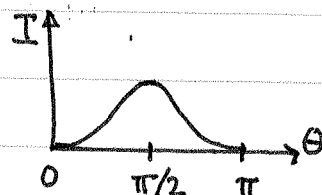
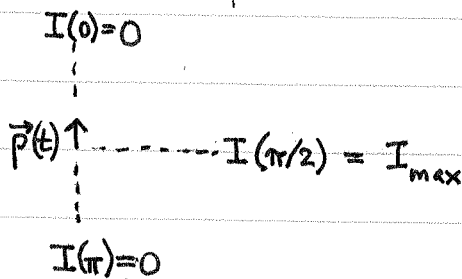


Vinkelfordeling: $E(\theta) \sim \sin \theta \Rightarrow I(\theta) \sim \sin^2 \theta$
(långt unna)



Frekvensavhengighet:

Maxwells lign. $\Rightarrow I \sim \langle \ddot{z} \rangle^2$; $\langle \ddot{z} \rangle$ = ladningenes midlere akselerasjon

$$z(t) = z_0 \sin \omega t \Rightarrow \ddot{z}(t) = -\omega^2 z_0 \sin \omega t \Rightarrow \langle \ddot{z} \rangle \sim \omega^2$$

$$\Rightarrow I \sim \omega^4$$

Utstrålt effekt: $\langle P \rangle = \frac{P_0^2 \omega^4}{12\pi \epsilon_0 c^3} \sim \omega^4$ (se f.eks. Griffiths) (Øving 8)

Tilsvarende for oscillerende magnetisk dipol:



$$\vec{m}(t) = I(t) \vec{A} = I_0 \vec{A} \sin \omega t = \vec{m}_0 \sin \omega t$$

Utstrålt effekt: $\langle P \rangle = \frac{\mu_0 m_0^2 \omega^4}{12\pi c^3} \sim \omega^4$ (se f.eks. Griffiths)

For ^{atom-}kjerne og atomer: $\langle P \rangle_{\text{elek.}} \gg \langle P \rangle_{\text{magn.}}$

For makroskopiske antenner: $\langle P \rangle_{\text{elek.}} \sim \langle P \rangle_{\text{magn.}}$