$$V_P = \sum_i V_P^i = -G \sum_i \frac{\Delta m_i}{r_i} \xrightarrow{\Delta m_i \rightarrow 0} -G \int \frac{dm}{r}$$

Eks 1: Kuleskall



$r > R$ :



$$\begin{aligned} \text{Pythagoras} \Rightarrow l^2 &= (R \sin \alpha)^2 + (r - R \cos \alpha)^2 \\ &= R^2 \sin^2 \alpha + r^2 - 2rR \cos \alpha + R^2 \cos^2 \alpha \\ &= R^2 + r^2 - 2rR \cos \alpha \end{aligned}$$

Bidrag til  $V$  i punktet  $P$  fra stykke ring:

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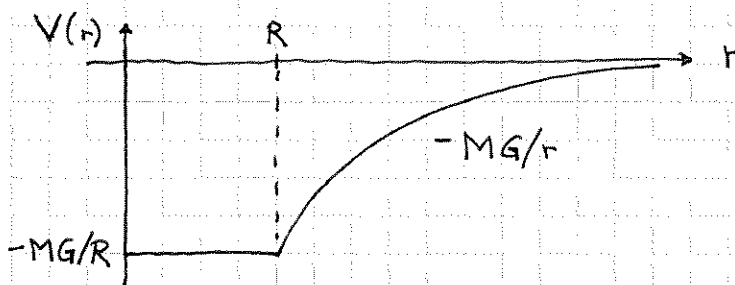
$$dV = -G \frac{dm}{l} ; \quad dm/M = dA/A = 2\pi R \sin\alpha \cdot R d\alpha / 4\pi R^2 \\ = \frac{1}{2} \sin\alpha d\alpha$$

Totalt potensial i  $P$ :

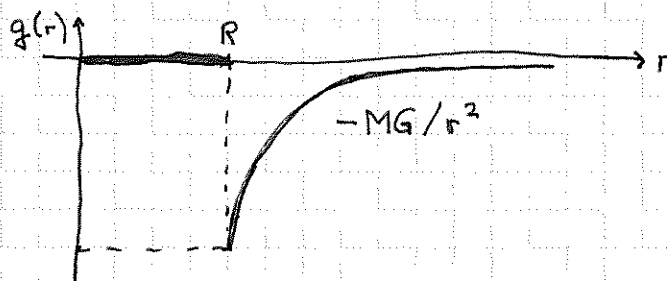
$$V(r) = \int dV = -G \int \frac{dm}{l} = -\frac{MG}{2} \int_0^\pi \frac{\sin\alpha d\alpha}{\sqrt{R^2+r^2-2rR\cos\alpha}} \\ = -\frac{MG}{2} \int_0^\pi \frac{\sqrt{R^2+r^2-2rR\cos\alpha}}{rR} \\ = -\frac{MG}{2rR} \left\{ \sqrt{R^2+r^2+2rR} - \sqrt{R^2+r^2-2rR} \right\} \\ = -\frac{MG}{2rR} \left\{ (R+r) - (r-R) \right\} = -\frac{MG}{r}$$

$r < R$ :

Alt blir som for  $r > R$ , bortsett fra at  $\sqrt{R^2+r^2-2rR} = R-r$ , slik at  $V(r) = \dots = -\frac{MG}{2rR} \left\{ (R+r) - (R-r) \right\} = -\frac{MG}{R}$

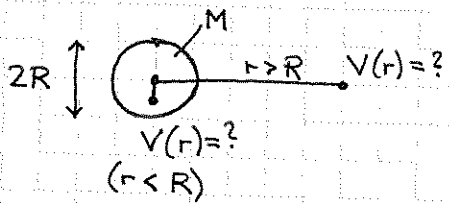


$$\text{Feltet: } \vec{g}(r) = -\nabla V(r) = -\hat{r} \frac{dV}{dr} = \begin{cases} 0 & ; r < R \\ -\hat{r} \frac{MG}{r^2} & ; r > R \end{cases}$$



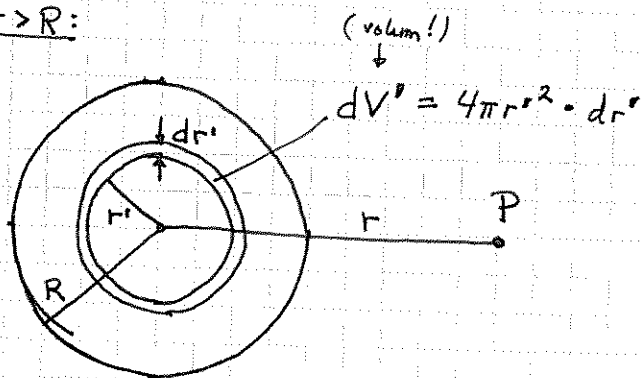
## Eks 2: Kompakt kule

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Strategi: Del opp kule i tynne kuleskall, bruk resultatet fra Eks 1, og integrer.

$r > R$ :



$$\begin{aligned} \frac{dm'}{M} &= \frac{dV'}{V} = \frac{4\pi r'^2 dr'}{\frac{4}{3}\pi R^3} \\ &= \frac{3}{R^3} r'^2 dr' \end{aligned}$$

Bidrag fra tynnt kuleskall til potensialet i P:

$$dV = -G \frac{dm'}{r} = -\frac{3MG}{rR^3} r'^2 dr'$$

Totalt potensial i P:

$$V(r) = \int dV = -\frac{3MG}{rR^3} \int_0^R r'^2 dr' = -\frac{3MG}{rR^3} \cdot \frac{1}{3} R^3 = -\frac{MG}{r}$$

$r < R$ :

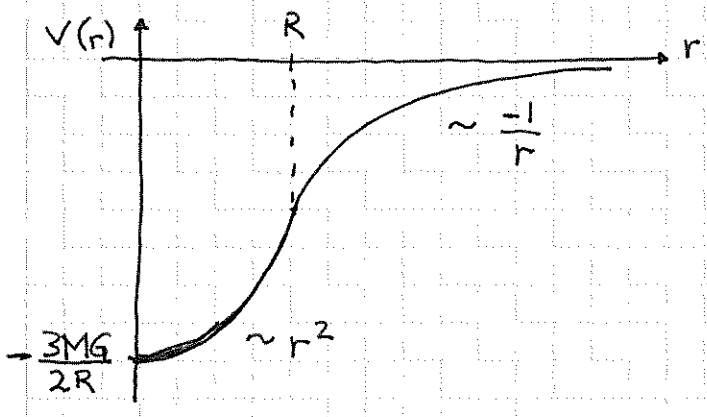
For tynne kuleskall med radius  $r' < r$ :  $dV = -\frac{G \cdot dm'}{r}$

————— " —————  $r' > r$ :  $dV = -\frac{G \cdot dm'}{r'}$

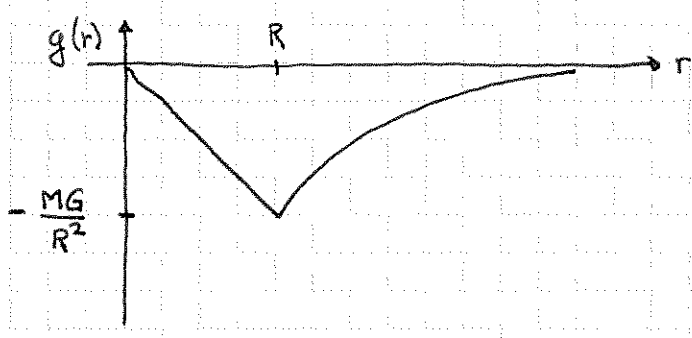
$$\Rightarrow V(r) = -\frac{3MG}{R^3} \left\{ \int_0^r \frac{r'^2 dr'}{r} + \int_r^R \frac{r'^2 dr'}{r'} \right\}$$

$$= -\frac{3MG}{R^3} \left\{ \frac{1}{3} r^2 + \frac{1}{2} R^2 - \frac{1}{2} r^2 \right\}$$

$$= -\frac{3MG}{2R} \left\{ 1 - \frac{r^2}{3R^2} \right\}$$



Feltet:  $\vec{g}(r) = -\nabla V(r) = -\hat{r} \frac{dV}{dr} = \begin{cases} -\hat{r} \frac{MGr}{R^3} & ; r < R \\ -\hat{r} \frac{MG}{r^2} & ; r > R \end{cases}$



⇒ Kraft på masse m ved  $r = R$ :

$$\vec{F}(r=R) = m \cdot \vec{g}(R) = -\hat{r} MmG/R^2,$$

som om hele massen M var plassert i sentrum av kula!

[ Denne antagelsen har vi da også gjort "hele tiden" for tyngdekraften på en masse m på jordoverflaten,  $\vec{F} = m\vec{g}$  med  $g = 9.81 \text{ m/s}^2$ . ]