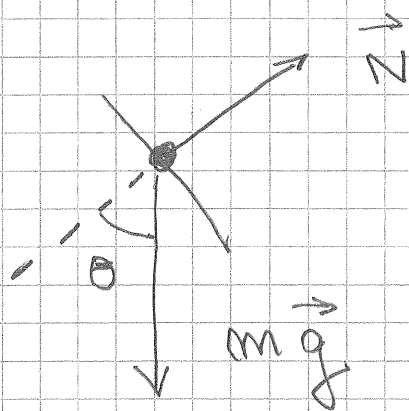


L1

a) Energierna bevaras i tryck deflater $\Rightarrow$

$$\frac{1}{2} M V^2 = m g \cdot R (1 - \cos \theta) \quad \text{"förel-höjden"}$$

$$V(\theta) = \sqrt{2 g R (1 - \cos \theta)}$$



Bewegung $\perp$ flate:

$$m g \cos \theta - N = m v \frac{V^2}{R}$$

$$N = m g \cos \theta - m v \frac{V^2}{R}$$

$$= m g \cos \theta - 2 m g (1 - \cos \theta)$$

$$= \underline{m g [3 \cos \theta - 2]}$$

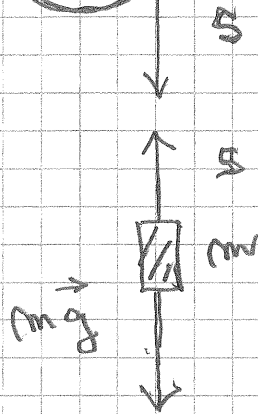
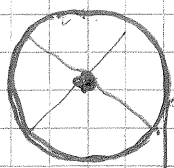
b) X läter fra flaten när $N = 0$:

$$\underline{\cos \theta_0 = \frac{2}{3}} \quad \theta_0 = 48,2^\circ$$

Farten är då

$$V(\theta_0) = \sqrt{2 g R (1 - \frac{2}{3})} = \underline{\underline{\sqrt{\frac{2}{3}} g R}}$$

$$\frac{I R}{M R}$$



a) Deler systemet og betrakter lod og hjul for seg:

Lodets translasjonsbevegelse.

$$m g - S = m A \quad (1)$$

Hjulets rotasjonsbevegelse.

$$\underbrace{S R}_{\text{Moment}} = I \alpha \quad (2)$$

Kinematiske betingelse

$$V = R \omega \Rightarrow A = R \alpha \quad (3)$$

$$(2) \text{ og } (3) \Rightarrow S R = \underbrace{M R^2}_{I} \cdot \underbrace{\frac{A}{R}}_{\alpha} \therefore S = M A$$

Insatt i (1):

$$m g - M A = m A \Rightarrow A = \frac{m g}{M + m}$$

$$S = \frac{m M g}{M + m}$$

Spesielt $m \ll M$

$$A \approx \frac{m}{M} g \ll g \quad S \approx m g$$

Rimelig: En tennisball vil ikke sette et kjerrehjul i merkelig rot. bevegelse.

b) Σ mechanische Energie : typisch falsch:

$$m g H = \underbrace{\frac{1}{2} m v^2}_{\text{Transl.}} + \underbrace{\frac{1}{2} I \omega^2}_{\text{Rot.}}$$

$$I = m R^2 \quad \omega = v/R \Rightarrow$$

$$m g H = \frac{1}{2} m v^2 + \frac{1}{2} M v^2 \Rightarrow$$

$$v = \sqrt{\frac{2 m g H}{m + M}}$$

$$K_{\text{ball}} = \frac{1}{2} m v^2 = \left(\frac{m}{m + M} \right) m g H$$

$$K_{\text{hül}} = \frac{1}{2} M v^2 = \left(\frac{M}{m + M} \right) m g H.$$

Kontrolle $m \gg M$.

$$v \approx \sqrt{2 g H} \quad K_{\text{ball}} \approx m g H \gg K_{\text{hül}}$$

sonst vortel. "100 kg-Ball mit
händelschnelle".

alt: Da $A = \text{konst.}$, er

$$v = A t \quad H = \frac{1}{2} A t^2 \Rightarrow$$

$$v = \sqrt{2 A H} = \sqrt{\frac{2 m g H}{m + M}}$$

sonst overflow.

L34

a) Ingen yttre krafter fortsatt
frå reaktioner från rot. axel:
Kraftmoment om axel lika
noll: Spinnkomp. långs
axel bevarat. Då bevaras
momentet om axel utan
må vinkelhastigheten öka.

$$\underbrace{I\omega_0 + mRv_0}_{\text{"För"}} = \underbrace{I\omega + mRv}_{\text{"Efter"}}$$

$$v_0 = R\omega_0, \quad v = R\omega \quad \Rightarrow$$

$$(I + mR^2)\omega_0 = (I + mR^2)\omega$$

$$\omega = \omega_0 \frac{I + mR^2}{I + mR^2} = \omega_0 \frac{\frac{1}{2}MR^2 + mR^2}{\frac{1}{2}MR^2 + mR^2}$$

$$\omega = \omega_0 \frac{1 + 2\frac{m}{M}}{1 + 2\frac{m}{M}\left(\frac{R}{R}\right)^2}$$

Kontroll: $r \rightarrow R \Rightarrow \omega \rightarrow \omega_0$

$$m \rightarrow 0 \Rightarrow \omega \rightarrow \omega_0$$

Spes: $r \rightarrow 0 \quad \underline{\underline{\omega \rightarrow \omega_0 \left(1 + \frac{2m}{M}\right)}}$

b)

$$\Delta K = K(r=0) - K(r=R)$$

$$= \frac{1}{2} I \omega_0^2 - \frac{1}{2} [I + mR^2] \omega_0^2 \Rightarrow$$

$$\Delta K = \frac{1}{4} MR^2 \left(1 + \frac{2m}{M}\right)^2 \omega_0^2 - \frac{1}{2} \left[\frac{1}{2} MR^2 + mR^2\right] \omega_0^2 \Rightarrow$$

$$\underline{\Delta K = \frac{1}{2} m \omega_0^2 R^2 \left[1 + \frac{2m}{M}\right]}$$

$$c) F_D = \frac{m v^3}{r} = m \omega^2 r$$

$$F_D(r) = m \omega_0^2 \left[\frac{1 + 2 \frac{m}{M}}{1 + \frac{2m}{M} \frac{r^2}{R^2}} \right]^2 r$$

$$W = \int_{r=R}^0 (-F_D) dr = \int_0^R F_D dr$$

$$= m \omega_0^2 \left(1 + 2 \frac{m}{M}\right)^2 \int_0^R \frac{r dr}{1 + \frac{2m}{M} \frac{r^2}{R^2}}$$

$$Z = \frac{2m}{MR^2} r^2 \Rightarrow r dr = \frac{MR^2}{4m} dZ$$

$$W = m \omega_0^2 \left(1 + 2 \frac{m}{M}\right)^2 \cdot \frac{MR^2}{4m} \int_0^{\frac{2m}{M}} \frac{dZ}{(1+Z)^2}$$

$$\int_Z = \left[-\frac{1}{1+Z} \right]_0^{\frac{2m}{M}} = \frac{2m/M}{1+2m/M} \Rightarrow$$

$$\underline{W = \frac{1}{2} m \omega_0^2 R^2 \left(1 + \frac{2m}{M}\right)}$$