

16.3 Scalar instantons.

Evaluating $\square\phi$ for the ansatz gives

$$A^{-1} \frac{d^2\phi}{dr^2} = \frac{8\rho r^2}{(r^2 + \rho^2)^3} - \frac{2\rho}{(r^2 + \rho^2)^2}$$

and

$$A^{-1} \frac{3}{r} \frac{d\phi}{dr} = -\frac{6\rho}{(r^2 + \rho^2)^2}.$$

Summing the two terms, we find

$$\frac{d^2\phi}{dr^2} + \frac{3}{r} \frac{d\phi}{dr} = 8A\rho \left[\frac{r^2}{(r^2 + \rho^2)^3} - \frac{1}{(r^2 + \rho^2)^2} \right] = -8A \left[\frac{\rho}{(r^2 + \rho^2)} \right]^3.$$

The potential term gives

$$-\frac{dV}{d\phi} = -\lambda\phi^3.$$

This fixes A as $-8A - \lambda A^3 = 0$ or $A = \sqrt{-8/\lambda}$. Thus the coupling constant λ should be negative.

Next we calculate the action,

$$S = 2\pi^2 \int_0^\infty dr r^3 \left[-\frac{1}{2} \phi \left(\frac{d^2\phi}{dr^2} + \frac{3}{r} \frac{d\phi}{dr} \right) + V(\phi) \right]$$

With

$$\rho^4 \int_0^\infty dr \frac{r^3}{(r^2 + \rho^2)^4} = \frac{1}{12}$$

we find

$$S = \frac{\pi^2}{6} \left[4A^2 + \frac{\lambda}{4} A^4 \right] = \frac{\pi^2}{6} \left[4(-8/\lambda) + \frac{\lambda}{4} (-8/\lambda)^2 \right] = \frac{8\pi^2}{3\lambda}.$$

For background and generalisation to arbitrary spacetime dimension d see S. Fubini, “A New Approach to Conformal Invariant Field Theories,” *Nuovo Cim. A* **34**, 521 (1976).