11.10 The β function in $\lambda \phi^4$.

a.) In DR, we have to add a factor $\mu^{-\varepsilon}$ to the relation between renormalised and bare coupling, $\lambda(\mu) = \mu^{-\varepsilon} Z_{\phi}^2 Z_{\lambda}^{-1} \lambda_0$. Setting $\bar{Z}^{-1} \equiv Z_{\phi}^2 Z_{\lambda}^{-1}$ and using the definition of the beta function gives

$$\beta(\lambda) = \mu \frac{\partial}{\partial \mu} \left(\mu^{-\varepsilon} \bar{Z}^{-1} \right) \lambda_0 \tag{210}$$

$$= -\varepsilon\mu^{-\varepsilon}\bar{Z}^{-1}\lambda_0 - \mu\mu^{-\varepsilon}\bar{Z}^{-2}\frac{\partial Z}{\partial\mu}\lambda_0$$
(211)

$$= -\varepsilon\lambda - \frac{\mu}{\bar{Z}}\frac{\partial Z}{\partial\mu}\lambda \tag{212}$$

b.) Looking back at our one-loop calculations, we have $Z_{\phi} = 1$ and $Z_{\lambda}^{-1} = 1 - \frac{3\lambda}{16\pi^2 \varepsilon}$. Thus $\overline{Z} = 1 + 3\lambda/(16\pi^2 \varepsilon)$ and

$$\mu \frac{\partial Z}{\partial \mu} = \frac{3}{16\pi^2 \varepsilon} \mu \frac{\partial \lambda}{\partial \mu} = \frac{3}{16\pi^2 \varepsilon} \beta(\lambda) \,.$$

Using $\bar{Z}^{-1} = 1 + \mathcal{O}(\lambda)$ we can write this also as

$$\frac{\mu}{\bar{Z}}\frac{\partial\bar{Z}}{\partial\mu} = \frac{3}{16\pi^2\varepsilon}\beta(\lambda)\,.$$

Inserting in (3) gives

$$eta(\lambda) = -arepsilon\lambda - rac{3\lambda}{16\pi^2arepsilon}eta(\lambda)$$

Solving for β and taking the limit $\varepsilon \to 0$,

$$\beta(\lambda) = -\varepsilon\lambda\left(1 - \frac{3\lambda}{16\pi^2\varepsilon}\right) = \frac{3\lambda^2}{16\pi^2}.$$

Note that after this lengthy operation, we end up with a simple result which holds in general: The beta function is given by the coefficient of the pole in the coupling constant renormalisation constant.

c.) We set $d\lambda/d \ln \mu \equiv \dot{\lambda}$. Then $\dot{\lambda} = \beta(\lambda) = b_1 \lambda^2 + b_2 \lambda^3 + \dots$ is the perturbative expansion of the beta function in a given scheme. The couplings in two different schemes are connected by

$$\tilde{\lambda} = \lambda + c_2 \lambda^2 + \dots \tag{213}$$

$$\lambda = \tilde{\lambda} - c_2 \tilde{\lambda}^2 + \dots \tag{214}$$

Differentiating gives

 $\dot{\lambda} = \dot{\tilde{\lambda}} - 2c_2\tilde{\lambda}\dot{\tilde{\lambda}} + \ldots = (1 - 2c_2\tilde{\lambda})\dot{\tilde{\lambda}} + \ldots$

Solving for $\dot{\tilde{\lambda}}$, inserting first the expansion for $\dot{\lambda}$ and then for λ results in

$$\tilde{\lambda} = (1 + 2c_2\tilde{\lambda} + \dots)\dot{\lambda} \tag{215}$$

$$= (1 + 2c_2\tilde{\lambda} + \dots)(b_1\lambda^2 + b_2\lambda^3 + \dots)$$
(216)

$$= (1 + 2c_2\tilde{\lambda} + \dots)[b_1(\tilde{\lambda} - c_2\tilde{\lambda}^2 + \dots)^2 + b_2(\tilde{\lambda} - c_2\tilde{\lambda}^2 + \dots)^3 + \dots]$$
(217)

$$=b_1\tilde{\lambda}^2 + b_2\tilde{\lambda}^3 + \dots$$
(218)

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Thus c_2 cancels from the first two terms, and the β function is up to two loop independent from the renormalisation scheme.