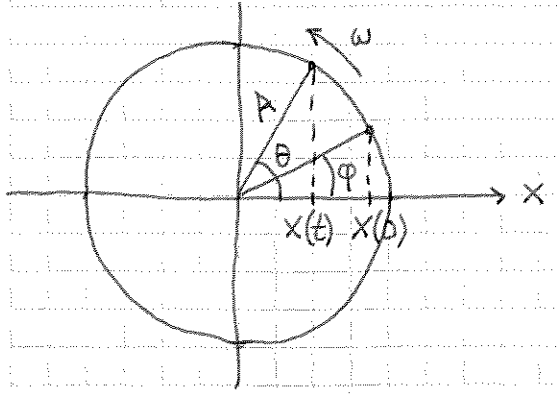


Analogi mellom 1D harm. osc. og sirkelbevegelse:



$$\theta(t) = \varphi + \omega t = \text{fasen}$$

A = amplituden

$$x(t) = A \cos \theta(t) = A \cos(\omega t + \varphi)$$

Energi, harm. osc. [YF 14.3, LL 9.4]

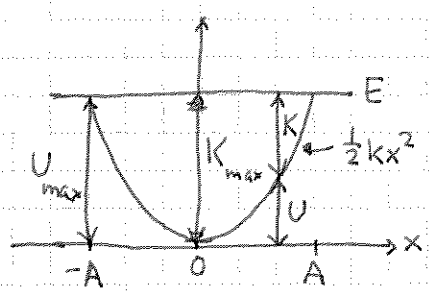
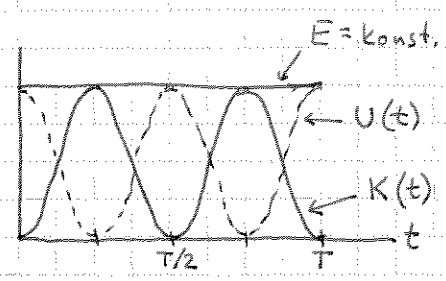
$$K(t) = \frac{1}{2} m \dot{x}(t)^2 = \frac{1}{2} m \omega^2 A^2 \sin^2(\omega t + \varphi) = \frac{1}{2} k A^2 \sin^2(\omega t + \varphi)$$

$$U(t) = \frac{1}{2} k x(t)^2 = \frac{1}{2} k A^2 \cos^2(\omega t + \varphi)$$

$$[U = - \int_0^x \vec{F} \cdot d\vec{l} = - \int_0^x (-kx) dx = \frac{1}{2} k x^2; \text{ se s. 34}]$$

Total energi: $E = K + U = \frac{1}{2} k A^2 = \frac{1}{2} m \omega^2 A^2 = \text{konst.}$

⇒ vi har konservativt system, (mekanisk) energi E er bevart (se s 27 - 34)



$$\vec{F} = -\nabla U$$

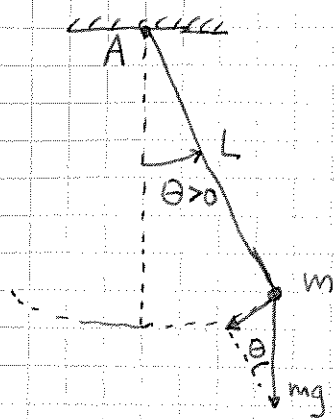
$$= -\hat{x} \partial U / \partial x$$

$$= -kx \hat{x}; \partial x$$

Harm. osc., oppsummering:

- F prop. med utsving fra likevekt
- U — " — kvadratet av utsvinget
- Bevegelsesligning: $\ddot{X} + \omega^2 x = 0$
- "Utsving" kan være lengde, vinkel, temperatur, trykk, ...

Eks 1: Matematisk pendel [VF 14.5, LL 9.6]



N2 || sirkelbuen: $-mg \sin \theta = m\ddot{s} = mL\ddot{\theta}$

Antar små utsving, $\theta \ll 1 \Rightarrow \sin \theta \approx \theta$

$\Rightarrow \ddot{\theta} + (g/L)\theta = 0$

\Rightarrow harm. osc. med $\omega = \sqrt{g/L}$

$\Rightarrow \theta(t) = \theta_0 \cos(\omega t + \varphi)$

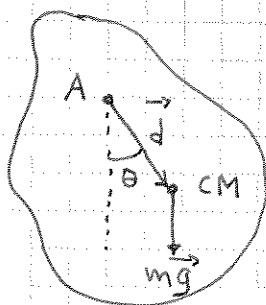
[se også øving 6, oppg. 3: numerisk løsning for vilkårlig θ]

N2 for rotasjon om A gir samme ligning:

$\tau = I\ddot{\theta}$, $I = mL^2$, $\vec{\tau} = \vec{L} \times m\vec{g}$, $\tau = -Lmg \sin \theta \approx -Lmg\theta$

[Demo: $L \sim 1.5\text{m} \Rightarrow T \sim 2.5\text{s}$, OK: $T = 2\pi\sqrt{L/g} \approx 2.5\text{s}$]

Eks 2: Fysisk pendel [VF 14.6, LL 9.6]



stivt legeme, masse m, svinger om akse gjennom A, treghetsmoment I mhp denne aksen

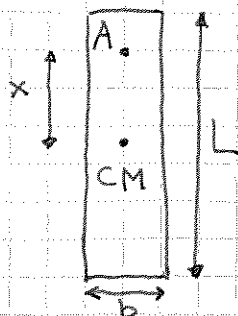
$\tau = I\ddot{\theta}$, $\tau = -d \cdot mg \cdot \sin \theta \approx -d \cdot mg \cdot \theta$ ($\theta \ll 1$)

$\Rightarrow \ddot{\theta} + \frac{mgd}{I} \theta = 0$

\Rightarrow harm. osc. med $\omega = \sqrt{mgd/I}$

Kontroll: punktmasse m i CM $\Rightarrow I = md^2 \Rightarrow \omega = \sqrt{g/d}$, OK.

Eks 3: Trefjøl (øving 12) [Demo]



$x =$ avst. fra CM til rotasjonsaksen A

$I(x) = m(L^2 + b^2)/12 + mx^2$

$T(x) = 2\pi/\omega(x) = \dots$ (se øving 12)

Dempede svingninger [YF 14.7, LL 9.7]

(90)

Friksjon \Rightarrow fri svingninger dempes

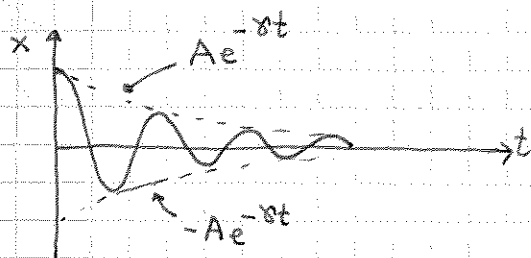
- 1) Langsom beregelse i fluid: $\vec{f} = -b \dot{\hat{x}}$
 - 2) Rask " " " " : $\vec{f} = -D \dot{\hat{x}}^2 \hat{x}$
 - 3) Torr friksjon: $\vec{f} = -\mu_k N \hat{x}$
- } se s. 18-20

Antar tilfelle 1) $\Rightarrow -kx - b\dot{x} = m\ddot{x}$

$$\Rightarrow \ddot{x} + 2\gamma \dot{x} + \omega_0^2 x = 0 \quad \left(\gamma = \frac{b}{2m}, \omega_0^2 = \frac{k}{m} \right)$$

Svak ("underkritisk") damping, $\gamma < \omega_0$:

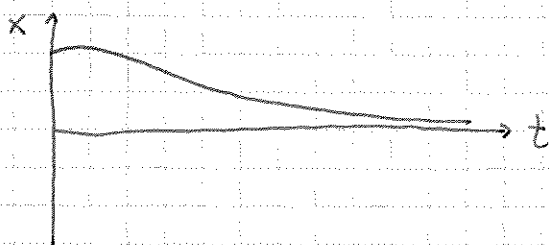
$$x(t) = A e^{-\gamma t} \sin(\omega t + \varphi); \quad \omega = \sqrt{\omega_0^2 - \gamma^2}$$



- redusert frekvens pga damping, $\omega < \omega_0$
- maks. utsving antar eksponentielt, $A e^{-\gamma t}$

Sterk ("overkritisk") damping, $\gamma > \omega_0$:

$$x(t) = A e^{-\alpha_1 t} + B e^{-\alpha_2 t}; \quad \alpha_{1,2} = \gamma \pm \sqrt{\gamma^2 - \omega_0^2}$$



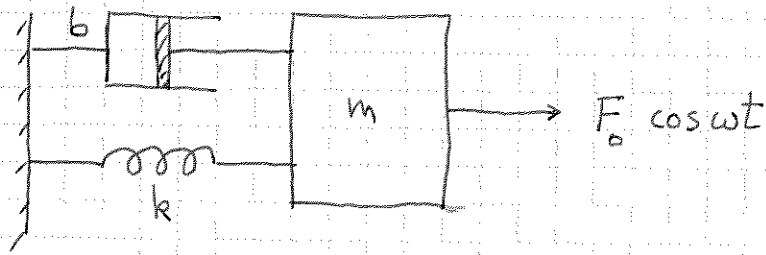
- ingen svingninger

Kritisk damping: $\gamma = \omega_0$ (Eks: støtdempere)

Ubestemte koef. fastlegges med gitte initialbetingelser.

Tvingne svingninger og resonans [YF 14.8, LL 9.9]

(91)



Anta harmonisk ytre kraft $F_y(t) = F_0 \cos \omega t$

$$\stackrel{N2}{\Rightarrow} -kx - b\dot{x} + F_0 \cos \omega t = m\ddot{x}$$

Anta at $F_y(t)$ har virket så lenge at m svinger med samme ω som F_y (f.eks. $t \gg 1/\delta$ med underkritisk damping).

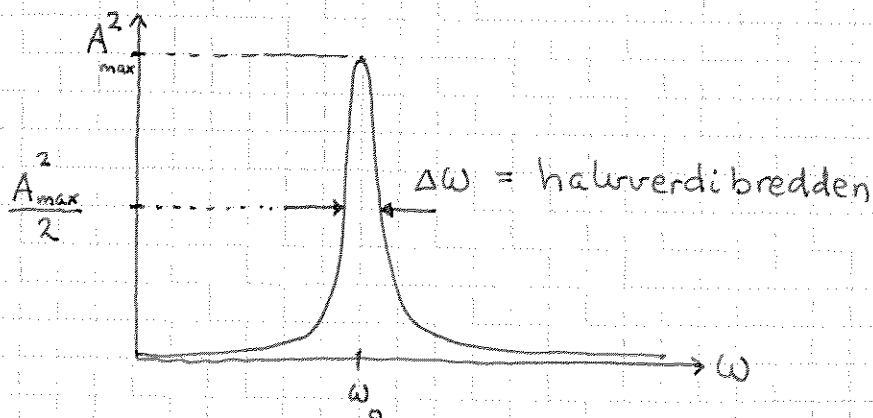
Da er løsningen:

$$x(t) = A(\omega) \sin(\omega t + \varphi) \quad [\text{der } \varphi = \varphi(\omega); \text{ se Bølgefysikk}]$$

med frekvensavhengig amplitude:

$$A(\omega) = \frac{F_0/m}{\sqrt{(\omega^2 - \omega_0^2)^2 + (2\delta\omega)^2}} \quad \left(\omega_0^2 = \frac{k}{m}, \quad 2\delta = \frac{b}{m}\right)$$

Resonans: Hvis δ er liten og $\omega \approx \omega_0$, blir A stor



Liten δ (svak damping) \Rightarrow skarp resonans, med $\Delta\omega = 2\delta$

$Q = \omega_0 / \Delta\omega = \omega_0 / 2\delta =$ resonansens "kvalitet" (Q-verdi)

HLE
13.11.12