

SOLUTION ASSIGNMENT 12

Question 1

Angular momentum:

$$\mathbf{L} = \mathbf{r} \times \mathbf{p}$$

Consider first $[p_i, L_j]$ and start with $i = j$, e.g., $i = 1$ (see p 147 in notater97.pdf):

$$\begin{aligned} [p_x, L_x] &= [p_x, yp_z - zp_y] \\ &= [p_x, y]p_z + [p_x, p_z]y - [p_x, z]p_y - [p_x, p_y]z \end{aligned}$$

We have

$$[p_i, x_j] = -\delta_{ij} \quad , \quad [p_i, p_j] = 0$$

and therefore

$$[p_x, L_x] = 0$$

Correspondingly:

$$[p_y, L_y] = [p_z, L_z] = 0$$

Next, consider $[p_i, L_j]$ with $i \neq j$, e.g., $i = 1$ og $j = 2$:

$$\begin{aligned} [p_x, L_y] &= [p_x, zp_x - xp_z] \\ &= [p_x, z]p_x + [p_x, p_x]z - [p_x, x]p_z - [p_x, p_z]x \\ &= 0 + 0 - (-1)p_z - 0 \\ &= p_z \end{aligned}$$

Cyclic change of x, y, z then gives

$$[p_y, L_z] = p_x \quad , \quad [p_z, L_y] = p_x$$

Interchanging indices only gives a change of sign, e.g.,

$$[p_z, L_y] = -p_x$$

In total,

$$[p_i, L_j] = \varepsilon_{ijk} p_k$$

With $L_j = \varepsilon_{jkl} x_k p_l$ this is obtained more directly:

$$\begin{aligned} [p_i, L_j] &= [p_i, \varepsilon_{jkl} x_k p_l] \\ &= [p_i, x_k] \varepsilon_{jkl} p_l \\ &= -\delta_{ik} \varepsilon_{jkl} p_l \\ &= -\varepsilon_{jil} p_l \\ &= \varepsilon_{ijl} p_l \end{aligned}$$

And now I guess we may handle $[x_i, L_j]$ the same way:

$$\begin{aligned} [x_i, L_j] &= [x_i, \varepsilon_{jkl} x_k p_l] \\ &= [x_i, p_l] \varepsilon_{jkl} x_k \\ &= \delta_{il} \varepsilon_{jkl} x_k \\ &= \varepsilon_{jki} x_k \\ &= \varepsilon_{ijk} x_k \end{aligned}$$

Question 2

Hand written solution in Norwegian on next page.

$$1) \quad u_x = \frac{dx}{dt} = \frac{dx'}{\gamma(dt' + \frac{v}{c^2} dz')} = \frac{u_x'}{\gamma(1 + \frac{v}{c^2} \frac{dz'}{dx'})}$$

$$u_y = \frac{dy}{dt} = \frac{u_y'}{\gamma(1 + \frac{v}{c^2} \frac{dz'}{dx'})}$$

$$u_z = \frac{dz}{dt} = \frac{\gamma(dz' + v dt')}{\gamma(dt' + \frac{v}{c^2} dz')} = \frac{u_z' + v}{1 + \frac{v}{c^2} \frac{dz'}{dx'}}$$

Differenzieren:

$$du_x = \frac{du_x'}{\gamma(1 + \frac{v}{c^2} \frac{dz'}{dx'})} - \frac{u_x'}{\gamma(1 + \frac{v}{c^2} \frac{dz'}{dx'})^2} \cdot \frac{v}{c^2} du_z' \xrightarrow{(\vec{u}'=0)} \gamma^{-1} du_x'$$


Teilwende für $du_y = \gamma^{-1} du_y'$

$$du_z = \frac{du_z'}{1 + \frac{v}{c^2} \frac{dz'}{dx'}} - \frac{u_z' + v}{(1 + \frac{v}{c^2} \frac{dz'}{dx'})^2} \cdot \frac{v}{c^2} du_z' \rightarrow du_z' - \frac{v^2}{c^2} du_z' = (1 - \beta^2) du_z'$$

$$\text{Für } dt = \gamma dt' (1 + \frac{v}{c^2} \frac{dz'}{dx'}) \text{ für } dt = \gamma dt'$$

$$\Rightarrow \underline{a_x = \frac{du_x}{dt} = (1 - \beta^2) a_x'}, \text{ Teilwende } \underline{a_y = (1 - \beta^2) a_y'}$$

$$\underline{a_z = \frac{du_z}{dt} = (1 - \beta^2)^{3/2} a_z'}$$

2)  $v^0 = \frac{dN}{dt^0}$, hvor dN er antallet bølger emitteret i tiden dt^0 i det instantane hvilesystemet S' .
En har $dt^0 = d\tau$ (egensiden tid præpareret).
Hvis t er observatørens tid vil $dt = \gamma d\tau$.

Frekvensen ν målt i rummet S er

$$\underline{\nu = \frac{dN}{dt} = \frac{dN}{\gamma d\tau} = \frac{\nu^0}{\gamma} = \frac{\nu^0}{\sqrt{1 - \beta^2}}, \quad \beta = \frac{u}{c}}$$

Transversal Dopplereffekt.