

Arbeid og energi

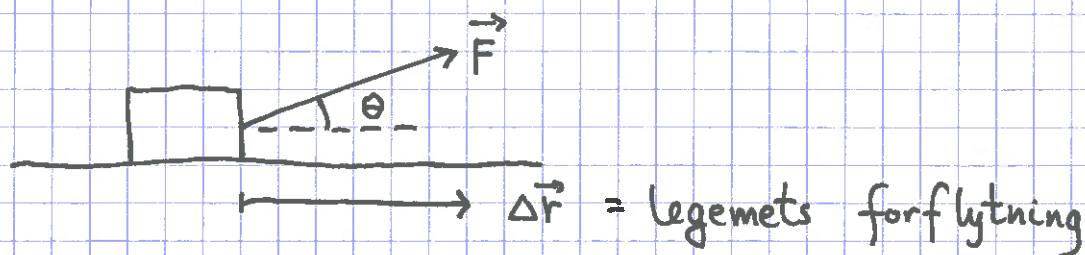
[YF 6, 7 ; LL 4]

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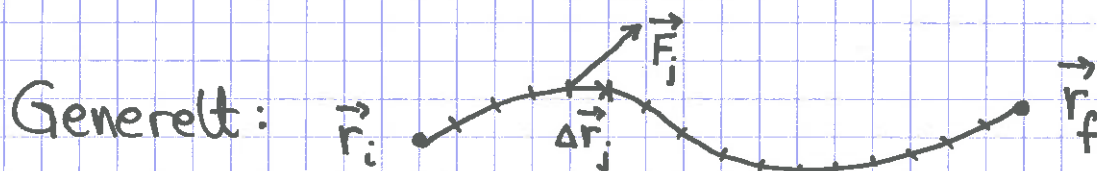
Arbeid

[YF 6.1-6.3 ; LL 4.1]



$$\Delta W = \vec{F} \cdot \Delta \vec{r} = F \cdot \Delta r \cdot \cos \theta = \text{arbeid utført av ytre kraft } \vec{F} \text{ på legemet}$$

$$[W] = \text{N} \cdot \text{m} = \text{J} \text{ (joule)}$$



Arbeid utført ved forflytning fra \vec{r}_i til \vec{r}_f :

$$W = \sum_j \Delta W_j = \sum_j \vec{F}_j \cdot \Delta \vec{r}_j \stackrel{\Delta r_j \rightarrow 0}{=} \int_{\vec{r}_i}^{\vec{r}_f} \vec{F} \cdot d\vec{r}$$

Effekt

[YF 6.4 ; LL 4.1]

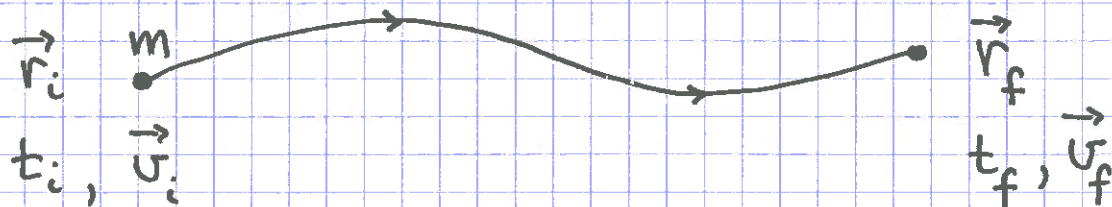
= arbeid (energi) pr tidsenhet

$$P = \frac{dW}{dt} = \frac{\vec{F} \cdot d\vec{r}}{dt} = \vec{F} \cdot \vec{v}$$

$$[P] = \text{J/s} = \text{W} \text{ (watt)}$$

Kinetisk energi [YF 6.2; LL 4.2]

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$$W = \int_{r_i}^{r_f} \vec{F} \cdot d\vec{r} \stackrel{N2}{=} \int_{t_i}^{t_f} m \frac{d\vec{u}}{dt} \cdot \vec{u} dt = m \int_{u_i}^{u_f} \vec{u} \cdot d\vec{u}$$

$$d(u^2) = d(\vec{u} \cdot \vec{u}) = d\vec{u} \cdot \vec{u} + \vec{u} \cdot d\vec{u} = 2\vec{u} \cdot d\vec{u}$$

$$\Rightarrow W = \frac{1}{2} m \int_{u_i^2}^{u_f^2} d(u^2) = \frac{1}{2} m u_f^2 - \frac{1}{2} m u_i^2$$

Definerer: $K = \frac{1}{2} m u^2 =$ kinetisk energi

Dermed:

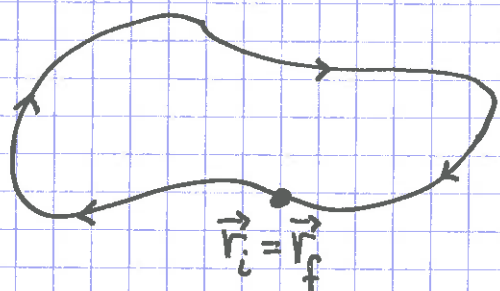
$$W = K_f - K_i = \Delta K$$

Arbeid W utført på legemet tilsvarer endringen ΔK i legemets kin. energi.

Konservativ kraft [YF 7.3 ; LL 4.4]

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System uten tap av mekanisk energi til andre energiformer (som varme) er konservativt.



Hvis \vec{F} er konservativ, er $K_f = K_i$ ($|\vec{v}_f| = |\vec{v}_i|$),
dvs $W = \Delta K = 0$

Dermed:

$$\oint \vec{F} \cdot d\vec{r} = 0$$

($\oint \dots$: Integral rundt lukket kurve.)

Potensiell energi [YF 7.1-7.4 ; LL 4.3-4.4]

Med kons. kraft \vec{F} er potensiell energi

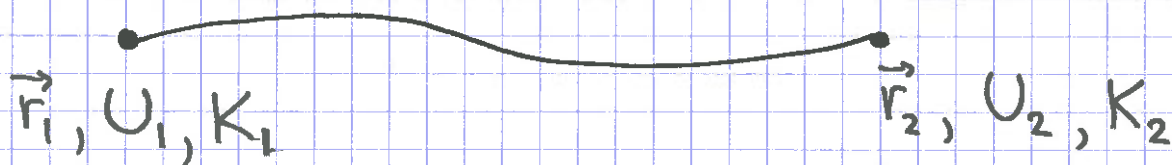
$$U(\vec{r}) = - \int_{\vec{r}_0}^{\vec{r}} \vec{F} \cdot d\vec{r}$$

Har her valgt $U(\vec{r}_0) = 0$.

Dvs: Kun forskjeller i pot. energi har fysisk betydning.

Mekanisk energibevarelse [YF 7.1-7.3; LL 4.5]

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$$\begin{aligned} U_1 - U_2 &= \int_{\vec{r}_1}^{\vec{r}_2} \vec{F} \cdot d\vec{r} + \int_{\vec{r}_2}^{\vec{r}_1} \vec{F} \cdot d\vec{r} \\ &= \int_{\vec{r}_1}^{\vec{r}_2} \vec{F} \cdot d\vec{r} = W = K_2 - K_1 \end{aligned}$$

$$\Rightarrow \boxed{K_1 + U_1 = K_2 + U_2}$$

Dvs: Total mekanisk energi,

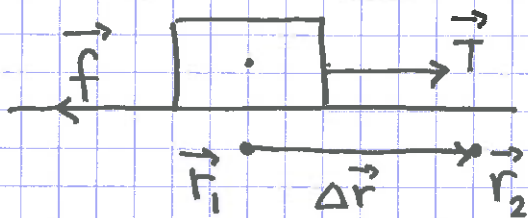
$$E = K + U,$$

er konstant ("bevart")

for et konservativt system.

Friksjonsarbeid

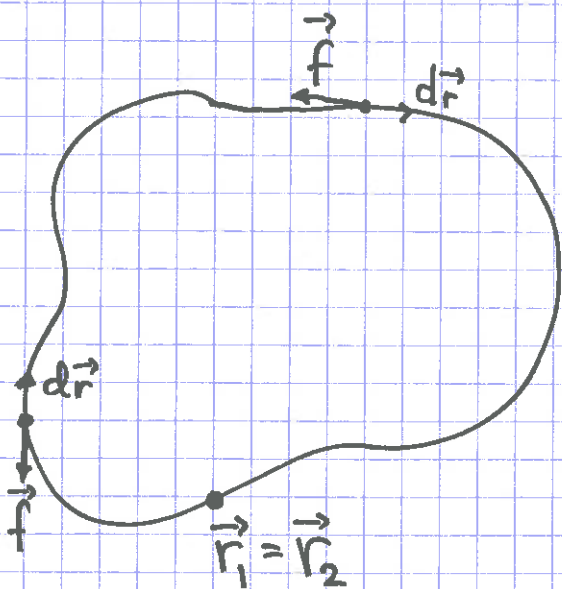
[YF 7.3 ; LL 4.5]



$$W_f = \int_{r_1}^{r_2} \vec{f} \cdot d\vec{r} < 0 \quad \text{fordi } \vec{f} \text{ alltid er}$$

rettet mot $d\vec{r}$. Friksjonsarbeidet W_f

"går tapt" : Mek. energi \rightarrow Varme

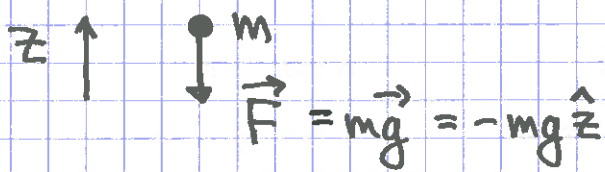


$$\Rightarrow \oint \vec{f} \cdot d\vec{r} < 0$$

\Rightarrow Friksjonskraften \vec{f} er ikke konservativ.

Eks: Fritt fall

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Anta $U(0)=0$, $v(0)=0$.
Bestem $U(z)$ og $v(z)$
for $z < 0$.

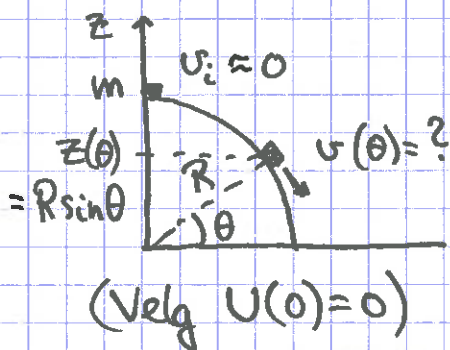
Løsn:
$$U(z) = - \int_0^z \underbrace{(-mg\hat{z})}_{\vec{F}} \cdot \underbrace{(\hat{z} dz)}_{d\vec{r}} = \underline{\underline{mgz}}$$

$$E(0) = U(0) + K(0) = 0$$

$$\Rightarrow U(z) + K(z) = mgz + \frac{1}{2} m v(z)^2 = 0$$

$$\Rightarrow v(z) = \underline{\underline{\sqrt{-2gz}}}$$

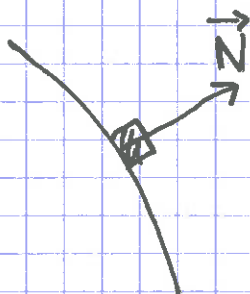
Eks: Gli på kuleflate



Løsn: E er bevart (uten friksjon)

$$\Rightarrow \frac{1}{2} m v(\theta)^2 + mgz(\theta) = mgR$$

$$\Rightarrow \underline{\underline{v(\theta) = \sqrt{2gR(1 - \sin\theta)}}}$$



- Myster kontakten med underlaget når $N=0$; "skrått kast" derfra. Hvor skjer dette?
- Hva med friksjon?
- Hva hvis legemet kan rulle?
Evt rulle og gli samtidig, dvs slure?

Impuls [YF8; LLS]

(= bevegelsesmengde, "(linear) momentum")

$$N2: \vec{F} = m \frac{d\vec{v}}{dt} \quad \text{hvis} \quad \frac{d}{dt} (m\vec{v})$$

$m = \text{konst.}$

impuls = masse · hastighet

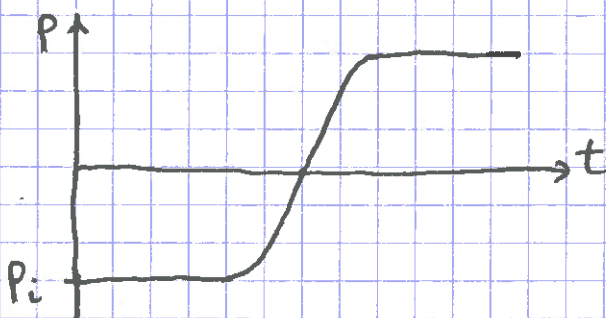
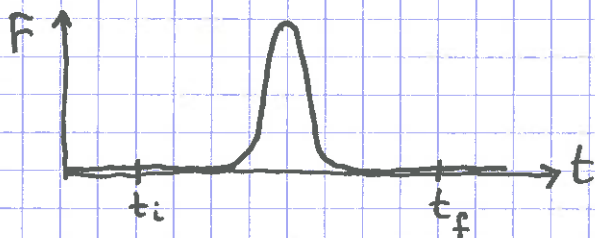
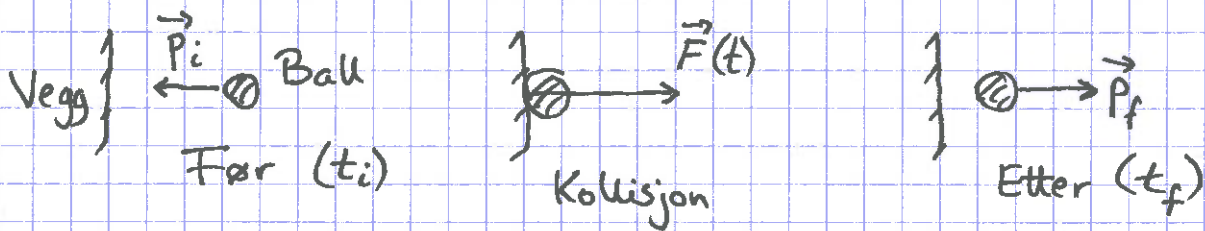
$$\vec{p} = m \cdot \vec{v} \quad [p] = \text{kg m/s}$$

$$\Rightarrow \boxed{\vec{F} = \frac{d\vec{p}}{dt}} \quad N2$$

⇒ Lov om impulsbevarelse:

Hvis sum av ytre krefter på et legeme er null, er legemets impuls bevart: $\vec{F} = 0 \Rightarrow \vec{p} = \text{konst.}$

Ytre $\vec{F} \Rightarrow$ impulsendring:



$$\begin{aligned} \Delta \vec{p} &= \vec{p}_f - \vec{p}_i \\ &= \int_{\vec{p}_i}^{\vec{p}_f} d\vec{p} \\ &= \int_{t_i}^{t_f} \vec{F}(t) dt \end{aligned}$$

Kollisjoner [YF 8.3, 8.4; LL 5.3]

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Elastisk støt: $\Delta K = 0$ (energi bevart)

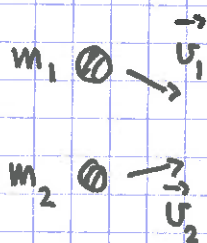
Uelastisk støt: $\Delta K < 0$ (— ikke bevart)

Fullstendig uelastisk støt: Legemene henger sammen, felles hastighet etter kollisjonen, max. energitap $|\Delta K|$.

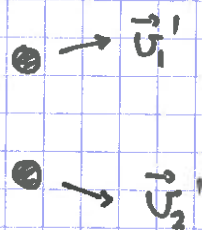
Tapt mek. energi $\Delta K \rightarrow$ deformasjon, lyd, varme...

Hvis $\vec{F}_{\text{ytre}} = 0$, er $\Delta \vec{p} = 0$ for alle typer kollisjoner.

Indre krefter \vec{F}_{ij} endrer ikke systemets totale impuls:



Kollisjon
(kort varighet Δt)



Etter

$$N3: \vec{F}_{21} = -\vec{F}_{12} \quad \stackrel{N2}{\Rightarrow} \quad \frac{d\vec{p}_1}{dt} = -\frac{d\vec{p}_2}{dt}$$

$$\Rightarrow \frac{d}{dt} (\vec{p}_1 + \vec{p}_2) = \frac{d}{dt} \vec{p}_{\text{tot}} = 0$$

$$\Rightarrow \vec{p}_{\text{tot}} = \vec{p}_1 + \vec{p}_2 = \text{konst.}$$

Eks: Tennis-serve; anslå midlere kraft (28)

$\langle F \rangle$ og sammenlign med mg .

Løsn: $m = 57 \text{ g}$, $\Delta v = 263.4 \text{ km/h} = 73.2 \text{ m/s}$, $\Delta t \sim 7 \text{ ms}$

$$\Rightarrow \langle F \rangle = \frac{m \Delta v}{\Delta t} = \frac{57 \cdot 10^{-3} \text{ kg} \cdot 73.2 \text{ m/s}}{7 \cdot 10^{-3} \text{ s}} \approx \underline{600 \text{ N}}$$

$$\langle F \rangle / mg = \frac{\Delta v}{g \Delta t} = \frac{\langle a \rangle}{g} \approx \frac{10^4 \text{ m/s}^2}{10 \text{ m/s}^2} = \underline{1000}$$

\Rightarrow OK å neglisjere den ytre kraften mg i kollisjonen

Sentralt støt [YF 8.2-8.4; LL 5.3]

Før (i): $m \rightarrow v$ $V \leftarrow M$

Etter (f): $v' \leftarrow m$ $M \rightarrow V'$

Fortegn: $\rightarrow +$ (dvs $v, V' > 0$ og $v', V < 0$ i fig.)

$$\Delta p = 0 \Rightarrow \underbrace{mv + MV}_{P_i} = \underbrace{mv' + MV'}_{P_f}$$

(a) Fullstendig uelastisk støt (enklest):

$$v' = V' = \frac{mv + MV}{m + M}$$

(b) Delvis uelastisk støt: Har 1 lign. ($\Delta p = 0$) for

2 ukjente (v', V') \Rightarrow Trenger en opplysning til

for å fastlegge v' og V' .

(c) Elastisk støt, $\Delta K = 0$:

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$$\underbrace{\frac{1}{2} m v^2 + \frac{1}{2} M V^2}_{K_i} = \underbrace{\frac{1}{2} m v'^2 + \frac{1}{2} M V'^2}_{K_f}$$

"Tricks" (lett omskiving):

$$m(v + v')(v - v') = M(V' + V)(V' - V) \quad (1) \quad (\Delta K = 0)$$

$$m(v - v') = M(V' - V) \quad (2) \quad (\Delta p = 0)$$

(1) delt på (2) gir:

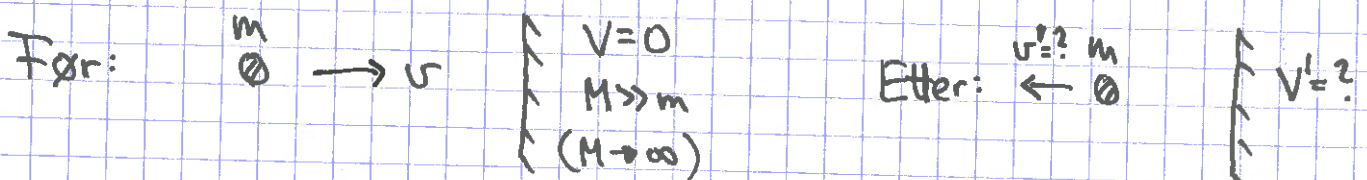
$$\begin{aligned} v + v' &= V + V' \\ \Rightarrow v' - V' &= -(v - V) \quad (3) \end{aligned}$$

Fra (2) og (3) fås:

$$v' = \frac{M}{m+M} (2V + v \cdot \frac{m-M}{M})$$

$$V' = \frac{m}{M+m} (2v + V \cdot \frac{M-m}{m})$$

Eks: Ball mot vegg, elastisk koll.



- Bestem v' og V'
- Sjekk at $\Delta p = 0$ og $\Delta K = 0$

Løsn: $v' = \frac{M}{m+M} (0 + v \cdot \frac{m-M}{M}) \approx \frac{M}{M} \cdot v \cdot (-\frac{M}{M}) = -v$ (30)

$V' = \frac{m}{M+m} (2v + 0) \approx \frac{m}{M} \cdot 2v = 0$ } Som vendet

Impulsbevarelse:

$$p = mv, P = MV = 0, p' = mv' = -mv,$$

$$P' = MV' \approx M \cdot \frac{m}{M} \cdot 2v = 2mv$$

$$\Rightarrow \Delta p = -mv + 2mv - mv - 0 = 0; \text{ OK}$$

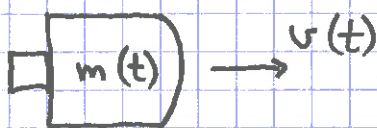
Energi bevarelse:

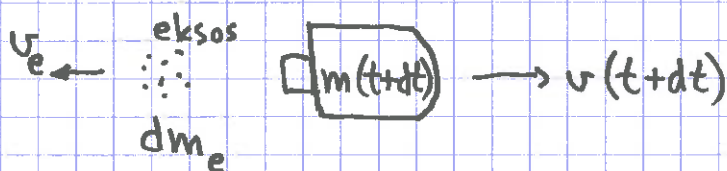
$$K_m = \frac{1}{2}mv^2, K_M = \frac{1}{2}MV^2 = 0, K_m' = \frac{1}{2}m(v')^2 = \frac{1}{2}mv^2,$$

$$K_M' = \frac{1}{2}MV'^2 = \frac{1}{2}M \left(\frac{m}{M} \cdot 2v\right)^2 = \frac{2m^2v^2}{M} = 0$$

$$\Rightarrow \Delta K = \frac{1}{2}mv^2 - \frac{1}{2}mv^2 = 0; \text{ OK}$$

Rakettprinsipp [YF 8.6; LL 5.4]

"Før" (t):  $p(t) = m(t)v(t)$

"Etter" (t+dt): 

$$p(t+dt) = \underbrace{m(t+dt)}_{m(t)+dm} \underbrace{v(t+dt)}_{v(t)+dv} + \underbrace{dm_e}_{-dm} \cdot \underbrace{v_e(t)}_{v(t)+u}$$

med $u =$ eksosens hastighet relativt raketten ($u < 0$)

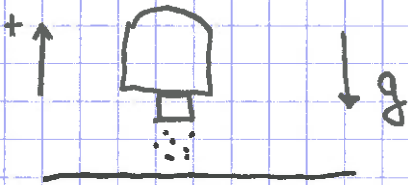
$dm =$ rakettenes masseendring fra t til $t+dt$ ($dm < 0$)

$$\Rightarrow p(t+dt) = \underbrace{m(t)v(t)}_{= p(t)} + m(t)dv + \underbrace{dm \cdot v(t) - dm \cdot v(t)}_{= 0} - dm \cdot u$$

"Outer space" : $F_{ytre} = 0$

$$\begin{aligned} \stackrel{N_2}{\Rightarrow} p(t+dt) &= p(t) \\ \Rightarrow m(t)dv &= u dm \\ \Rightarrow m \frac{dv}{dt} &= u \frac{dm}{dt} = u \dot{m} \\ \Rightarrow m \cdot a &= F_{skyr}, \text{ med} \\ \text{skyrkraft } F_{skyr} &= u \dot{m} > 0 \end{aligned}$$

I tyngdefeltet :



$$F_{ytre} = -mg$$

Total kraft på "rest-raketten" :

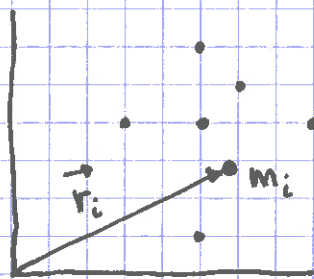
$$F_{skyr} + F_{ytre} = u \dot{m} - mg$$

$$\stackrel{N_2}{\Rightarrow} u \dot{m} - mg = ma \quad (\text{Øving})$$

Til nå: Punktmasser.

Nå: Partikkelsystemer. Stive legemer.

Massesenter. Tyngdepunkt [YF 8.5, oppg 8.115+8.116; LL 5.6, 5.8, 6.1]

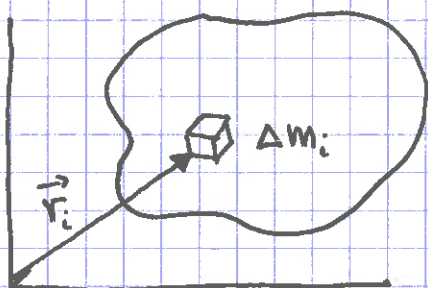


Massesenter (CM, "center of mass") for N punktmasser m_1, m_2, \dots, m_N i posisjoner $\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N$:

$$\vec{R}_{CM} = \frac{\sum_{i=1}^N m_i \vec{r}_i}{\sum_{i=1}^N m_i} = \frac{1}{M} \sum_i m_i \vec{r}_i$$

Total masse: $M = \sum_i m_i$

Kontinuerlig massefordeling:



$$\vec{R}_{CM} = \frac{\sum_i \Delta m_i \vec{r}_i}{\sum_i \Delta m_i} \xrightarrow{\Delta m_i \rightarrow 0} \frac{\int \vec{r} dm}{\int dm}$$

$$= \frac{1}{M} \int \vec{r} dm$$

(integrerer over der vi har masse!)

Masselement:

$dm = \rho dV$, ρ = masse pr volumenhet, dV = volumenelement (3D)

$dm = \sigma dA$, σ = " " flate " ", dA = flate " " (2D)

$dm = \lambda dl$, λ = " " lengde " ", dl = lengde " " (1D)