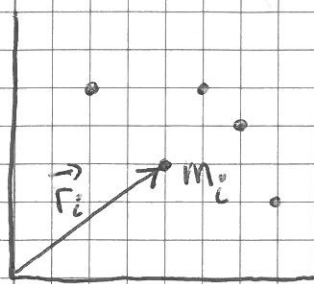


Massecenter

[YF 8.5, oppg 8.115+116; LL 5.6, 5.8, 6.1]

(28)

(= tyngdepunkt hvis g = konstant i hele systemet)



Massecenter (CM) for N punktmasser
 $m_1, m_2, \dots, m_N$ i pos. $\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N$:

$$\vec{R}_{CM} = \frac{\sum_i m_i \vec{r}_i}{\sum_i m_i} = \frac{1}{M} \sum_i m_i \vec{r}_i$$

$$\text{Total masse: } M = \sum_i m_i$$

Kontinuerlig massefordeling: $m_i \rightarrow \Delta m_i \xrightarrow{\Delta m_i \rightarrow 0} dm$ (masseelement)
 $\Rightarrow \vec{R}_{CM} = \frac{1}{M} \int \vec{r} dm$; $M = \int dm$
↑ integral over der vi har masse

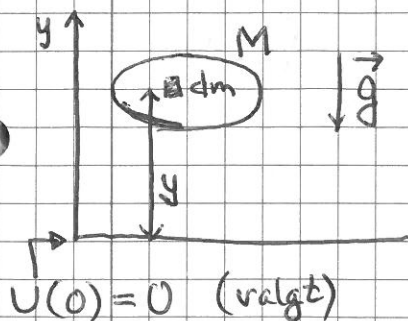
Må typisk uttrykke dm ved koordinater:

$$dm = \rho dV \text{ (3D)}, \quad dm = \sigma dA \text{ (2D)}, \quad dm = \lambda dl \text{ (1D)}$$

ρ, σ, λ = masse pr. hhv. volum-, flate-, lengdeenhet
 dV, dA, dl = hhv. volum-, flate-, lengdeelement

Hvis massen er uniformt fordelt, er $\frac{dm}{M} = \frac{dV}{V}$ etc.

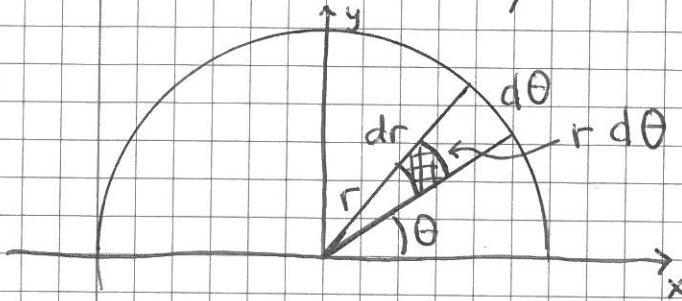
Eks: Pot. energi i tyngdefeltet (antar g = konstant)



$$U = \int g y dm = g \int y dm = g M Y_{CM}$$

Dvs: Total U som om hele massen $M = \int dm$ er samlet i høyden $Y_{CM} = \frac{1}{M} \int y dm$,
f.eks. i $\vec{R}_{CM}$

Eks: Halvsirkulær tynn skive, radius R



$$\frac{dm}{M} = \frac{dA}{A}$$

$$A = \pi R^2 / 2, \quad dA = dr \cdot r d\theta$$

$$y = r \sin \theta$$

$$X_{cm} = 0 \text{ pga symmetri} \Rightarrow \vec{R}_{cm} = Y_{cm} \hat{y}$$

$$Y_{cm} = \frac{1}{M} \int y dm = \frac{1}{A} \int y dA$$

$$= \frac{2}{\pi R^2} \int_{r=0}^R \int_{\theta=0}^{\pi} r \sin \theta r d\theta dr$$

$$= \frac{2}{\pi R^2} \int_0^R r^2 dr \int_0^{\pi} \sin \theta d\theta = \frac{2}{\pi R^2} \cdot \frac{1}{3} R^3 \cdot \underbrace{\int_0^{\pi} (-\cos \theta) d\theta}_{=2}$$

$$= \underline{\underline{\frac{4}{3\pi} R \approx 0.42 R}}$$

[Vis at $Y_{cm} = 2R/\pi$ for bølge (1D)
og $Y_{cm} = 3R/8$ for halvkule (3D)]

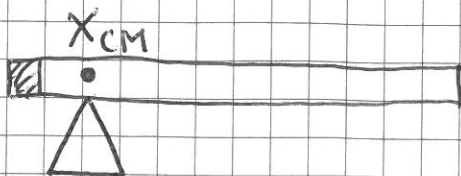
Eks: Rør med lite lodd i enden



$$m = 165g \quad M = 305g$$

$$X_{cm} = \frac{1}{m+M} \left\{ M \cdot 0 + \int_0^L x dm \right\} = \frac{1}{m+M} \int_0^L x m \frac{dx}{L} = \underline{\underline{\frac{m}{2(m+M)} L}}$$

$$= \frac{165}{940} L \approx \underline{\underline{0.18 L}}$$

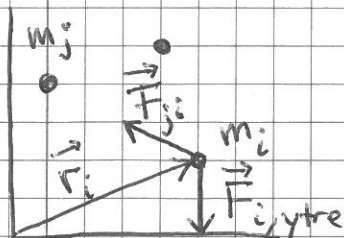


Tyngdepunktbevegelsen

[YF 8.5; LL 5.8]

(30)

- Exp. (kast) viser at ~~CM~~ CM beveger seg som punktmasse i CM! (Parabel, se s. 4)



N punktmasser, $\vec{R}_{CM} = \frac{1}{M} \sum_i m_i \vec{r}_i$, $M = \sum_i m_i$

N2 for m_i :

$$m_i \ddot{\vec{r}}_i = \vec{F}_{i,ytre} + \underbrace{\sum_{j \neq i} \vec{F}_{ji}}_{\text{total indre kraft p\aa } m_i}$$

total ytre kraft p\aa m_i

total indre kraft p\aa m_i

Legger sammen N2 for alle $i=1,2,\dots,N$:

$$\sum_i m_i \ddot{\vec{r}}_i = \underbrace{\sum_i \vec{F}_{i,ytre}}_{\substack{\text{total ytre kraft} \\ \vec{F}_{ytre} \text{ p\aa systemet}}} + \underbrace{\sum_i \sum_{j \neq i} \vec{F}_{ji}}_{\substack{= \vec{F}_{21} + \vec{F}_{12} + \dots + \vec{F}_{N,N-1} + \vec{F}_{N-1,N} \\ = 0}}$$

$$\frac{d^2}{dt^2} \left(\sum_i m_i \vec{r}_i \right) = \frac{d^2}{dt^2} (M \vec{R}_{CM}) = M \ddot{\vec{R}}_{CM}$$

Konklusjon:

$$\boxed{M \ddot{\vec{R}}_{CM} = \vec{F}_{ytre}}$$

CM beveger seg som om hele M er samlet i CM og utsettes for summen av alle ytre krefter p\aa systemet, $\vec{F}_{ytre}$

I tillegg kommer rotasjon om CM og

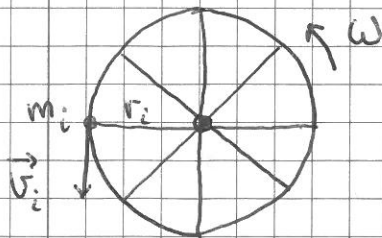
vibrasjon om CM.

ROTASJON [YF 9,10 ; LL 6 (5)]

(31)

• Innledende observasjoner:

- Ren rotasjon (hjul)



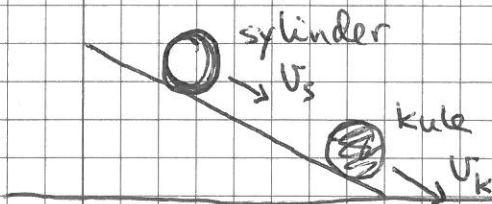
$$(CM \Rightarrow i \text{ ro} \Rightarrow K_{\text{trans}} = \frac{1}{2} M \vec{R}_{CM}^2 = 0$$

men $K_{\text{rot}} \neq 0$

$$\text{Total impuls } \vec{P} = \sum_i m_i \vec{v}_i = 0$$

men dreieimpulsen $\neq 0$

- Rulling



Hvorfor oppstår rotasjon?

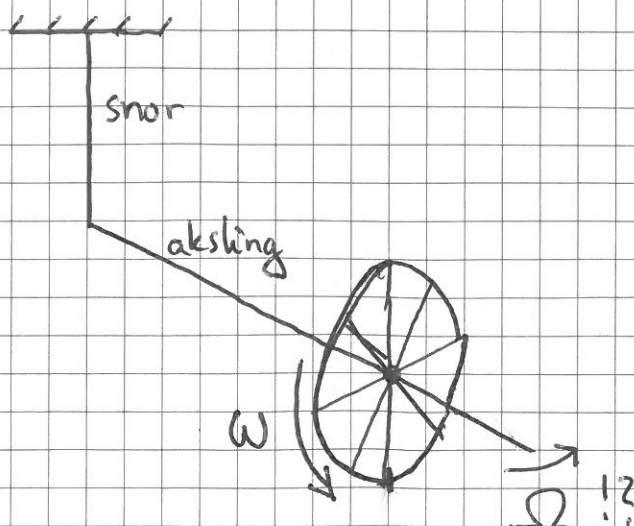
Hvor angriper kreftene?

Må ha dreiemoment!

Hvorfor blir $v_k > v_s$?

Har vi fiksjon her?

- Kompleks dynamikk



Gyroskop.

Presesjon

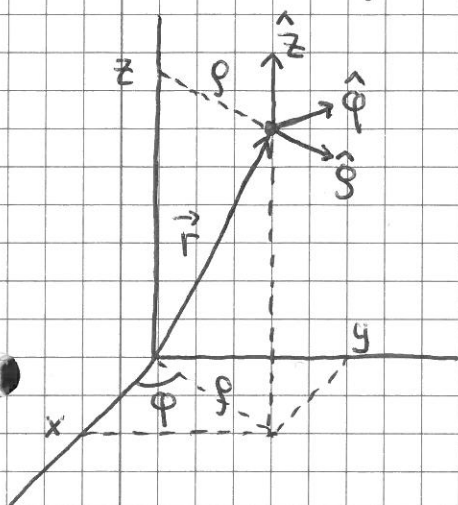
Sirkelbevegelse [YF 9.1-9.3; LL 1.8]

(32)

- Generalisering av s. 5-6.

Anta rotasjon om z-aksen

⇒ Bruker sylinderkoord. (= polarkoord. + z)



$$\vec{r} = x\hat{x} + y\hat{y} + z\hat{z} = \rho\hat{\rho} + z\hat{z}$$

$$x = \rho \cos \varphi, \quad y = \rho \sin \varphi, \quad z = z$$

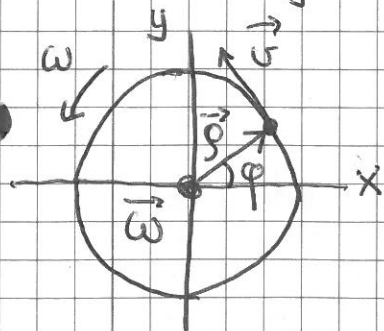
$$\rho = \sqrt{x^2 + y^2}, \quad \tan \varphi = y/x$$

$$r = \sqrt{x^2 + y^2 + z^2} = \sqrt{\rho^2 + z^2}$$

Bruker vinkelhastigheten som vektor til å "markere" rot.aksen:

$$\omega \rightarrow \vec{\omega} = \omega \hat{z} = \omega \hat{\omega}$$

Sett ned langs z-aksen:



$$\vec{v} = \frac{d\vec{s}}{dt} = \frac{\rho d\varphi}{dt} \hat{\varphi} = \rho \omega \hat{\varphi}$$

$$\vec{\omega} = \omega \hat{z}, \quad \vec{s} = \rho \hat{\rho}$$

(ut av planet)

$$\Rightarrow \boxed{\vec{v} = \vec{\omega} \times \vec{s}}$$

Retning: $\hat{\varphi} = \hat{z} \times \hat{\rho}$ OK

Abs.verdi: $v = \omega \rho$ OK