

KAP 5: Beregningslign. for stive legemer

- $\vec{L} = \vec{I} \vec{\omega}$; $I_{jk} = \int_V \rho(\vec{r}) (r^2 \delta_{jk} - x_j x_k) dV$
- $T = \frac{1}{2} \vec{\omega}^T \vec{I} \vec{\omega}$; $I = \vec{n}^T \vec{I} \vec{n}$
- Hovedakser $x_1, x_2, x_3 \Rightarrow \vec{I} = \begin{pmatrix} I_1 & 0 & 0 \\ 0 & I_2 & 0 \\ 0 & 0 & I_3 \end{pmatrix} \Rightarrow L_i = I_i \omega_i$
 $T = \sum_j \frac{1}{2} I_j \omega_j^2$
- Eulerligningene: $\frac{d\vec{L}}{dt} + \vec{\omega} \times \vec{L} = \vec{N}$ (sett fra roterende system)
- Fri rotasjon - presesjon ; Gyroskopeffekt ; Snurrebass

KAP 6: Små oscillasjoner

- små utsving fra stabil likevekt $\Rightarrow V \approx \frac{1}{2} \left(\frac{\partial^2 V}{\partial q_i \partial q_j} \right)_0 \eta_i \eta_j \equiv \frac{1}{2} V_{ij} \eta_i \eta_j$
- kin. energi : $T = \frac{1}{2} m_{ij} \dot{\eta}_i \dot{\eta}_j \equiv \frac{1}{2} T_{ij} \dot{\eta}_i \dot{\eta}_j$
- $L = T - V$; $\frac{d}{dt} \frac{\partial L}{\partial \dot{\eta}_i} - \frac{\partial L}{\partial \eta_i} = 0 \Rightarrow T_{ij} \ddot{\eta}_j + V_{ij} \eta_j = 0$
- $\eta_i = A_i e^{-i\omega t} \Rightarrow \underbrace{\det(V - \omega^2 T)}_{\text{sekulær ligning}} = 0 \Rightarrow \text{egenfrekvenser } \omega_\alpha$
- $A_{i\alpha} \sim \Delta_{i\alpha} \sim (i, \alpha)\text{-minoren til } |V - \omega_\alpha^2 T|$
- $\text{Re } \eta_i = \sum_\alpha \text{Re } \eta_{i\alpha} = \text{Re } \sum_\alpha C_\alpha \Delta_{i\alpha} e^{-i\omega_\alpha t} = \sum_\alpha \Delta_{i\alpha} \Theta_\alpha(t)$
 $\Rightarrow \dots \Rightarrow \ddot{\Theta}_\alpha + \omega_\alpha^2 \Theta_\alpha = 0$ (dekoblet!)
 $\uparrow$ Normalkoordinater
- Systemet oscillerer i normalmode α med egenfrekvens ω_α